

NLO Analysis of SLAC E154 Data

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- **Motivation**
- **Formalism**
- **Experimental and Theoretical Uncertainties**
- **Results and Interpretation**
- **Summary**

Motivation

- **E154** $g_1^n(x)$ large and strongly x dependent at low $x \leq 0.1$ at $Q^2 = 5 \text{ GeV}^2$
What does it mean?
- **Evolution based on $g_1/F_1 = (\text{constant in } Q^2)$ assumption**
 - * **consistent within errors with experimental data**
 - * **incompatible with pQCD**
 - * **no associated systematic error**
- **Low x extrapolation for the sum rules**
 - * **Regge-like, $\sim \text{constant in } x$**
 - * **Power-like, $\sim x^{-0.9 \pm 0.2}$****model dependent.**
- **Reasonable kinematic range and precision of polarized data**

- Known NLO spin-dependent anomalous dimensions
R. Mertig and W. L. van Neerven (1996), W. Vogelsang (1996)
- Careful study of statistical and systematic uncertainties
- New E154 data on g_1^n
- Global analysis of x and Q^2 dependence
- Allows spin decomposition of structure functions onto valence and sea quarks and gluons
- Standard in unpolarized DIS data analyses

Factorization Theorem

Flavor Decomposition of g_1

$$g_1^p(x, Q^2) = \frac{1}{2} \sum_q^{N_f} e_q^2 [C_q \otimes \Delta q + \frac{1}{N_f} C_G \otimes \Delta G]$$

and for the neutron and deuteron

$$g_1^n = g_1^p(u \leftrightarrow d)$$

$$g_1^d = \frac{1}{2}(g_1^n + g_1^p)(1 - \frac{3}{2}\omega_D), \quad \omega_D = 0.05 \pm 0.01$$

Convolution \otimes is defined as

$$g_1(x) \equiv (C \otimes q)(x) = \int_x^1 \frac{dy}{y} C\left(\frac{x}{y}\right) q(y)$$

Mellin transform

$$f(n) = \int_0^1 dx x^{n-1} f(x)$$

$$g_1^p(n, Q^2) = \frac{1}{2} \sum_q^{N_f} e_q^2 [C_q \Delta q + \frac{1}{N_f} C_G \Delta G]$$

- * $C_{q,G}$ – **perturbative coefficients**
- * $\Delta q = \delta q + \delta \bar{q}$ – **non-perturbative input**
- * **Factorization scheme dependent in NLO**

- **Common Schemes**

- * **$\overline{\text{MS}}$ – gluons do not contribute to the integral of**
 $g_1, C_G^{(1)}(n=1)=0$
- * **Adler-Bardeen – include axial anomaly contribution,**
 $C_G^{(1)}(n=1)=-N_f$

$$\int_0^1 dx g_1^p(x) = \frac{1}{2} \left[\frac{4}{9} \Delta u + \frac{1}{9} \Delta d + \frac{1}{9} \Delta s \right] - \frac{\alpha_s}{6\pi} \Delta G$$

- **Consequence: ambiguity in definition of the "quark spin" contribution**

$$\Delta\Sigma(\text{AB}) = \Delta\Sigma(\overline{\text{MS}}) + \frac{N_f}{2\pi}\alpha_s \Delta G$$

- **Product $\alpha_s(Q^2)\Delta G(Q^2)$ is independent of Q^2 to $O(\alpha_s^2)$** J. Kodaira (1980)
 $\Rightarrow \Delta\Sigma$ is defined up to a constant
 \Rightarrow Asymptotically $\Delta G \sim 1/\alpha_s \rightarrow \infty$
- **Angular momentum sum rule revised:**
X. Ji (1996)

$$S_p = \frac{1}{2} = \left(\frac{1}{2}\Delta\Sigma + \langle L_z^q \rangle \right) + \left(\Delta G + \langle L_z^G \rangle \right)$$

