

# Large Rapidity Gap Events

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- the conventional variables  
are used:

$$\beta = \frac{c^2}{M^2 + Q^2} \quad (t=0)$$
$$x_p = \frac{M^2 + Q^2}{w^2 + Q^2}$$

- closely related work:

[ Bartels,  
Levin,  
Mueller,  
Nikolaev, Zakharov  
Diehl,

⋮

]

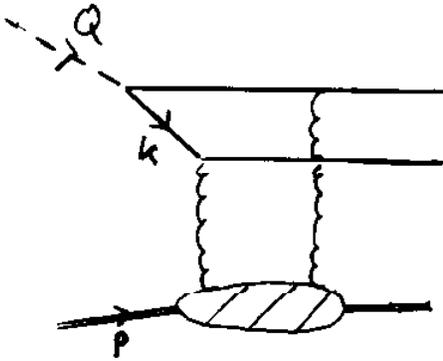
# Dipole Picture

## 1) Quark-Antiquark Dipole :

$r$  : Photon helicity  
 $h$  : Quark helicity

wave function:

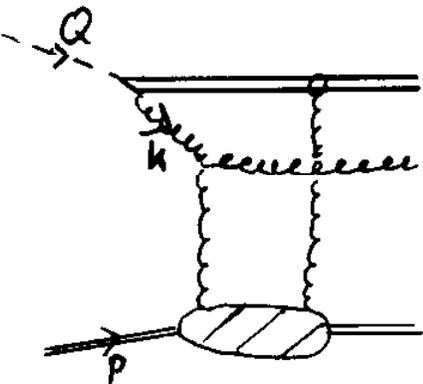
$$\Psi_h^r(\alpha, k_t) = \begin{cases} \frac{\sqrt{2}(\alpha-1)k_t}{|k_t|^2 + \alpha(1-\alpha)Q^2} & (r, h) \\ \frac{\sqrt{2}\alpha k_t}{|k_t|^2 + \alpha(1-\alpha)Q^2} & (+, +) \\ \frac{\sqrt{2}\alpha k_t^*}{|k_t|^2 + \alpha(1-\alpha)Q^2} & (+, -) \\ \frac{\sqrt{2}(\alpha-1)k_t^*}{|k_t|^2 + \alpha(1-\alpha)Q^2} & (-, +) \\ \frac{\sqrt{2}(\alpha-1)k_t^*}{|k_t|^2 + \alpha(1-\alpha)Q^2} & (-, -) \\ 2 \frac{\alpha(1-\alpha)Q}{|k_t|^2 + \alpha(1-\alpha)Q^2} & (0, \pm) \end{cases}$$



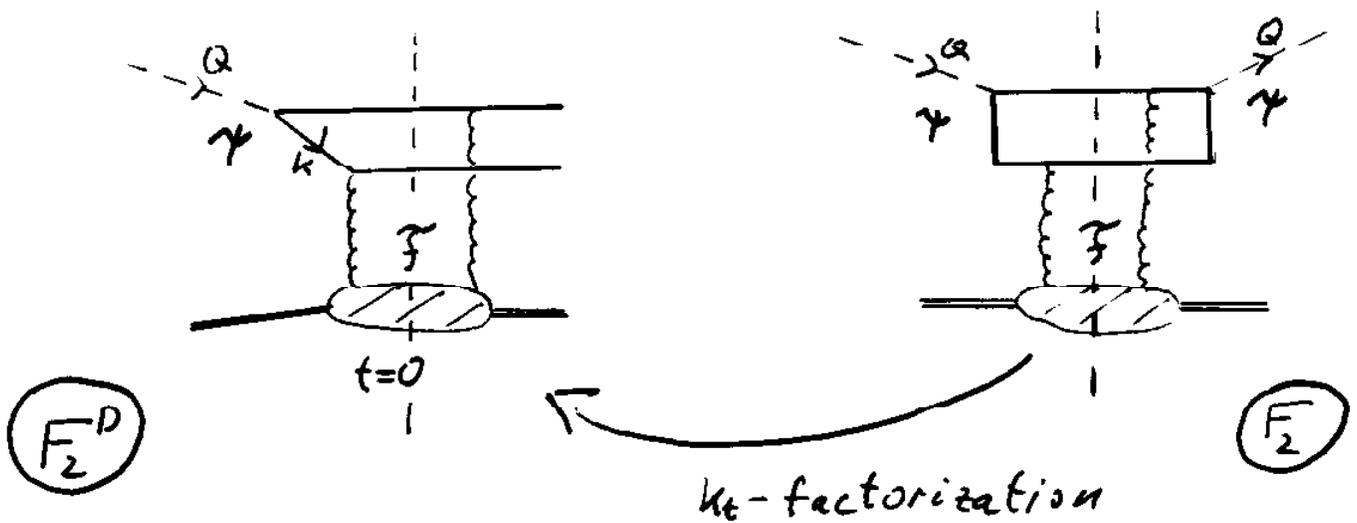
## 2) Effective Gluon-Gluon Dipole :

