

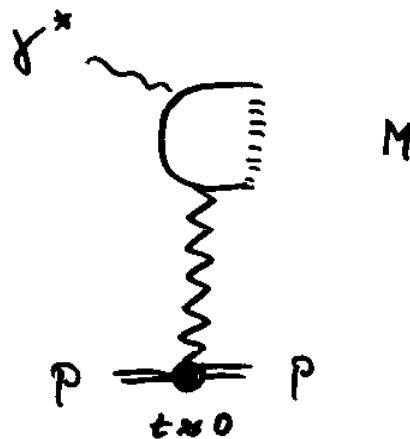
DIS 97, Chicago April '97
J. Bartels

Higher Twist in

Diffractive Dissociation

Motivation for studying higher twist in DIS diff. dis.

Experimental evidence that



contains also "hard Poweron": $\gamma^* p \rightarrow Vp$, $\gamma^* p \rightarrow jet + p$.

How to separate "soft" and "hard" components in the integrated diffractive cross section: by

expansion in powers of $1/Q^2$ ("twist expansion")

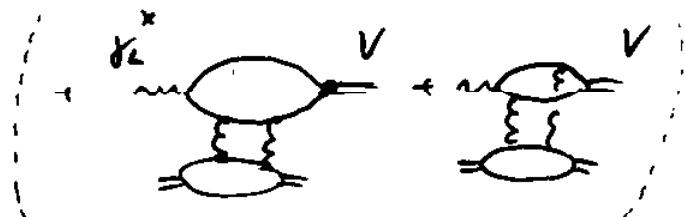
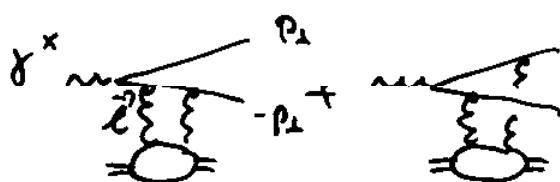
$$\frac{d^4\sigma}{dx dQ^2 dx_p dt} \Big|_{t=0} = \frac{4\pi \alpha_{em}}{x Q^4}$$

$$\cdot \left\{ \frac{1 + (1-y)^2}{2} \left(\bar{F}_L^D + \frac{m^2}{Q^2} \Delta \bar{F}_L^D \right) + (1-y) \left(\bar{F}_L^D + \frac{m^2}{Q^2} \Delta \bar{F}_L^D \right) \right\}$$

How to address this question:

- $q\bar{q}$ production (reliable for small x_{scat} , $\beta \approx 1$)
- $q\bar{s}$ ($u\bar{g}$) in the region $Q^2 \ll M^2 \ll W^2$

A. Diffractive $q\bar{q}$ -production:



Start from large- p_t region where ρ QCD is applicable
and extrapolate down to finite/small p_t^2 :

JTB, Ellis, Kowalski,
Wüsthoff

$$\frac{d\sigma_T}{dt^2 dp_t^2 dt} \Big|_{t=0} \approx \sum_f e_f^2 \frac{\alpha_m \pi^2 \alpha_s^2}{3} \frac{1 - \frac{2p_t^2}{M^2}}{\sqrt{1 - \frac{4p_t^2}{M^2}}} \frac{4 M^4 \alpha^4}{(M^2 - t \alpha^2)^6} \left[x_p g(x_p, \frac{p_t^2}{t-p}) \frac{1}{p_t^2} \right]^2$$

Integral over p_t^2 (at fixed β): dominance of small p_t^2 ;
AJM configuration; $(1-\beta)$ factor.

Expect soft (nonperturbative) Powers

Expand in powers of $1/\alpha^2$:

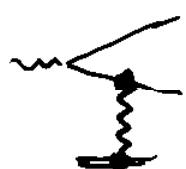
leading-twist: nonperturbative, needs model assumption

Buchmiller
et al

e.g. unintegrated gluon structure function

$$\phi \sim \frac{1}{k_0^2} \left(\frac{k_0^2}{\ell^2} \right)^{v(\Delta)}, \quad v \begin{cases} 1 & \text{large quark virtuality} \\ 0 & \text{small " " } \end{cases}$$

nonplanar diagram << planar diagram:



→ picture of Powers structure function emerges

$$\left. \frac{d^2\sigma_t}{dt dM^2} \right|_{t=0} \sim \frac{1}{Q^6} \beta^3 (1-\beta) \left[x_F g(x_F, \frac{Q^2}{4\beta}) \right]^2$$

- large ℓ^2 , but no $\log Q^2$ calculable in pQCD
- vanishes as $1-\beta$
- enhancement due to square of gluon structure function
"hard Pomeran contribution"

The same for the longitudinal cross section:

$$\left. \frac{d\sigma_L}{dM^2 dp_T^2 dt} \right|_{t=0} \approx \sum e_f^2 \frac{\alpha_{em} \pi \bar{s}^2}{3} \frac{4}{\sqrt{1 - \frac{4p_T^2}{M^2}}} \frac{\alpha^2 (Q^2 - M^2)^2}{(M^2 + Q^2)^6} \frac{1}{p_T^2} \left[x_F g(x_F, \frac{p_T^2}{1-\beta}) \right]^2$$

Integral over p_T^2 : no enhancement at small p_T
Expansion starts with higher twist (twist 4)
no twist-2 contribution