Fitting Procedure

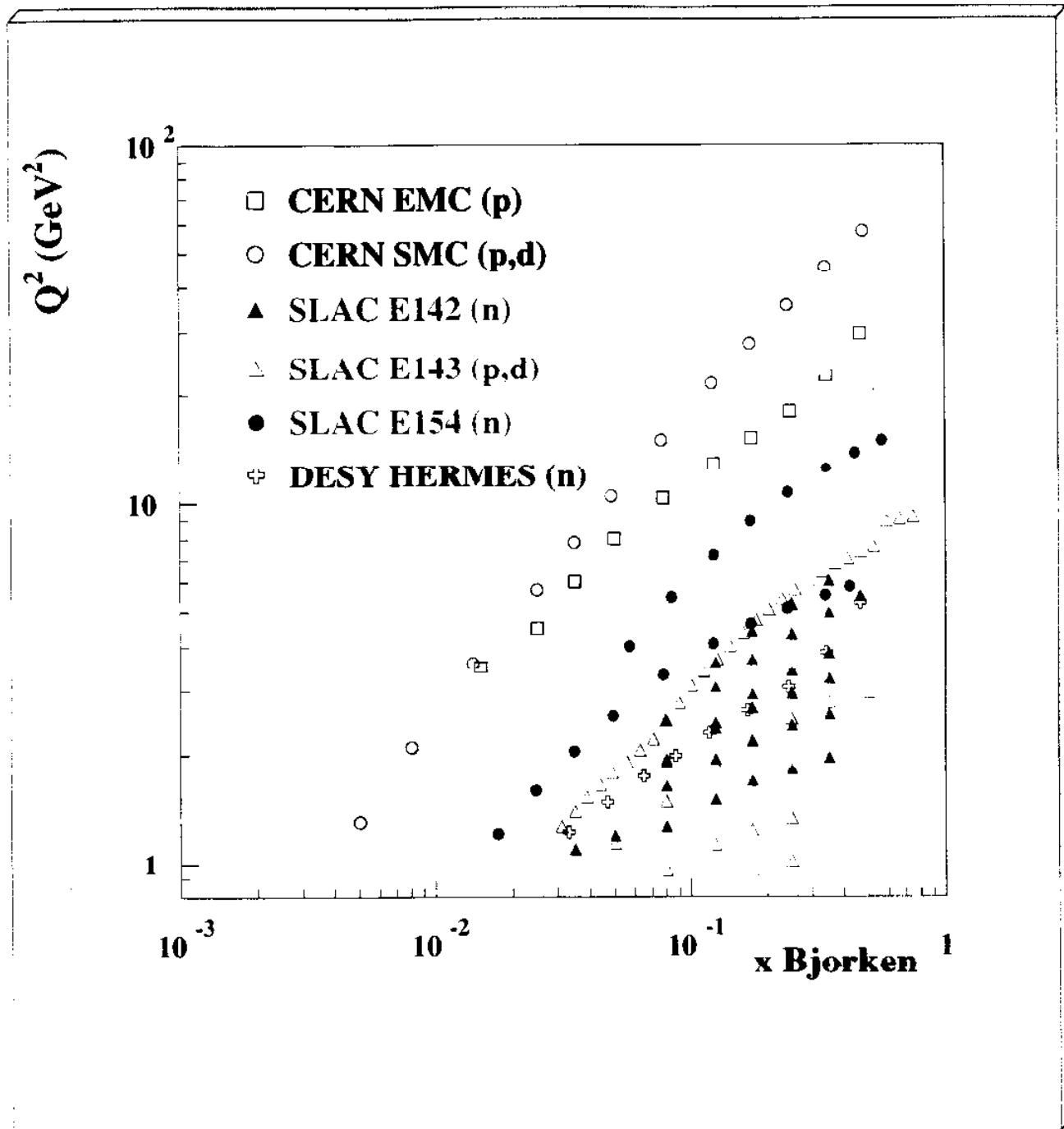
- Start with parametrized parton distributions $\Delta q(x, Q_0^2)$ at low Q_0^2
- Evolve up to experimental Q^2 using NLO DGLAP equations
- Compute g_1 at experimental (x, Q^2) and compare with data. Contributions to χ^2 weighted by statistical error
- Constrain 8 free parameters to give the best fit

Notice that we do not assume

$$\Delta q_3 = g_A = 1.2601 \pm 0.0025$$

$$\Delta q_8 = 0.579 \pm 0.025$$

Kinematic coverage



Initial Distributions

at $Q_0^2 = 0.34 \text{ GeV}^2$

- Assume the x dependence

$$\Delta f(x, Q_0^2) = A_f x^{\alpha_f} f(x, Q_0^2) (1 - x)^{\beta_f}$$

where

- * $\Delta f = \Delta u_V, \Delta d_V, \Delta \bar{Q}, \Delta G$
- * $f(x, Q_0^2)$ – unpolarized distributions of GRV95

- Positivity constraint

$$|\Delta f(x)| \leq f(x) \Rightarrow \alpha_f \geq 0, \beta_f \geq 0$$

- Helicity retention property

$$\frac{\Delta f(x)}{f(x)} \xrightarrow{x \rightarrow 1} 1 \Rightarrow \beta_f = 0$$

The fits are compatible with this assumption

- Assume sea isospin symmetry

$$\Delta \bar{u} = \Delta \bar{d}$$

- Then

$$g_1^{\text{sea}} \sim \frac{\Delta \bar{u} + \Delta \bar{d}}{2} + \frac{1}{5} \Delta \bar{s}$$

- * DIS does not probe light and strange sea separately
- * Sensitivity to the difference only via evolution, beyond present experimental precision

Thus parametrize

$$\Delta \bar{Q} \equiv \frac{\Delta \bar{u} + \Delta \bar{d}}{2} + \frac{1}{5} \Delta \bar{s}$$