$$\Psi^{mn}(\alpha, k_t) = \frac{1}{\sqrt{\alpha(1-\alpha)Q^2}} \frac{k_t^2 \delta^{mn} - 2k_t^m k_t^n}{k_t^2 + \alpha(1-\alpha)Q^2}$$

tensor (transverse) ;  $m, n = 1, 2$



# Phenomenology



- the wave function  $\Psi$  has to be folded with the unintegrated structure function  $\mathcal{F}$ .  
 $\rightarrow$  this can be done in momentum space or impact parameter space.

$$F(x_B, l_t^2, Q_0^2) = \frac{G(x_B, Q^2/Q_0^2)}{l_t^2 + Q_0^2}$$

$$G(x_B, Q^2/Q_0^2) = 0.877 \left(\frac{x_0}{x_B}\right)^{1-\alpha_P(Q^2)} \left[\ln\left(\frac{Q^2}{Q_0^2}\right) + 1\right]^{-0.596} \quad \left[\exp(3t)\right] \text{ LPS-data}$$

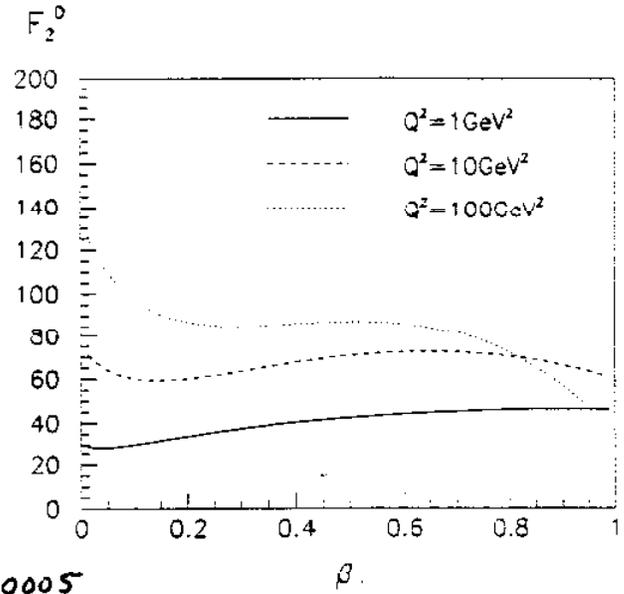
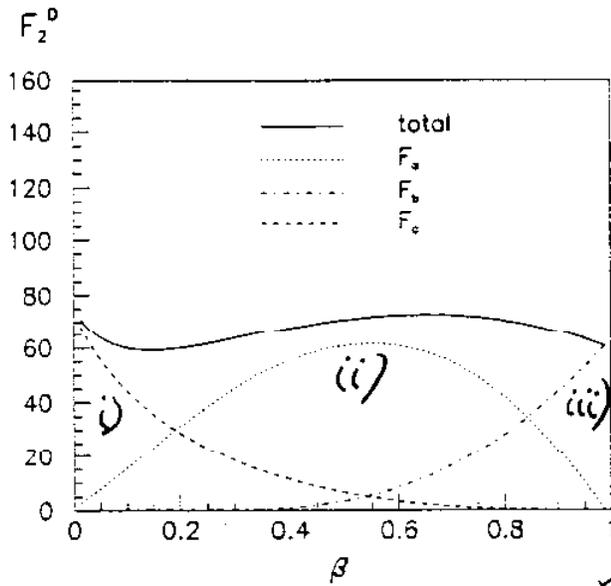
$$\alpha_P(Q^2) = 0.085 + \begin{cases} 0.133 \ln[\ln(Q^2/Q_0^2) + 1] & \text{if } Q^2 > Q_0^2 \\ 0 & \text{if } Q^2 \leq Q_0^2 \end{cases}$$

$$Q_0 = 1 \text{ GeV}$$

- the parameters are determined by a fit to  $F_2$

$\rightarrow$  the scale  $Q^2$  has to be replaced by  $k^2 \equiv \frac{k_t^2}{1-A}$