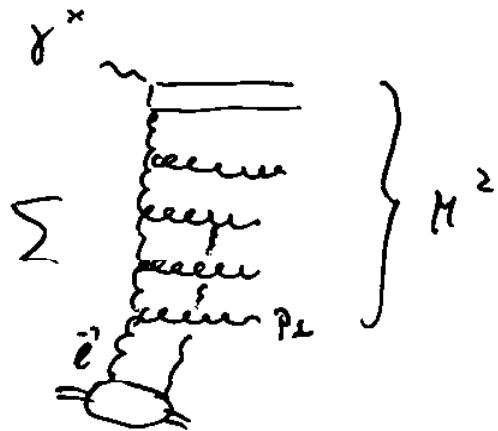
Twist-4 term:

$$\left. \frac{d^2\sigma_L}{dt dM^2} \right|_{t=0} \sim \frac{1}{Q^6} \ln \frac{Q^2}{4\beta} \cdot \left[x_F g(x_F, \frac{Q^2}{4\beta}) \right]^2$$

- extra $\log Q^2$, strong ordering in transverse momenta
e.g. $\ell^2 \ll p_T^2 \ll Q^2$; calculable in p8
- nonvanishing at $\beta = 1$
- enhancement due to square of gluon structure function: "hard Pomeran contribution".

Q^2 -evolution

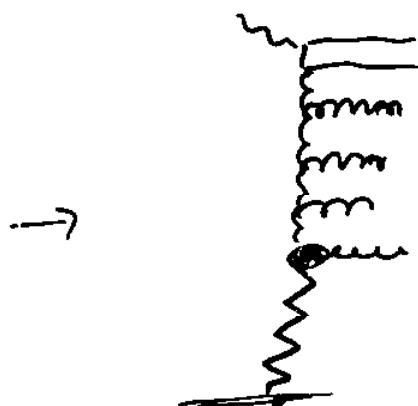
Argument based upon $\gamma^* p \rightarrow \underbrace{(q\bar{q} = u\bar{d})}_{\mu^2} p$, in the region
 $Q^2 < M^2$ (small β)



Leading-twist:

- soft $\vec{\ell}^2 \rightarrow$ nonperturbative Power
- DGLAP - evolution

Soper, Sereva
Kuraev



Kwieciński et al
Golec-Biernat + Stiriba
H1
CERN NA3
..

Twist-4: very different characteristics

- $\langle \ln \frac{\bar{Q}^2}{Q_0^2} \rangle \sim \langle \ln \frac{Q^2}{Q_0^2} \rangle \cdot \xi(W^2, M^2), \quad 0 < \xi < 1$

→ Power ~ gluon structure function

→ calculable in pQCD!

Lipatov

- QCD evolution for twist 4: nonlinearward DGLAP-splitting functions, new kernels



- interesting possibility:

evolution equation has inhomogeneous term

\neq Lenz et al.

$$\frac{\partial f(\beta, q^2)}{\partial \ln q^2} = f_0(\beta, q^2) + (K \otimes f)(\beta, q^2)$$

↑ ↑
 twist-4 evolution kernels

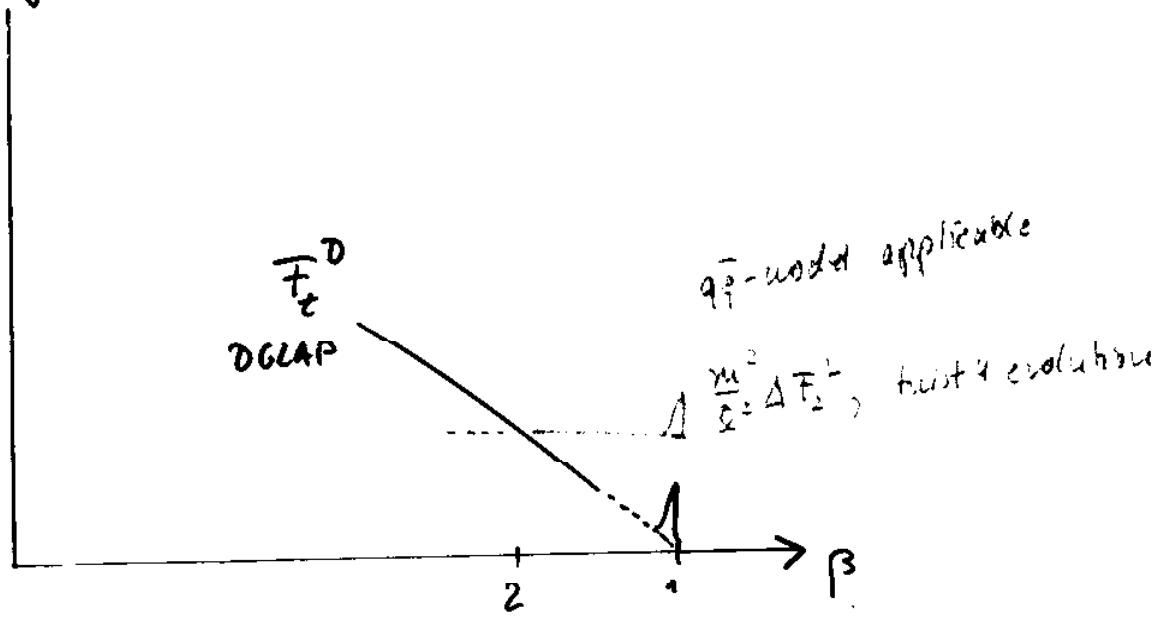
$$\sim [x_F g(x_F, q^2)]^2$$

(analogy with pointlike photon?)

- need to calculate twist-4 evolution kernels and mixing of twist-4 operators



Summary: expect



- $1/\alpha_2$ -expansion good tool for separating "soft" and "hard" components in \bar{F}_t^D
- twist-4 visible in \bar{F}_t^D near $\beta=1$
- analysis of data should include \bar{F}_t , higher twist,
in particular near $\beta=1$. → Wüsthoff's talk