- Assume $\Delta s(x, Q_0^2)$

$$\Delta s = \Delta \bar{s} = \lambda_s \frac{\Delta \bar{u} + \Delta \bar{d}}{2} = \frac{\lambda_s}{1 + \lambda_s/5} \Delta \bar{Q}$$

with parameter $\lambda_s = 0 - 1$

- Initial distributions

$$\Delta u_V(x, Q_0^2) = A_u x^{\alpha_u} u_V(x, Q_0^2)$$

$$\Delta d_V(x, Q_0^2) = A_d x^{\alpha_d} d_V(x, Q_0^2)$$

$$\Delta \bar{Q}(x, Q_0^2) = A_Q x^{\alpha_Q} \bar{Q}(x, Q_0^2)$$

$$\Delta G(x, Q_0^2) = A_G x^{\alpha_G} G(x, Q_0^2)$$

specified in terms of 8 free parameters.

Unpolarized distributions of GRV95

$$\begin{aligned} u_V(x, Q_0^2) &= 0.988 x^{-0.457} (1 + 1.58 \sqrt{x} + 2.58 x + 18.1 x^{3/2}) (1 - x)^{3.380} \\ d_V(x, Q_0^2) &= 0.182 x^{-0.684} (1 + 2.51 \sqrt{x} + 25.0 x + 11.4 x^{3/2}) (1 - x)^{4.113} \\ \bar{Q}(x, Q_0^2) &= 0.654 x^{-0.70} (1 + 2.65 x) (1 - x)^{8.33} \\ G(x, Q_0^2) &= 26.20 x^{0.9} (1 - x)^{4.0} \end{aligned}$$

NLO Evolution ($\overline{\text{MS}}$)

- DGLAP equations in n-moments space

$$Q^2 \frac{d}{dQ^2} \Delta q_{NS}^\eta = \frac{\alpha_s(Q^2)}{2\pi} \gamma_{NS}^\eta \cdot \Delta q_{NS}^\eta, \quad \eta = \pm 1,$$

$$Q^2 \frac{d}{dQ^2} \begin{pmatrix} \Delta\Sigma \\ \Delta G \end{pmatrix} = \frac{\alpha_s(Q^2)}{2\pi} \begin{pmatrix} \gamma_{qq} & \gamma_{qG} \\ \gamma_{Gq} & \gamma_{GG} \end{pmatrix} \cdot \begin{pmatrix} \Delta\Sigma \\ \Delta G \end{pmatrix}$$

with the NS distributions defined as

$$\left. \begin{array}{lcl} \Delta u_V & = & \delta u - \delta \bar{u} \\ \Delta d_V & = & \delta d - \delta \bar{d} \end{array} \right\} \eta=1$$

$$\left. \begin{array}{lcl} \Delta q_3 & = & \Delta u - \Delta d \\ \Delta q_8 & = & \Delta u - \Delta d - 2\Delta s \end{array} \right\} \eta=-1$$

and the singlet

$$\Delta\Sigma = \Delta u + \Delta d + \Delta s$$

- **Fixed flavor scheme, $N_f=3$**
- **Heavy flavors contribution included in running of two-loop α_s**

$$\alpha_s(m_q^2, f) = \alpha_s(m_q^2, f + 1)$$

- **Quark masses: $m_c=1.5$ GeV, $m_b=4.5$ GeV**
- **For consistency with the unpolarized fit**

$$\alpha_s(M_Z^2) = 0.109 \quad \Rightarrow \quad \alpha_s(5 \text{ GeV}^2) = 0.237$$

- **Thus**

$$\Lambda_{f=3,4,5} = 248, 200, 131 \text{ MeV}$$

Experimental and Theoretical Uncertainties

- **Statistical Errors**

- ✗ **Added in quadrature contributions from experimental points**

- **Systematic Errors**

- ✗ **Usually dominated by normalization errors (beam and target polarizations, dilution factor)**
- largely correlated from point to point
 - ✗ **Assume 100% correlation within experiment**
 - ✗ **Add contributions from any given experiment linearly**

- **Theoretical Uncertainties**

- * **α_s from fixed target DIS experiments,**

$$\alpha_s(M_Z^2) = 0.108 - 0.116,$$

has the biggest effect (renormalization and factorization scale dependence)

- * **Variation of current quark masses - small effect**

$$m_c = 1 - 2 \text{ GeV} \quad m_b = 4 - 5 \text{ GeV}$$

- * **Initial strange sea distribution - small effect**

$$\lambda_s = 0 - 1$$

- * **The x dependence of initial distributions and starting Q_0^2 scale: fit with MRS96 initial unpolarized distributions at $Q_0^2 = 1 \text{ GeV}^2$. Results consistent within statistical errors**
- * **Factorization scheme dependence: $\overline{\text{MS}}$ and AB schemes results consistent within errors**
- * **Possible higher twist effects neglected. Expected to drop as $1/W^2$ and $W^2 > 4 \text{ GeV}^2$ for all experimental data with majority $W^2 > 8 \text{ GeV}^2$**

