

# DIS97

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## EXCLUSIVE VECTOR MESON PRODUCTION AT HERA

Frank Gaede  
MPI für Physik München

for the H1 Collaboration



- introduction
- experimental results in photoproduction and at high  $Q^2$
- conclusion

## introduction

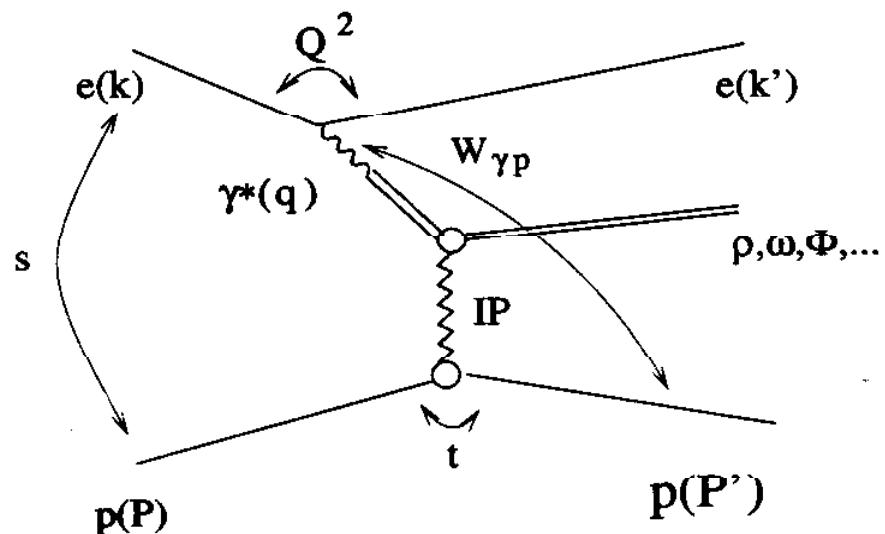
- different theoretical predictions for the dependence on kinematic variables of the differential cross sections
  - phenomenological model:  
Regge theory – pomeron exchange
  - perturbative QCD calculations
- which is the appropriate description at HERA energies ( $W_{\gamma p} \approx 100 \text{ GeV}$ ) ?
- analysed decay modes

$$\begin{aligned}\rho^0 &\rightarrow \pi^+ \pi^- \\ (\omega &\rightarrow \pi^+ \pi^- \pi^0) \\ \phi &\rightarrow K^+ K^- \\ J/\Psi &\rightarrow l^+ l^- (l = e, \mu) \\ \Psi' &\rightarrow J/\Psi \pi^+ \pi^- \\ \rho' &\rightarrow \pi^+ \pi^- \pi^+ \pi^-\end{aligned}$$

- study transition region from soft to hard physics on two scales:  
mass of the vector meson  
virtuality of the photon  $Q^2$

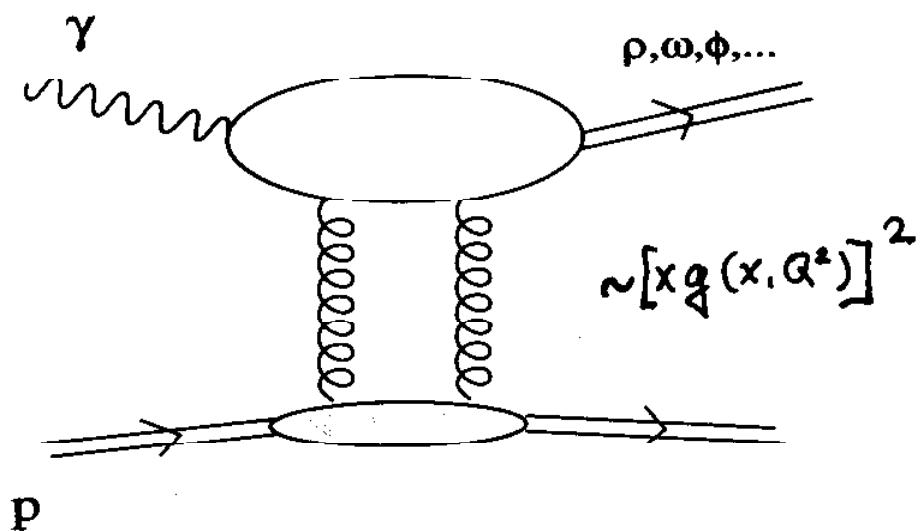
## kinematics and production mechanisms

- Regge model - pomeron exchange (VDM)



(Donnachie - Landshoff)

- pQCD models



(Brodsky et.al., Ryskin, ...)

## dependence on kinematic variables

- $W_{\gamma p}$

$$\frac{d\sigma}{dt} \Big|_{t=0} \propto W_{\gamma p}^{4*0.0808} \quad (\text{REGGE})$$

$$\frac{d\sigma}{dt} \Big|_{t=0} \propto \frac{[x g(x, Q^2)]^2}{Q^6} \quad (\text{pQCD})$$

$$W_{\gamma p}^2 = \frac{1-x}{x} Q^2 \rightarrow \propto W_{\gamma p} \approx 4*0.2$$

- $Q^2$

$$\frac{d\sigma}{dQ^2} \propto (m_{VM}^2 + Q^2)^{-n} \quad (n \approx 2 - 3)$$

- $t$

$$\frac{d\sigma}{dt} \propto e^{-b|t|}$$

$$b = b_0 + 4\alpha' \ln \frac{W_{\gamma p}}{W_0} \quad (\text{REGGE})$$

→ shrinkage

(pQCD) little prediction, if then small shrinkage

- decay angular distribution

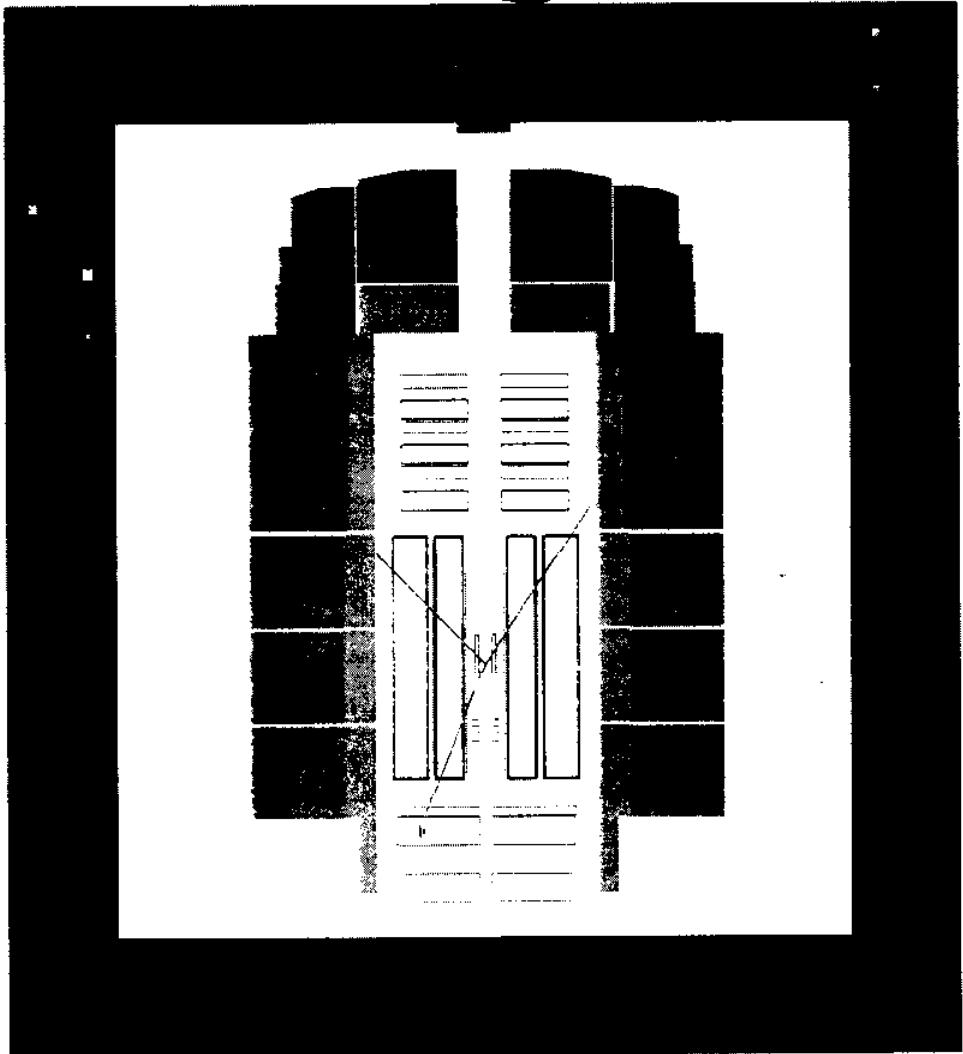
$$\frac{d\sigma}{d \cos(\Theta^*)} \propto (3 r_{00}^{04} - 1) \cos^2(\Theta^*) + (1 - r_{00}^{04})$$

(SCHC)