# $\beta$ -spectrum



$x_p = 0.0005$

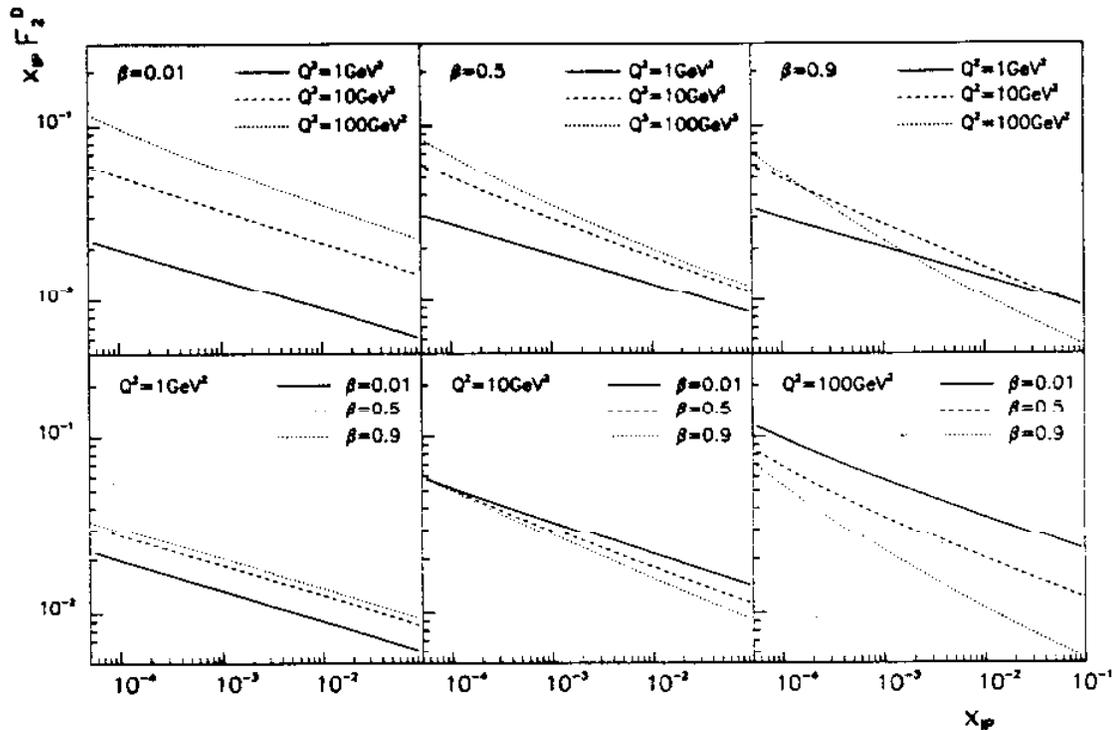
- the  $\beta$ -spectrum is rather flat around  $Q^2 \approx 10 \text{ GeV}^2$

→ it consists of three major contributions:

- i) small  $\beta$ :  $q\bar{q}$  ( $F_E$ )
- ii) medium  $\beta$ :  $q\bar{q}$  (transverse) ( $F_a$ )
- iii) large  $\beta$ :  $q\bar{q}$  (longitudinal) ( $F_b$ )  
↳ higher twist

- with increasing  $Q^2$  the  $\beta$ -distribution starts tilting towards large  $\beta$   
(the low  $\beta$ -regime increases and the large  $\beta$ -regime decreases → higher twist)