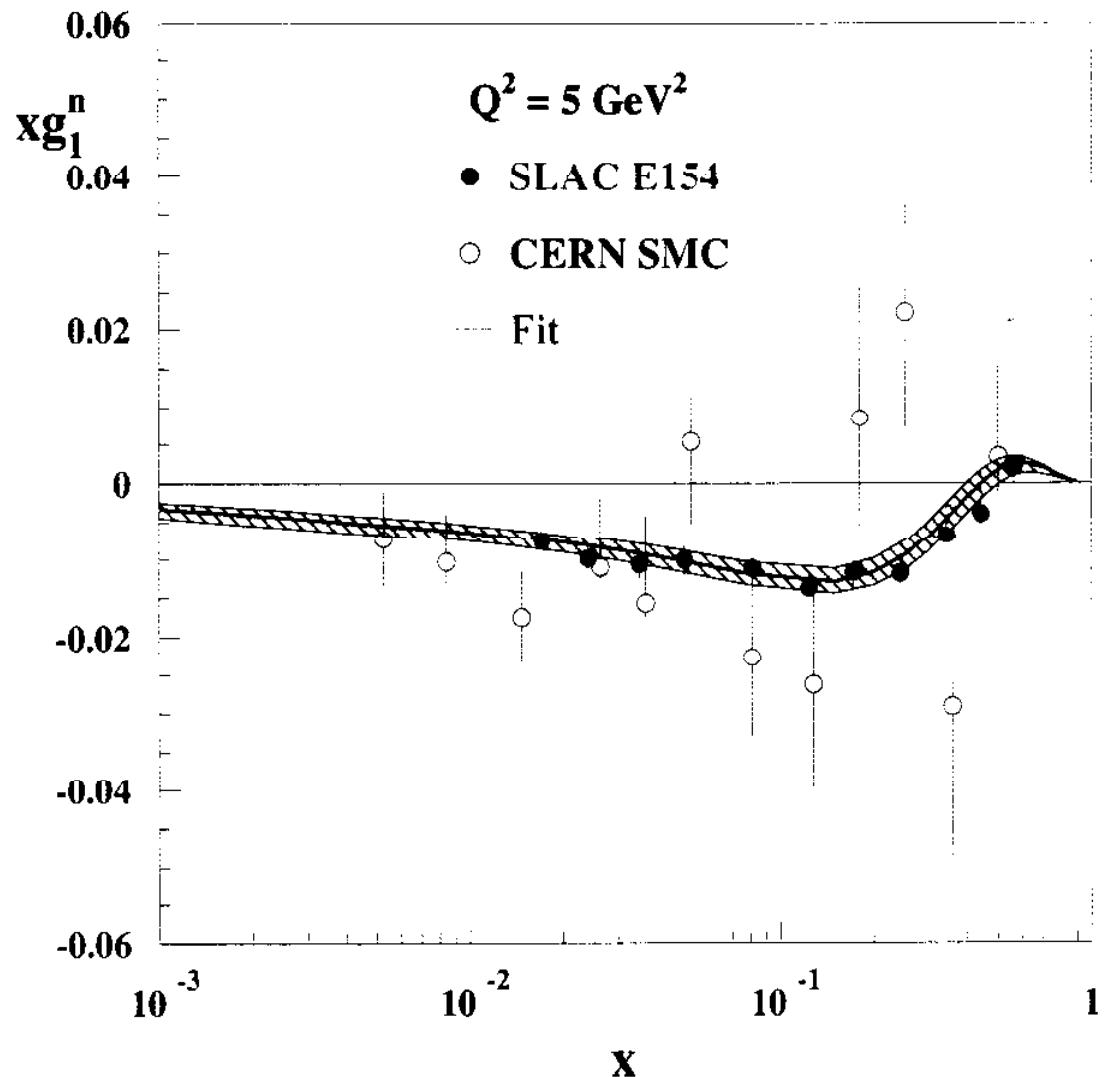
NLO Results

- Very good fits,

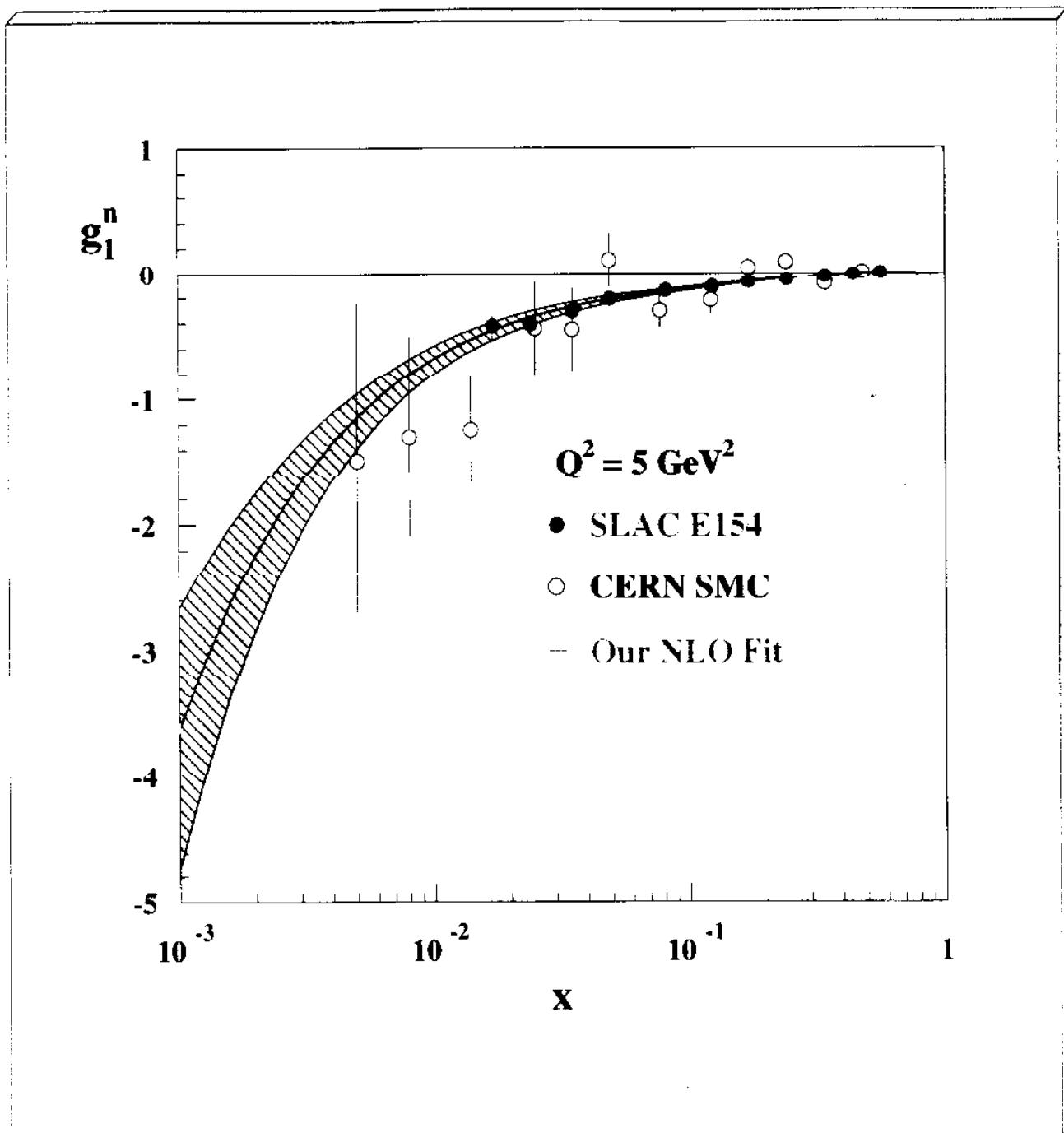
$$\chi^2 = 146 \text{ for } 168 \text{ points } (\overline{\text{MS}})$$
$$148 \quad (\text{AB})$$

- Low x behavior of valence distributions at the initial Q_0^2 scale consistent with Regge prediction
- First moments of valence distributions well determined
- First moments of sea quarks and gluons only qualitatively constrained
- Contribution of systematic errors comparable to the statistical
- Theoretical uncertainty quite large. Could be reduced by simultaneous analysis of unpolarized and polarized DIS data with α_s as one of the parameters