 Run 120704 Event 3250 Class: 3 11 15 20 22 24 Date 16/05/1996

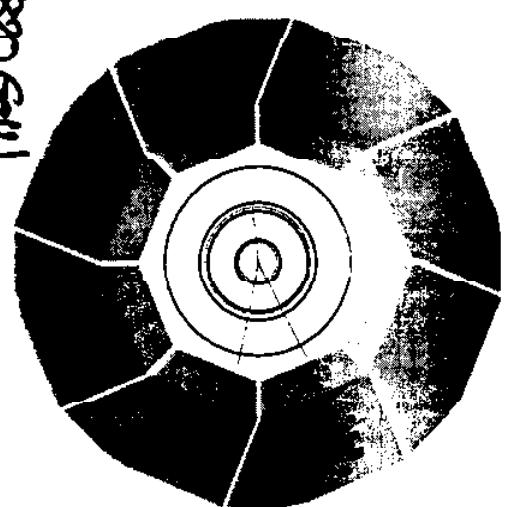
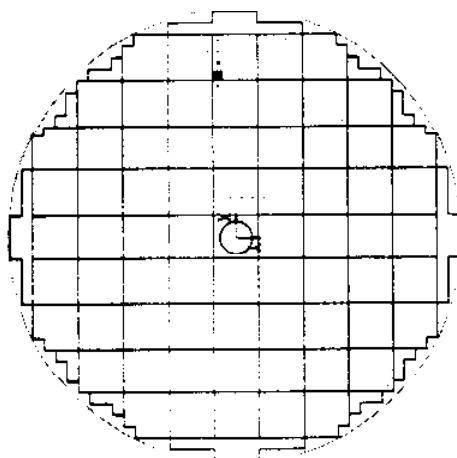
$\frac{3}{4} \rightarrow \mu^+ \mu^-$

(ERSS)



$\bar{Z}$   
R

P(80 GeV)

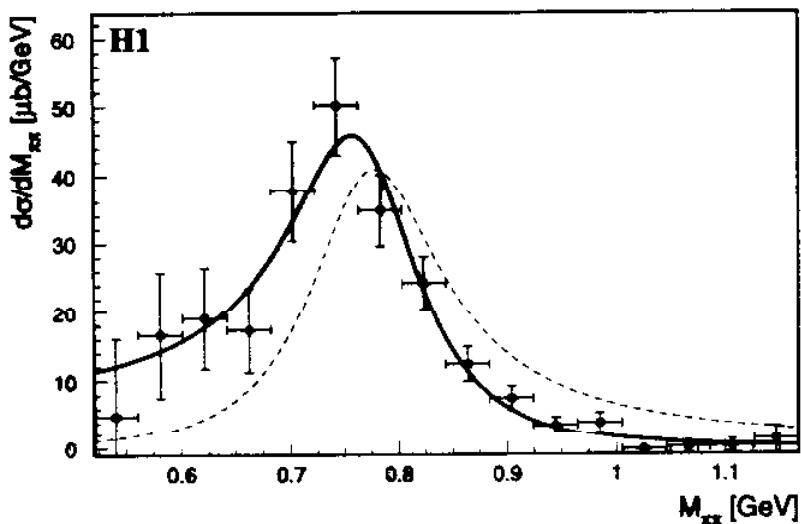


## common analysis cuts

- exact two charged tracks in central tracker with unlike signs, fitted to the primary vertex (+ 1 track for the scattered electron)
- no additional energy deposition in calorimeters
- forward cuts against proton dissociation  
no signal (above noise) in forward detectors:
  - proton tagger (+24 m)
  - forward muon system
  - forward part of LAr
$$\left. \begin{array}{l} \text{– proton tagger (+24 m)} \\ \text{– forward muon system} \\ \text{– forward part of LAr} \end{array} \right\} \Rightarrow M_Y \tilde{<} 1.6 \text{ GeV}$$
- **t cut:**  
 $|t| < 0.5 - 0.6 \text{ GeV}^2$
- for high  $Q^2$  events:
  - good electromagnetic shower in backward calorimeter
  - associated hit in backward track chamber
  - cut against radiation:  
 $\sum_i E_i - p_{zi} > 45 - 48 \text{ GeV}$

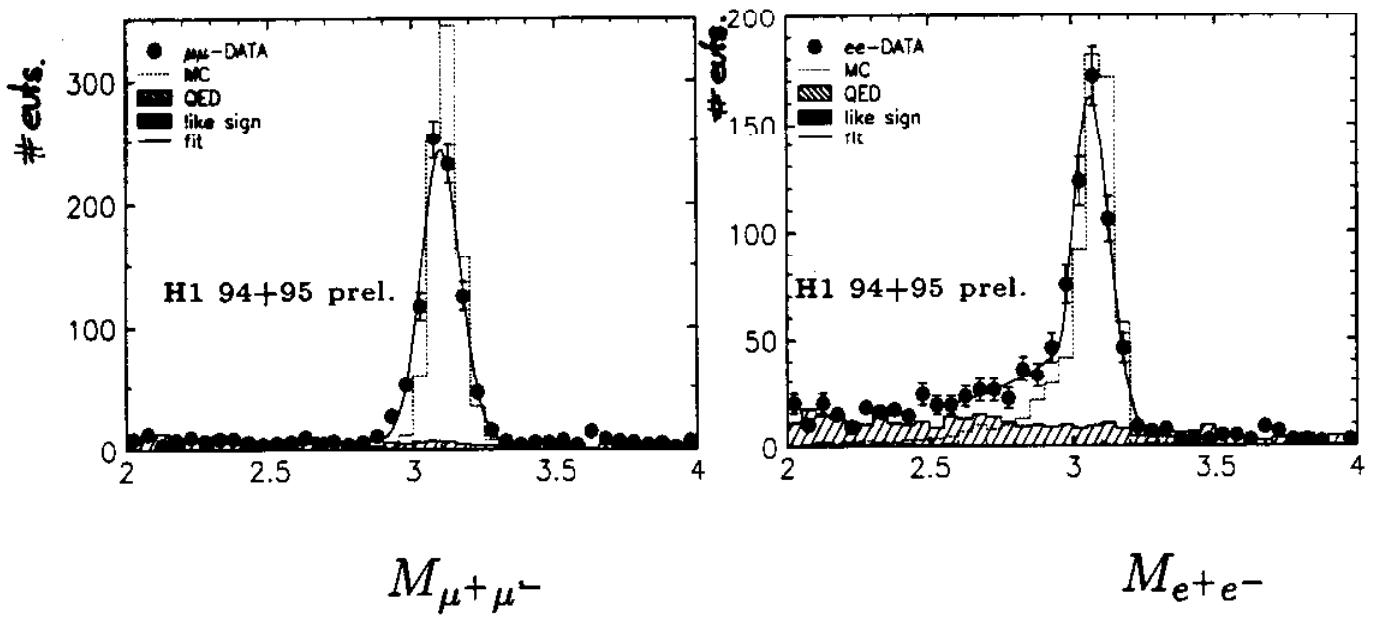
# mass signals in photoproduction

- $\rho^0 \rightarrow \pi^+ \pi^-$

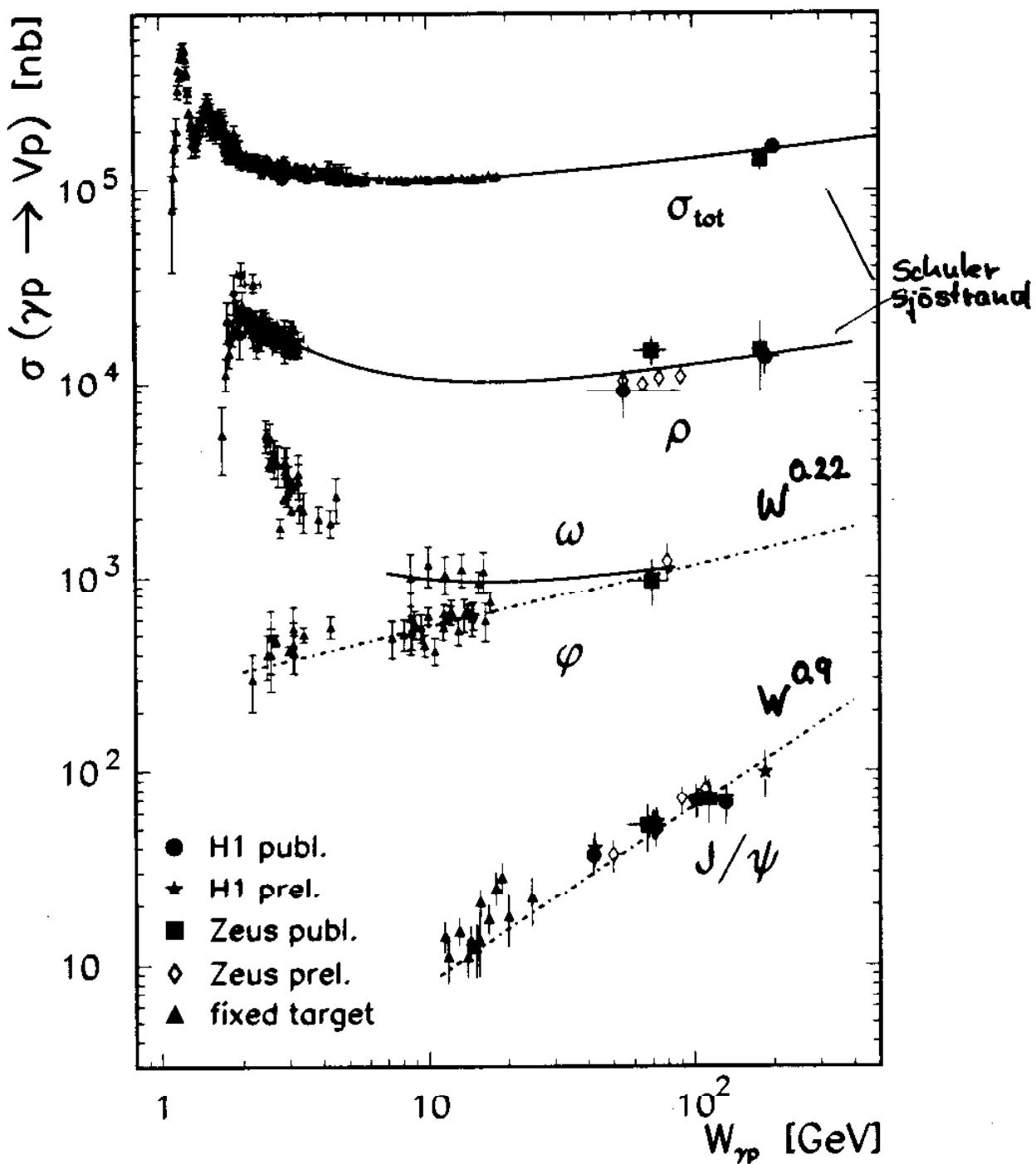


skewed mass (Ross-Stodolsky)