# $x_p$ - distribution

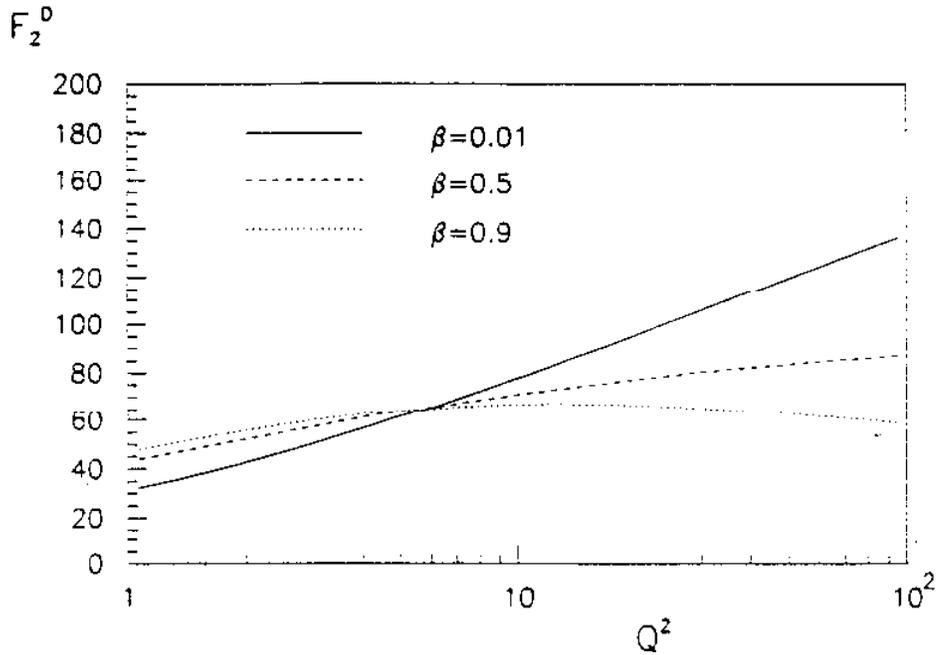


- increase of  $x_p$  with increasing  $\beta$  and  $Q^2$

→ breaking of Regge-factorization

- at low  $\beta$  and  $Q^2$  the Pomeron is purely soft ( $\alpha_P \approx 1.085$ )

## $Q^2$ -dependence



— all curves are rising at low  $Q^2$  and diverge beyond  $Q^2 = 10 \text{ GeV}^2$ :

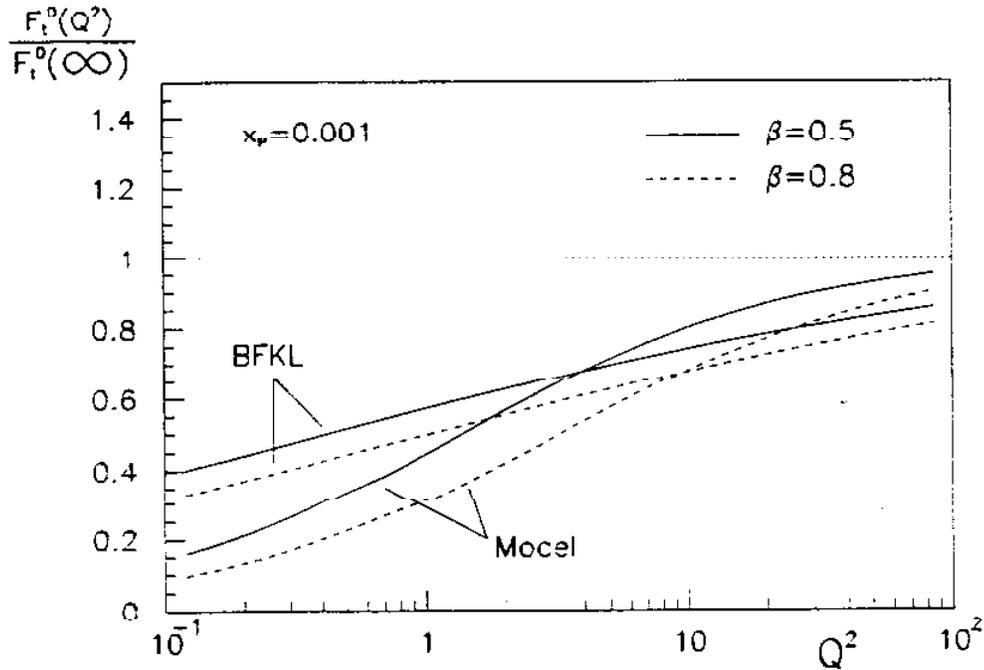
small  $\beta$  : keeps rising

medium  $\beta$  : flattens out

large  $\beta$  : turns down

↓  
higher twist

$q\bar{q}$  - final state with transverse polarized photons



- the asymptotic regime (leading twist) is approached rather slowly.