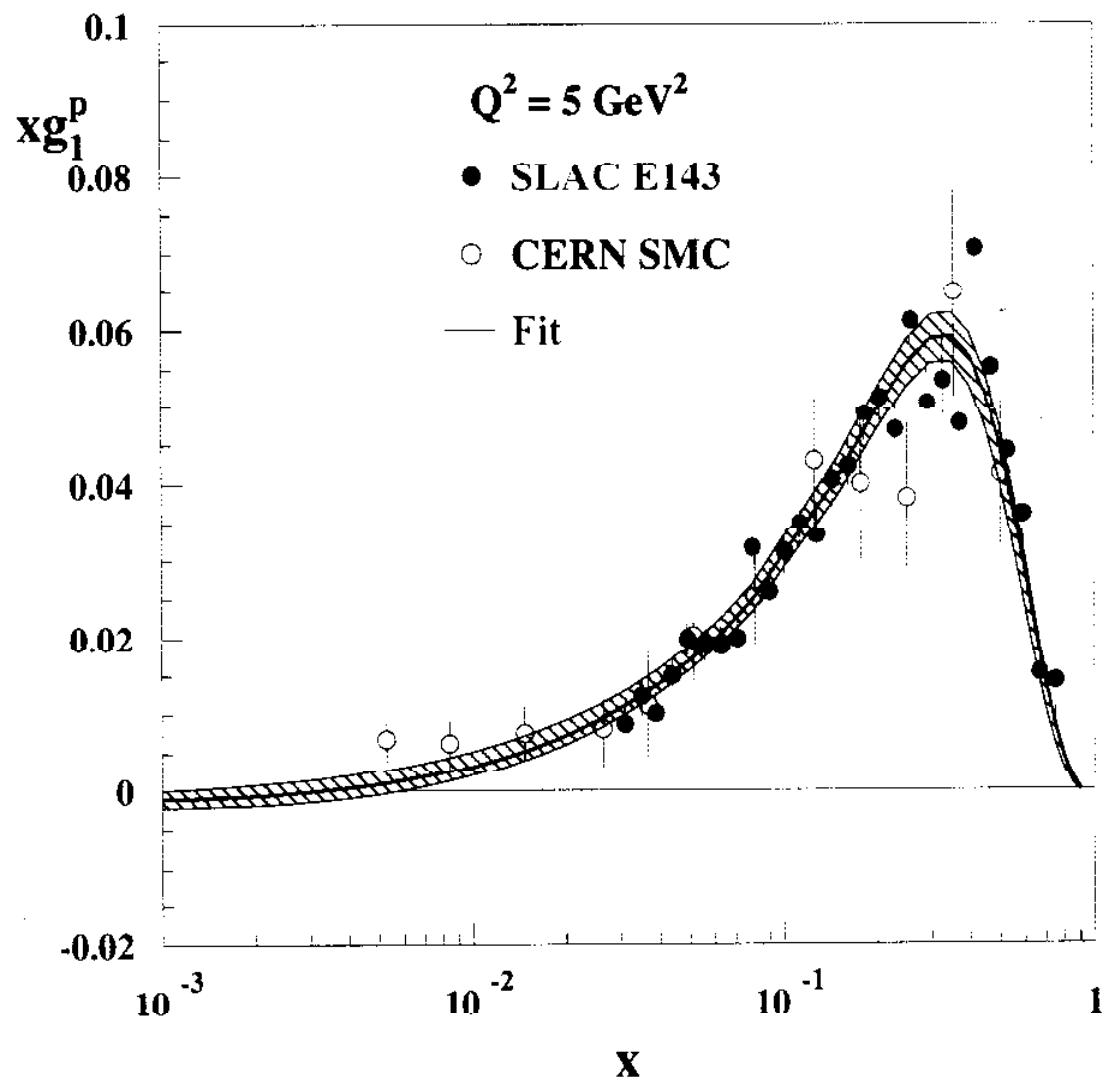
Neutron



Neutron



Proton



Fitted Values of Parameters ($\overline{\text{MS}}$)

	Value	Stat.	Syst.	Theory
A_u	0.99	+0.08 -0.08	+0.04 -0.05	+0.97 -0.11
A_d	-0.78	+0.14 -0.20	+0.05 -0.05	+0.05 -1.28
A_Q	-0.02	+0.03 -0.06	+0.01 -0.02	+0.01 -0.35
A_G	1.6	+1.1 -0.9	+0.6 -0.6	+0.2 -1.3
α_u	0.63	+0.06 -0.07	+0.04 -0.05	+0.36 -0.06
α_d	0.28	+0.15 -0.11	+0.05 -0.03	+0.75 -0.03
α_Q	0.04	+0.29 -0.03	+0.12 -0.03	+0.55 -0.01
α_G	0.8	+0.4 -0.5	+0.3 -0.3	+0.1 -0.6

	Value	Stat.	Syst.	Theory
Δu_V	0.69	+0.03 -0.02	+0.05 -0.04	+0.14 -0.01
Δd_V	-0.40	+0.03 -0.04	+0.03 -0.03	+0.07 -0.00
$\Delta \bar{Q}$	-0.02	+0.01 -0.02	+0.01 -0.01	+0.00 -0.03
ΔG	1.8	+0.6 -0.7	+0.4 -0.5	+0.1 -0.6
Δq_3	1.09	+0.03 -0.02	+0.05 -0.05	+0.06 -0.01
Δq_8	0.30	+0.06 -0.05	+0.05 -0.05	+0.23 -0.01
$\Delta \Sigma$	0.20	+0.05 -0.06	+0.04 -0.05	+0.01 -0.01
Γ_1^p	0.112	+0.006 -0.006	+0.008 -0.008	+0.009 -0.001
Γ_1^n	-0.056	+0.006 -0.007	+0.005 -0.006	+0.002 -0.001
Γ_1^d	0.026	+0.005 -0.006	+0.005 -0.006	+0.005 -0.001
Γ_1^{p-n}	0.168	+0.005 -0.004	+0.008 -0.007	+0.007 -0.001