- $J/\Psi \rightarrow l^+ l^-$

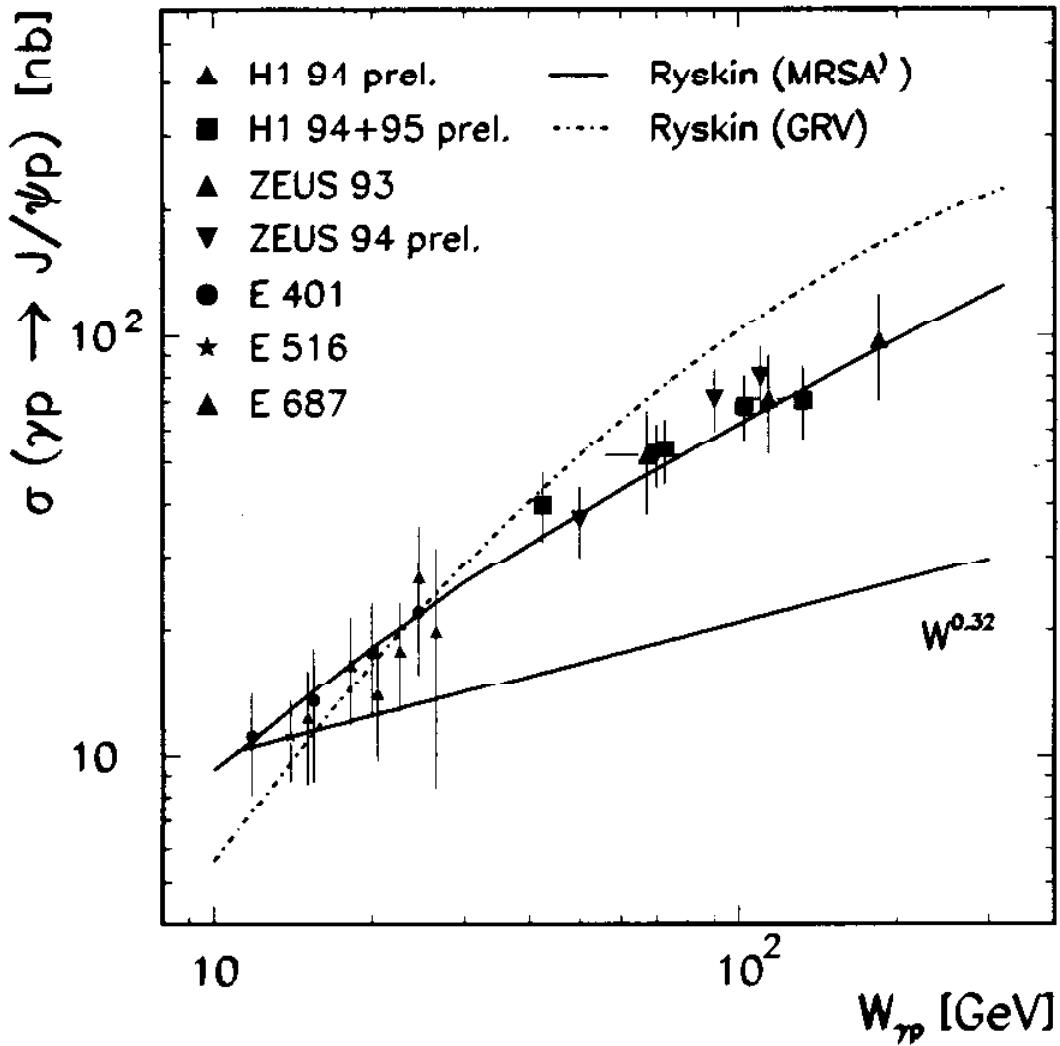


# photoproduction cross sections



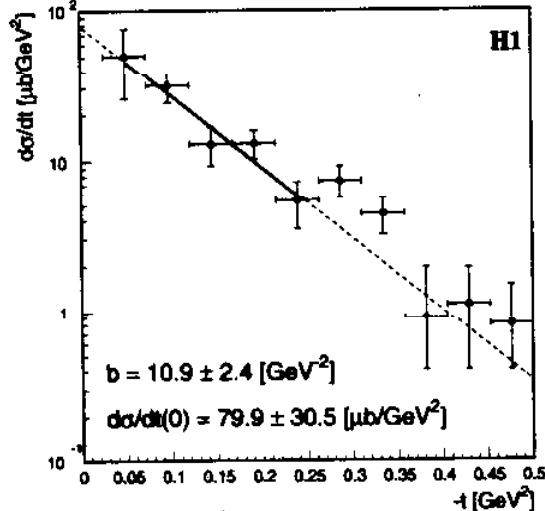
- light vector mesons described by soft pomeron
- $J/\Psi$  steeper rise  $\rightarrow$  'hard' pomeron

# pQCD description of $J/\Psi$ photoproduction

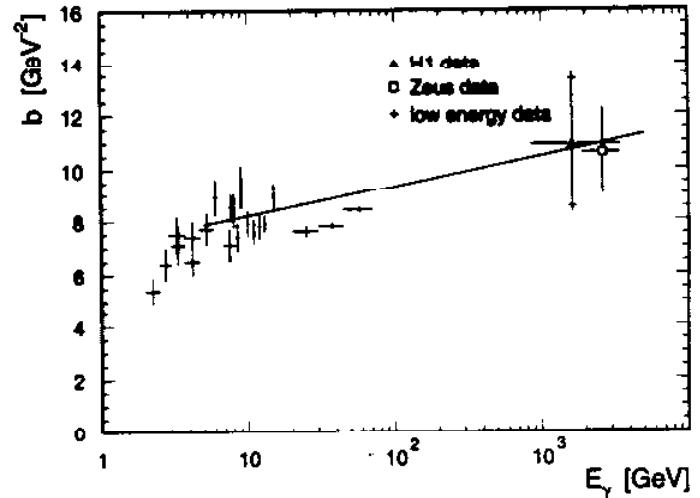


- Ryskin model with MRS( $A'$ ) parametrisation of the gluon density agrees well with the data  
(normalisation uncertain  $\rightarrow$  compare shape only)
- hard scale due to large charm quark mass

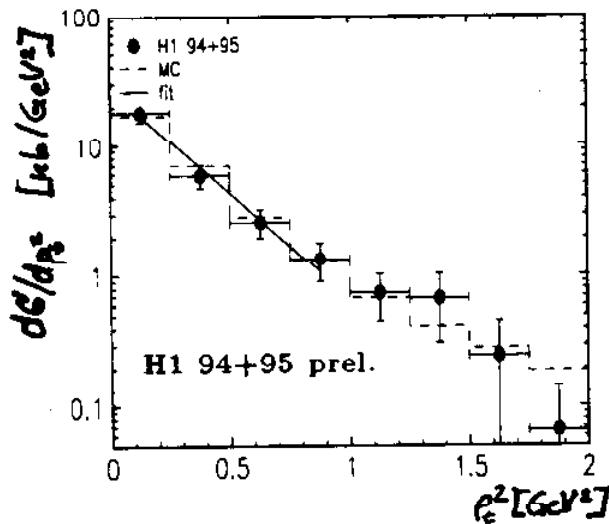
# $t$ distribution in photoproduction



$\rho^0$  : diffractive peak

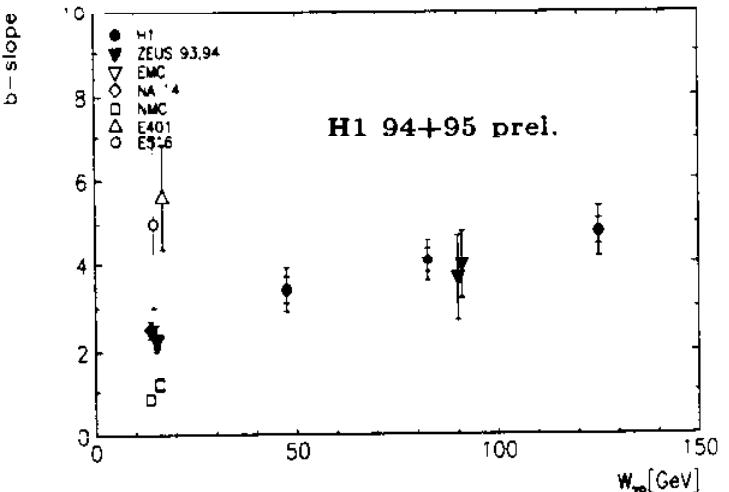


shrinkage



$J/\Psi$  :