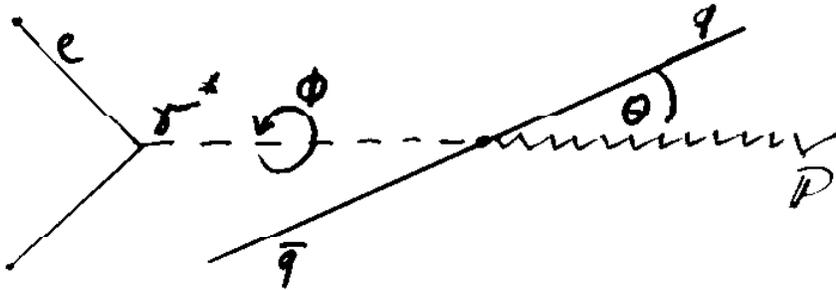
Even at  $Q^2 = 10 \text{ GeV}^2$   $F_1^D$  is 30% below its asymptotic value.

→ preasymptotic regime

→ higher twist corrections

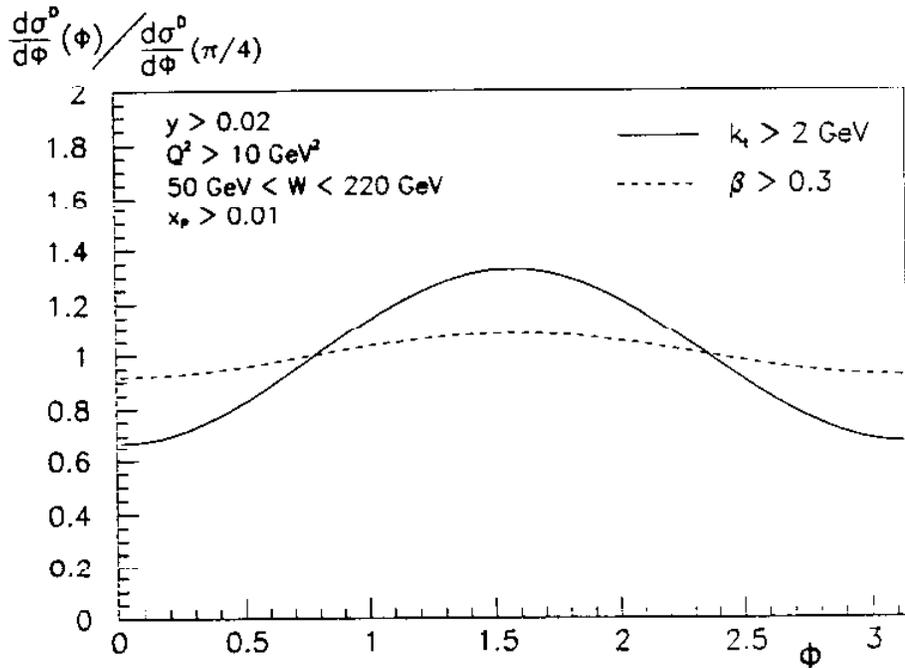
(Bartels)

# Azimuthal - distribution



$\phi$  = azimuthal angle

$\theta$  = polar angle



————— : with a cut on  $k_t$  a strong oscillation is seen (peak at  $\frac{\pi}{2}$ )

- - - - : a slight oscillation can also be seen in the inclusive case ( $\beta > 0.3$ )

## Conclusions

- rather flat total  $\beta$ -spectrum
- breaking of Regge-factorization
  - increasing  $\kappa_D$  with increasing  $\beta$  and  $Q^2$
- \* positive slope in  $Q^2$ 
  - preasymptotic regime  
(higher twist effects)
- azimuthal distribution with peak at  $\frac{\pi}{2}$