- Notice that $\Delta q_3 = 1.09^{+0.09}_{-0.06}$ is 2σ lower than $g_A = 1.2601$. Higher order corrections could be $\sim 5\%$ at 5 GeV^2 bringing Δq_3 to a better agreement
- NLO evolution of data points

$$g_1^{exp}(x_i, \langle Q^2 \rangle) = g_1^{exp}(x_i, Q_i^2) - \Delta g_1^{fit}(x_i, Q_i^2, \langle Q^2 \rangle)$$

where

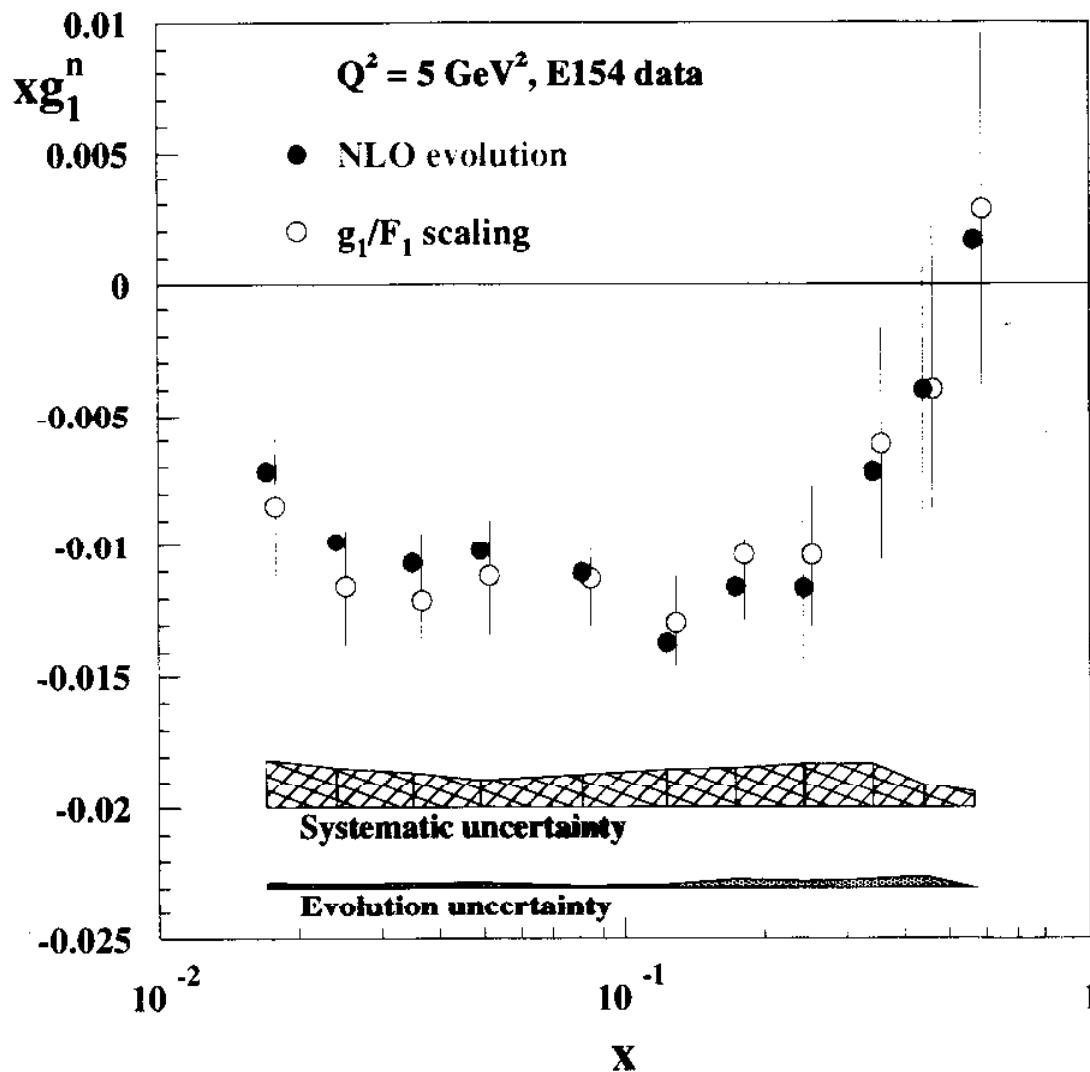
$$\Delta g_1^{fit}(x_i, Q_i^2, \langle Q^2 \rangle) = g_1^{fit}(x_i, Q_i^2) - g_1^{fit}(x_i, \langle Q^2 \rangle)$$

and

$$\sigma^2(g_1^{exp}(x_i, Q_i^2)) = \sigma^2(g_1^{exp})_{stat} + \sigma^2(g_1^{exp})_{syst} + \sigma^2(g_1)_{evol}$$