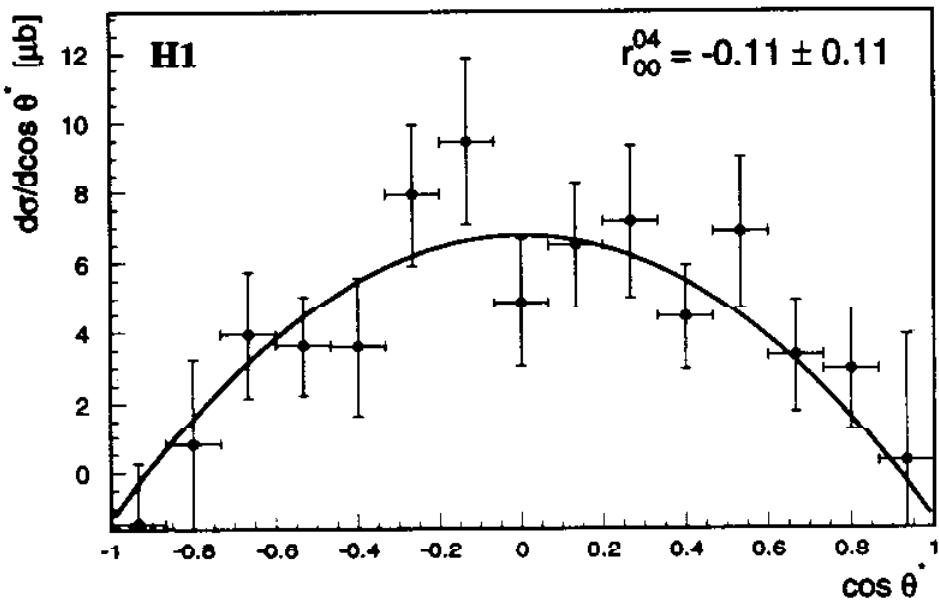
$$b = 4.1 \pm 0.2 \pm 0.45 \text{ GeV}^{-2}$$



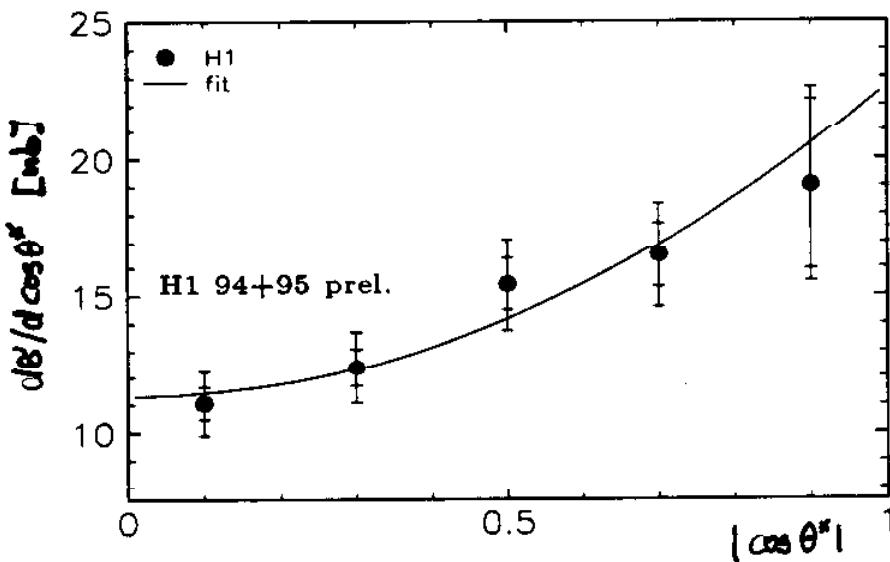
indication for shrinkage ?

$$\text{H1: } \alpha' = 0.36 \pm 0.21 \text{ GeV}^{-2}$$

## decay angular distribution



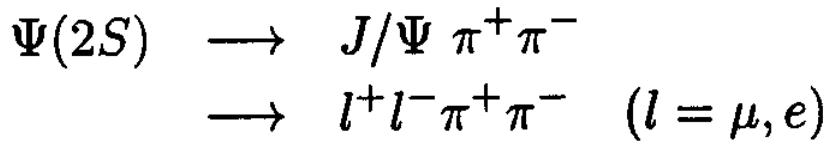
$$\rho^0 : r_{00}^{04} = -0.11 \pm 0.11$$



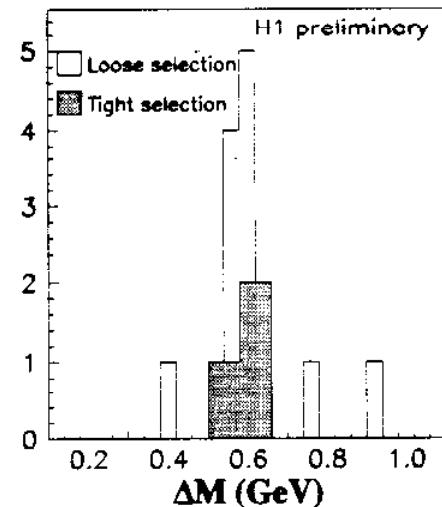
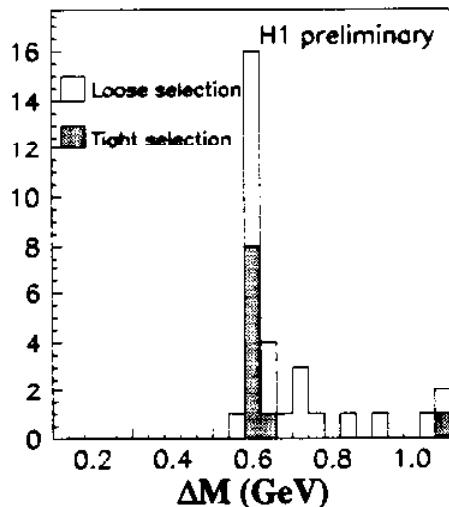
$$J/\Psi : r_{00}^{04} = -0.02 \pm 0.07$$

consistent with SCHC !

# photoproduction of $\Psi(2S)$ mesons



$$(\text{br} = 1.9 \pm 0.2 \%)$$



$$\mu : \sigma_{\gamma p} = 19 \pm 6 \pm 4 \text{ nb}$$

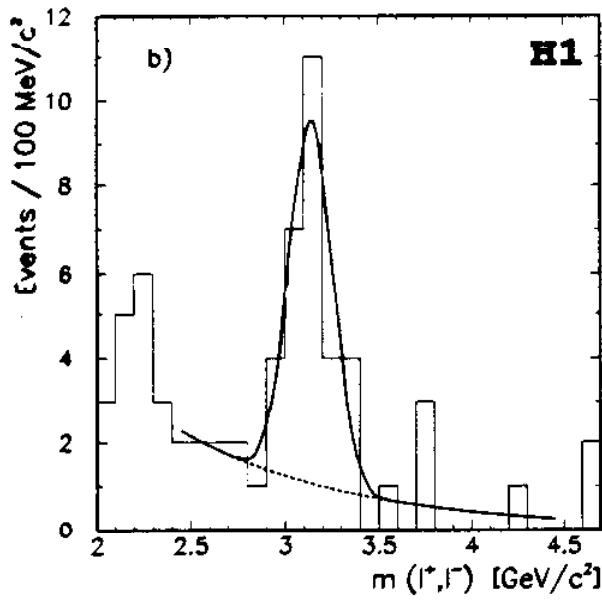
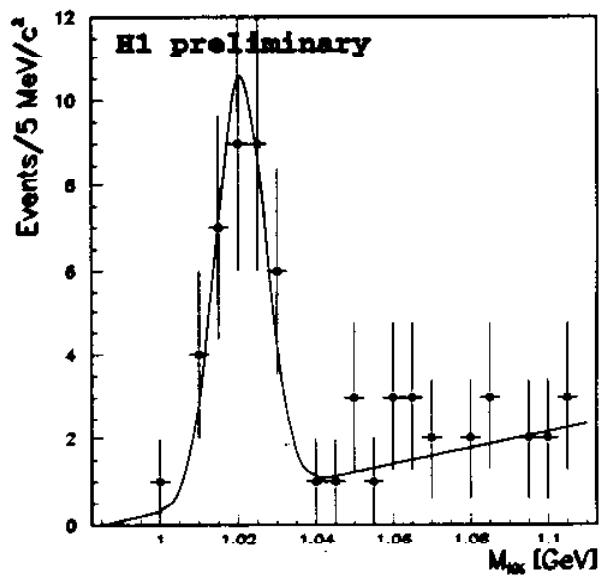
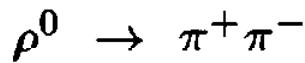
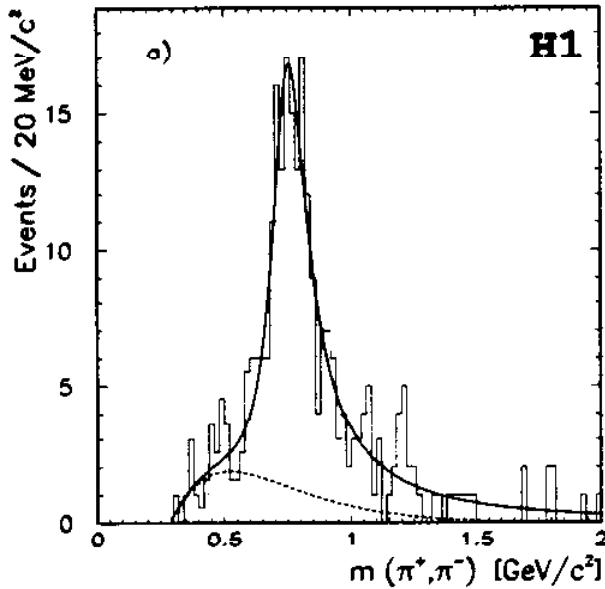
$$e : \sigma_{\gamma p} = 18 \pm 7 \pm 4 \text{ nb}$$

$$\Rightarrow \sigma_{\gamma p} = 19 \pm 5 \pm 4 \text{ nb}$$

$$\Rightarrow \frac{\sigma_{\gamma p}(\Psi(2S))}{\sigma_{\gamma p}(J/\Psi)} = 0.16 \pm 0.06$$

similar result in low energy experiments  
 $\Rightarrow$  no strong energy dependence of the ratio

# mass signals at high $Q^2$



$\rho^0$  :

rel. BW (m.d.w.) + BG  
( no skewing )

$\Phi$  :

rel. BW  $\otimes$  Gauss + BG

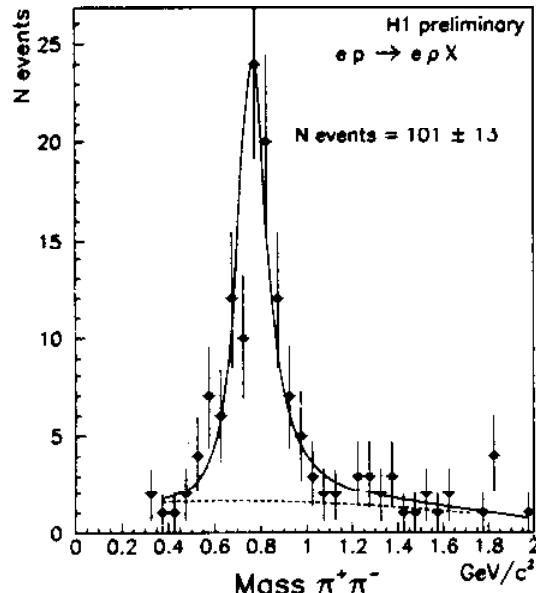
$J/\Psi$  :

Gauss + BG

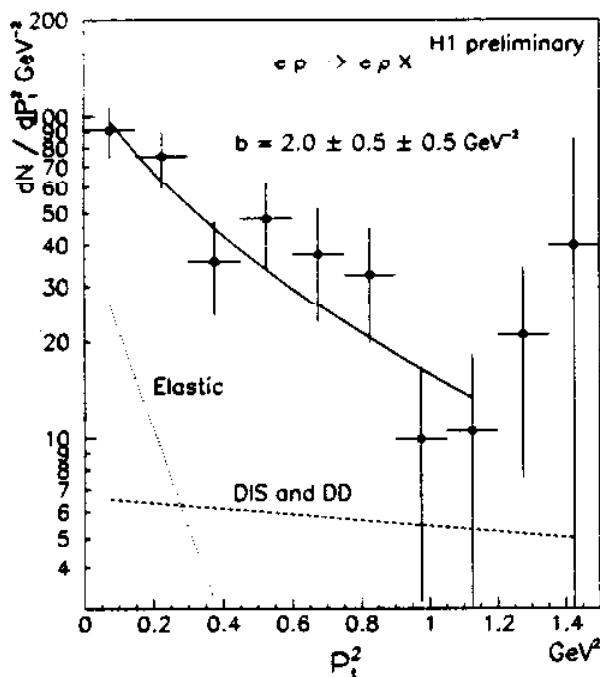
$Q^2 > 8 \text{ GeV}^2$  ( $> 6 \text{ GeV}^2$  for  $\Phi$ )

# $\rho^0$ electroproduction with proton dissociation

proton dissociation with forward detectors  $\rightarrow M_Y > 1.6 \text{ GeV}$



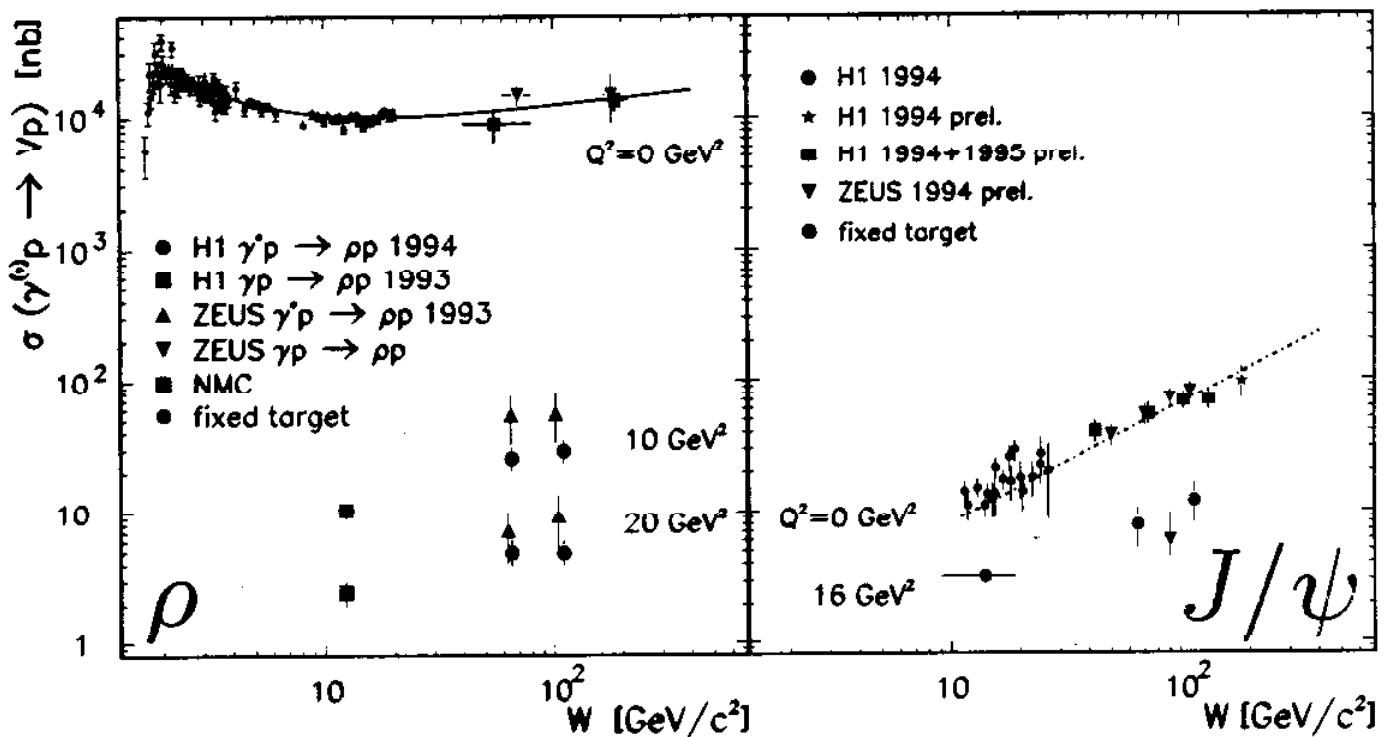
- $7 < Q^2 < 36 \text{ GeV}^2$
- $60 < W_{\gamma p} < 180 \text{ GeV}$
- $\sigma_{\gamma^* p \rightarrow \rho Y} / \sigma_{\gamma^* p \rightarrow \rho p} = 0.65 \pm 0.16 \pm 0.12$
- no dependence on  $W_{\gamma p}$  and  $Q^2$  observed



- $b_{\text{elastic}} / b_{\text{pdis}} \simeq 3.5$
- similar result as in  $J/\Psi$  photoproduction

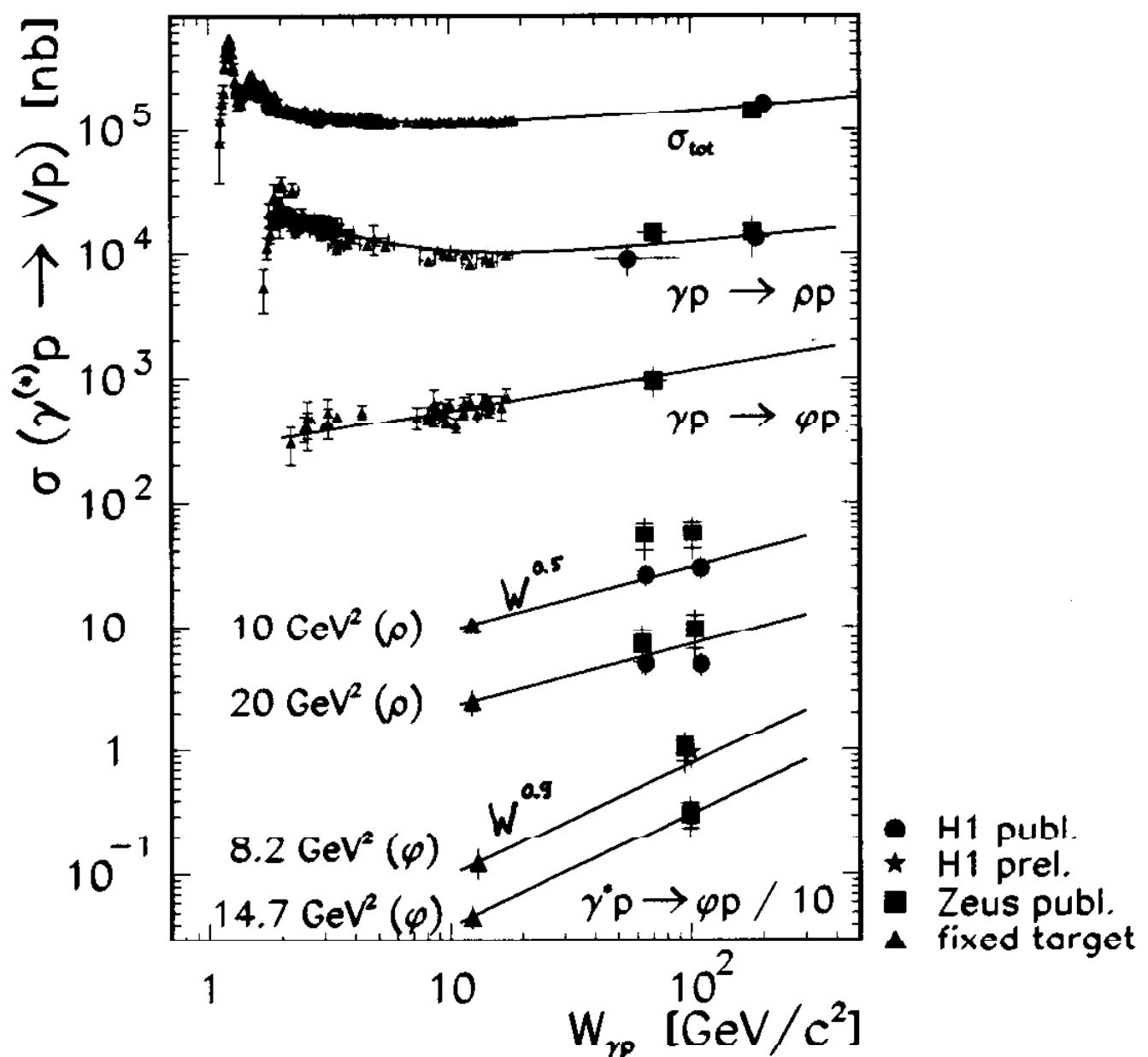
not shown:  $Q^2$  and  $\cos(\theta^*)$  dependence compatible within errors to elastic production

# $W_{\gamma p}$ dependence $\rho^0$ vs. $J/\Psi$ at high $Q^2$



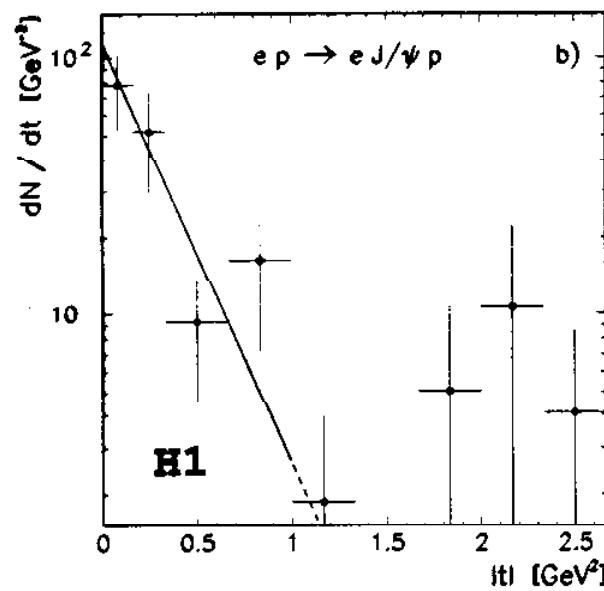
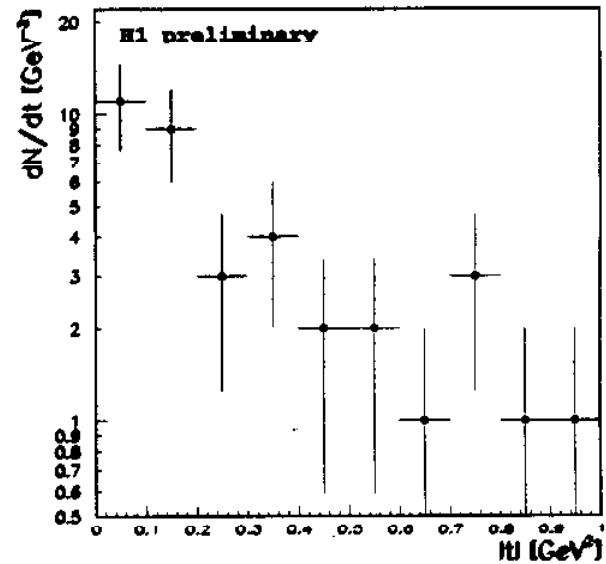
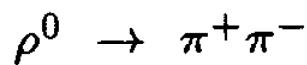
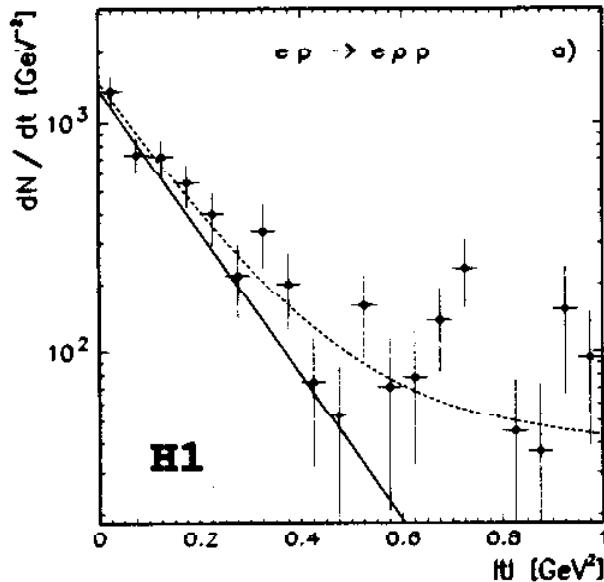
- $\rho^0$  : situation unclear
  - discrepancy between H1 and ZEUS
  - new E665 result (at  $Q^2 \approx 6 \text{ GeV}^2$ ) indicates consistency with soft pomeron
- soft pomeron not ruled out
- $J/\Psi$  : same as in photoproduction

## $W_{\gamma p}$ dependence $\rho^0$ vs. $\Phi$



- $\Phi$  : steeper rise with  $W_{\gamma p}$  than in photoproduction
- increase larger for  $\Phi$  than for  $\rho^0$

# t dependence at high $Q^2$



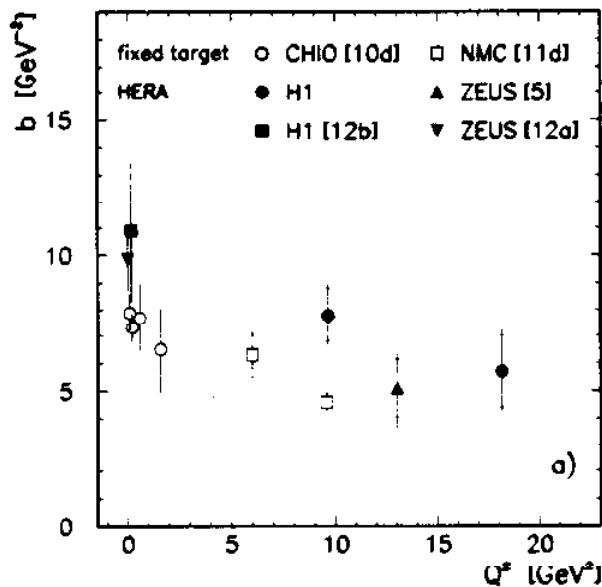
$$\rho^0 : b = 7.0 \pm 0.8 \pm 0.4 \pm 0.5 \text{ GeV}^2$$

$$\Phi : b = 5.2 \pm 1.6 \pm 1.0 \text{ GeV}^2$$

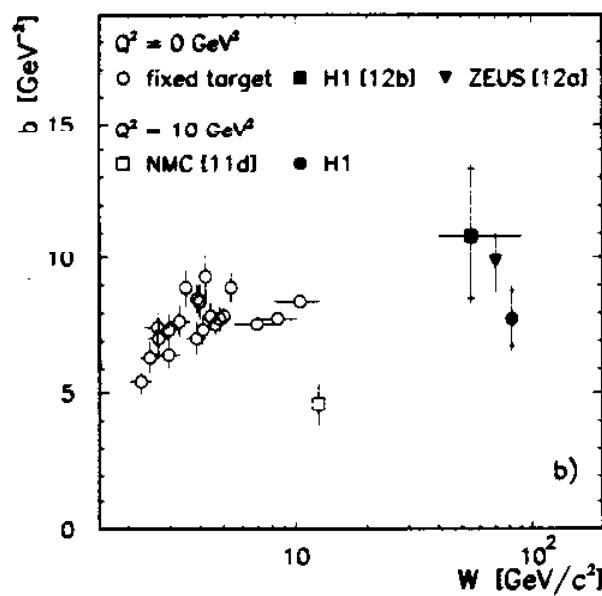
$$J/\Psi : b = 3.8 \pm 1.2^{+2.0}_{-1.6} \text{ GeV}^2$$

## evolution of the $t$ slope at high $Q^2$

- for  $\rho^0$

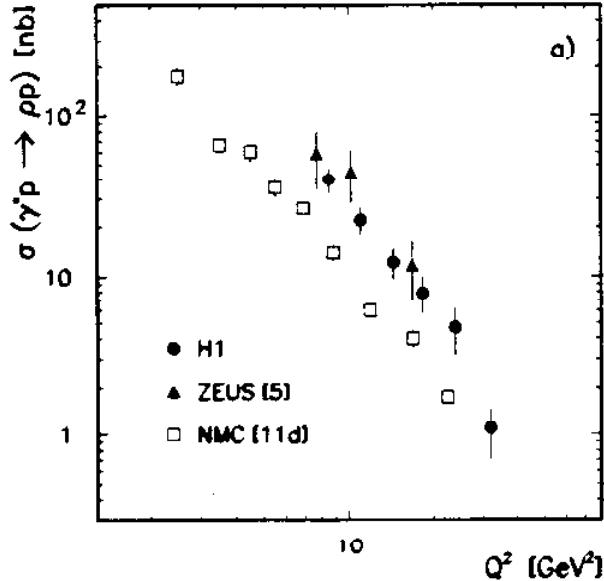


decrease with  $Q^2$   
as seen for low energy

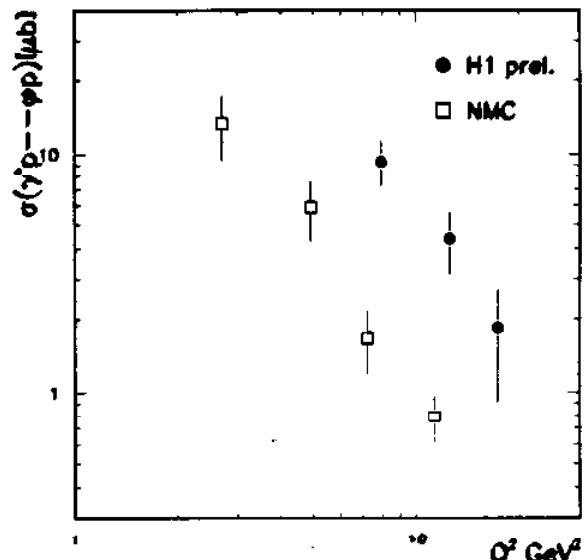


shrinkage with  $W_{\gamma p}$ ?  
relies on fixed targed point

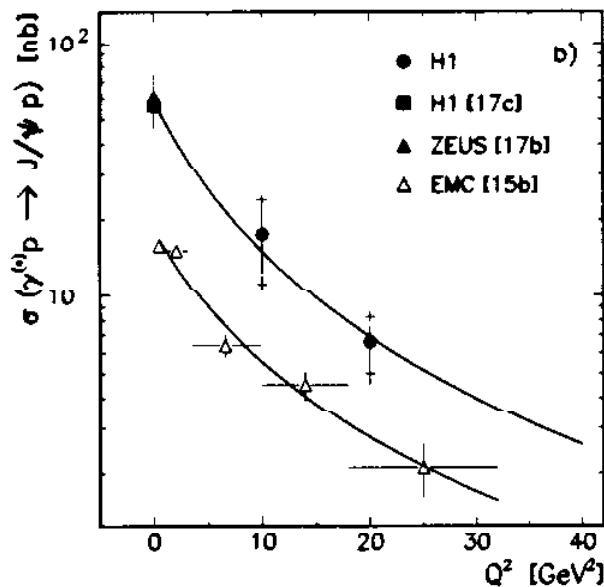
# $Q^2$ dependence at high $Q^2$



$$\rho^0 : \propto Q^{-2n}$$



$$\Phi : \propto Q^{-2n}$$



$$\rho^0 : n = 2.5 \pm 0.5 \pm 0.2$$

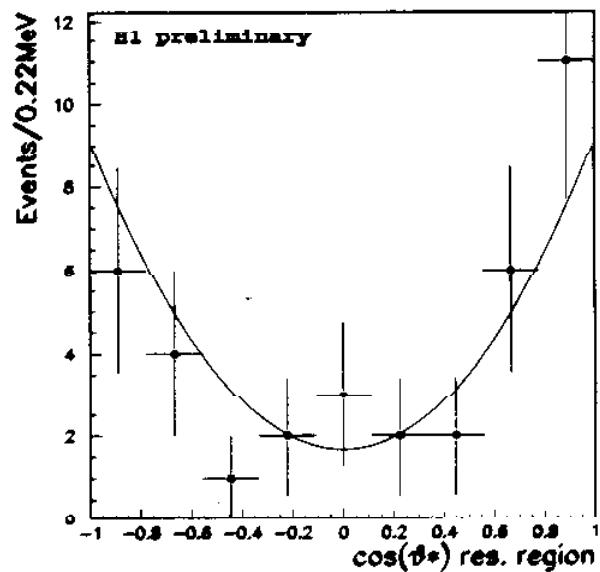
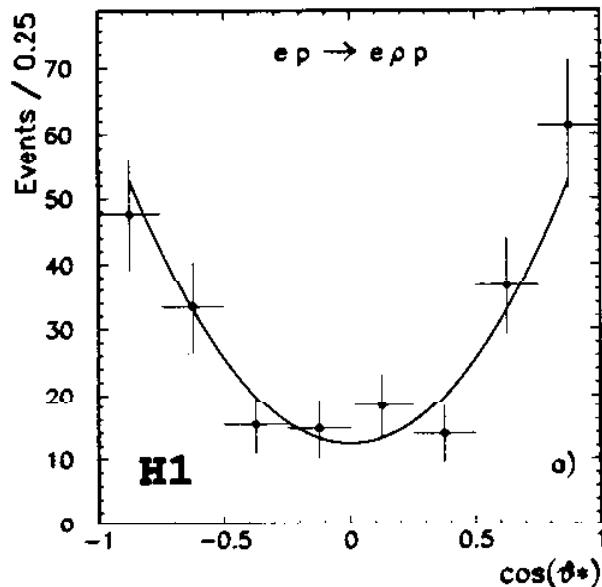
$$\Phi : n = 2.0 \pm 0.6 \pm 0.2$$

$$J/\Psi : n = 1.9 \pm 0.3$$

$$J/\Psi : \propto 1/(Q^2 + m_\Psi^2)^n$$

good agreement with low energy data

# decay angular distribution at high $Q^2$



$$\rho^0 : r_{00}^{04} = 0.73 \pm 0.05 \pm 0.02$$

$$\Phi : r_{00}^{04} = 0.77 \pm 0.11 \pm 0.01$$

assuming SCHC  $\Rightarrow$

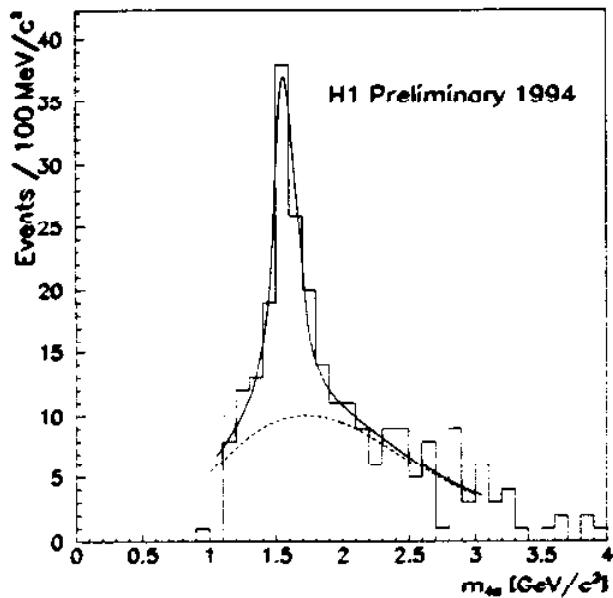
$$R = \sigma_L / \sigma_T = 2.7^{+0.7}_{-0.5} {}^{+0.3}_{-0.2}$$

$$R = 3.3^{+4.0}_{-1.4}$$

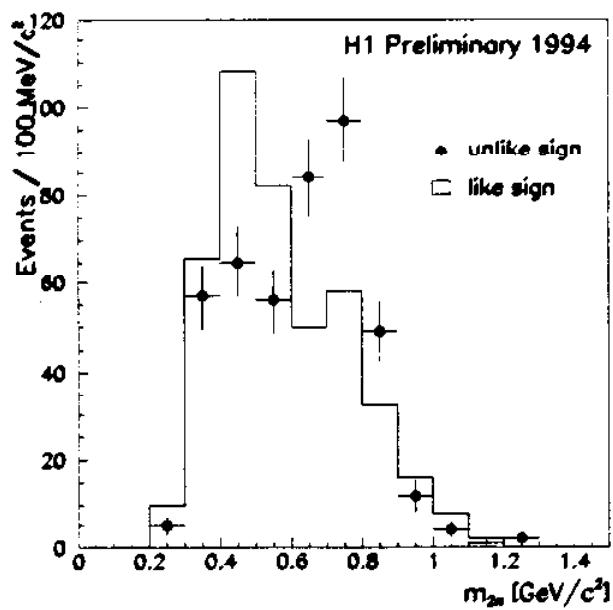
at high  $Q^2$  VM are predominantly longitudinally polarized

## $\rho'$ signal at high $Q^2$

$$\rho' \rightarrow (\rho \pi^+ \pi^-) \rightarrow \pi^+ \pi^- \pi^+ \pi^-$$



- clear enhancement in  $m_{4\pi}$
- rel. BW + BG ( $O(m^5)$ )
- $\rightarrow m_{\rho'} = 1.57 \pm 0.02 \pm 0.01 \text{ GeV}$
- $\rightarrow \Gamma_{\rho'} = 0.18 \pm 0.06 \pm 0.07 \text{ GeV}$

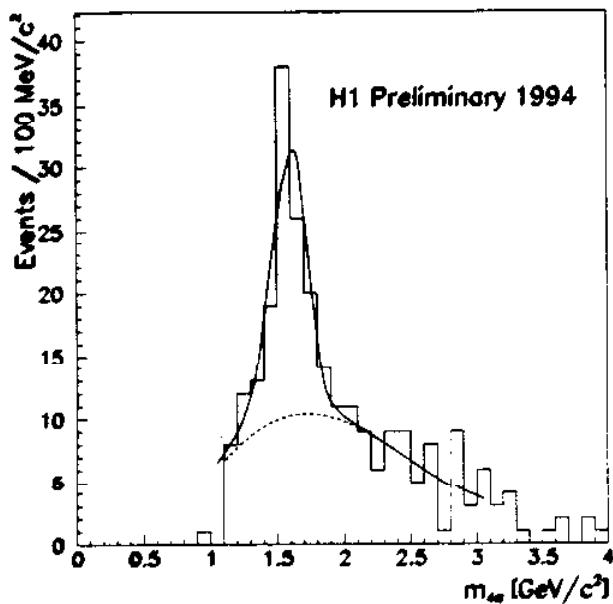


- $\rho$  signal in unlike sign distribution
- consistent with one  $\rho$  per event

$$4 < Q^2 < 50 \text{ GeV}^2 \quad 40 > W_{\gamma p} < 140 \text{ GeV}$$

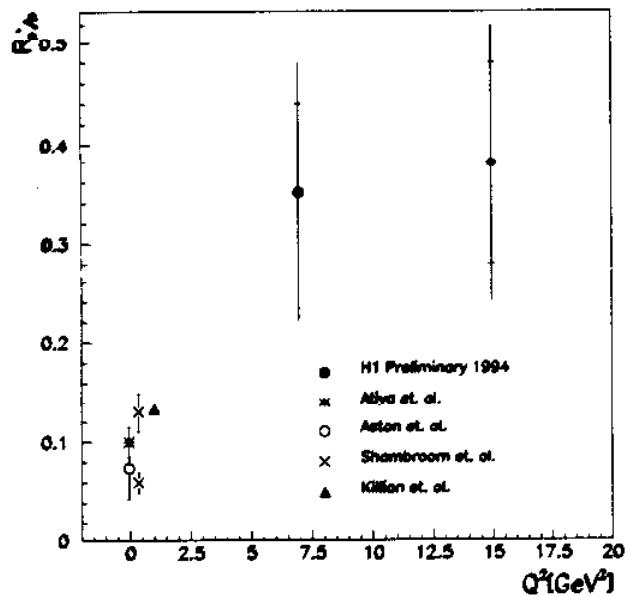
$\rho(1450)$  and/or  $\rho(1700)$  ?

Ansatz:  $|BW(\rho(1450)) + \alpha e^{i\beta} BW(\rho(1700))|^2$



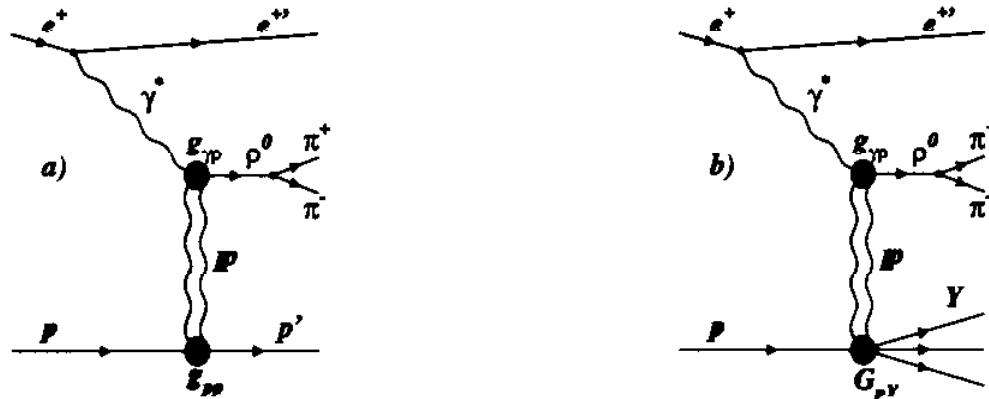
- $\rightarrow \alpha = 1.0 \pm 0.3$
- $\rightarrow \beta = 3.1 \pm 0.9$
- comparable  $\chi^2$  as for single resonance
- $\rightarrow$  no conclusion possible

ratio of  $\rho'$  to  $\rho$  production



- assuming single resonance
- $R_{\rho'}/\rho = 0.36 \pm 0.07 \pm 0.11$
- larger than at  $Q^2 = 0$   
( as expected by theory )

# test of a factorisation hypothesis



$$\frac{d^2 \sigma_{pdis}/dt dM_Y^2}{d\sigma_{el}/dt} = \left( \frac{G_{pY}(t, M_Y)}{g_{pp}(t)} \right)^2 = f(t, M_Y)$$

$$\frac{\sigma_{pdis} b_{pdis}}{\sigma_{el} b_{el}} = f_1(W) \quad b_{el} - b_{pdis} = f_2(W)$$

$f_1$  and  $f_2$  should be independent of the process

experiment	ISR,pp	H1 ( $\rho$ )	H1 ( $J/\Psi$ )
$E_{cms} / GeV$	53	60-180	30-150
$f_1$	$0.24 \pm 0.04$	$0.19 \pm 0.08$	$0.40 \pm 0.11$
$f_2 / GeV^{-2}$	$6.6 \pm 1.0$	$5.0 \pm 1.2$	$2.4 \pm 0.4$

⇒ factorisation seems to hold for  $\rho$   
perturbative effects for  $J/\Psi$  !?  
further confirmation needed !

## conclusion

- photoproduction

- slow rise of  $\sigma_{\gamma p}$  with  $W_{\gamma p}$  for the  $\rho^0$  well described by soft pomeron
- steeper rise for the  $J/\Psi$ , now seen from HERA data alone
  - best evidence yet for 'hard' pomeron
- $\rho^0$  and  $J/\Psi$  consistent with SCHC
- shrinkage observed for  $\rho^0$  ( indication for  $J/\Psi$  ?)

- high  $Q^2$

- $\Phi$  and  $J/\Psi$  show steep rise with  $W_{\gamma p}$  at high  $Q^2$  but still depending on low energy data
  - situation unclear for the  $\rho^0$ 
    - need  $W_{\gamma p}$  dependence from HERA data alone
- same for shrinkage
- $\rho^0$  and  $\Phi$  are predominantly longitudinally polarized
- clear signal of  $\rho'$  at high  $Q^2$ 
  - unclear whether it consists of one or two resonances
- $\rho^0$  electroproduction with proton dissociation shows much flatter t dependence

- what can we learn from  $\rho'$  and  $\Psi'$  production ?

## reconstruction of kinematic variables

$$\begin{aligned}
 Q^2 &= 4E_0^2 \frac{\sin \theta_V (1 + \cos \theta_e)}{\sin \theta_V + \sin \theta_e - \sin(\theta_V + \theta_e)} \\
 y &= \frac{\sin \theta_e (1 - \cos \theta_V)}{\sin \theta_V + \sin \theta_e - \sin(\theta_V + \theta_e)} \quad \left( Y = \frac{Q^2 - 0}{2CE - R_e} \right) \\
 E'_e &= [2E_0 - (E_V - |P_V| \cdot \cos \theta_V)] / (1 - \cos \theta_e) \\
 x &= \frac{Q^2}{sy} \quad W^2 = Q^2 \left( \frac{1}{x} - 1 \right) \\
 t &\simeq (\vec{p}_{tp})^2 = (\vec{p}_{te} + \vec{p}_{tV})^2 \\
 E - p_z &= (E_e + E_V) - (p_{ze} + p_{zV})
 \end{aligned}$$

## cross section definitions

$$\begin{aligned}
 \frac{d^2\sigma_{tot}(ep \rightarrow eVp)}{dy \ dQ^2} &= \Gamma_T \ \sigma_{tot}(\gamma^* p \rightarrow Vp) \\
 &= \Gamma_T \ \sigma_T(\gamma^* p \rightarrow Vp) (1 + \varepsilon \ R)
 \end{aligned}$$

$$\Gamma_T = \frac{\alpha_{em} (1 - y + y^2/2)}{\pi y Q^2}$$

$$\varepsilon = \frac{1 - y}{1 - y + y^2/2}. \quad ; \quad R = \sigma_L / \sigma_T$$