- Difference between NLO and $g_1/F_1 = \text{const}$ evolutions is on the order of present day experimental precision

Q^2 evolution



- The integrals over measured range less sensitive to the type of evolution if computed at the average $\langle Q^2 \rangle$ of the experiment. For E154 at 5 GeV²

$$\int_{0.014}^{0.7} dx g_1^n(x) = -0.035 \pm 0.003 \pm 0.005 \pm 0.001 \text{ NLO}$$

$$= -0.036 \pm 0.004 \pm 0.005 \quad g_1^n / F_1^n$$

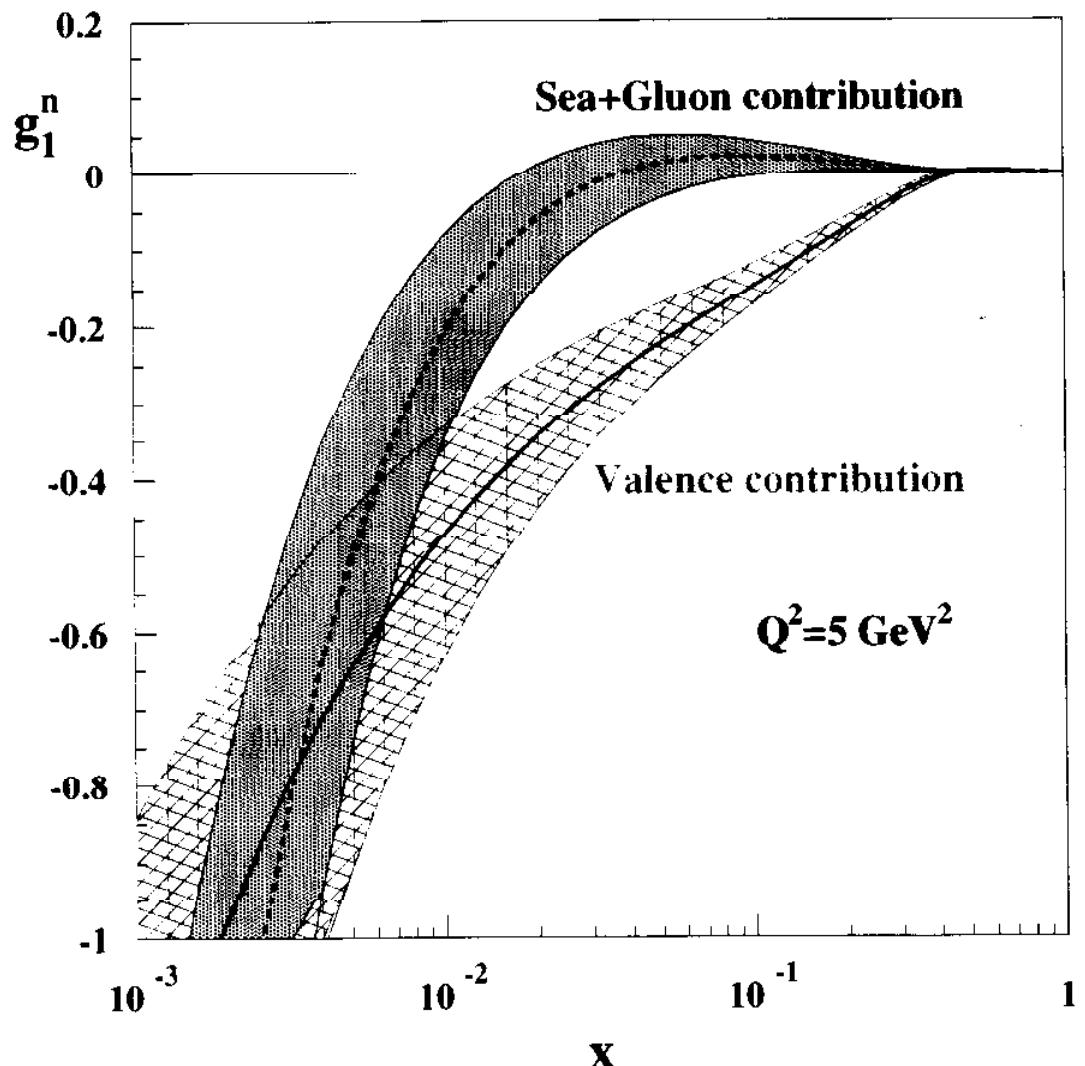
- Strong x behavior of

$$g_1^n \sim x^{-0.8}$$

for $x \leq 0.1$ at $Q^2 = 5$ GeV² largely due to the sea and gluon contributions

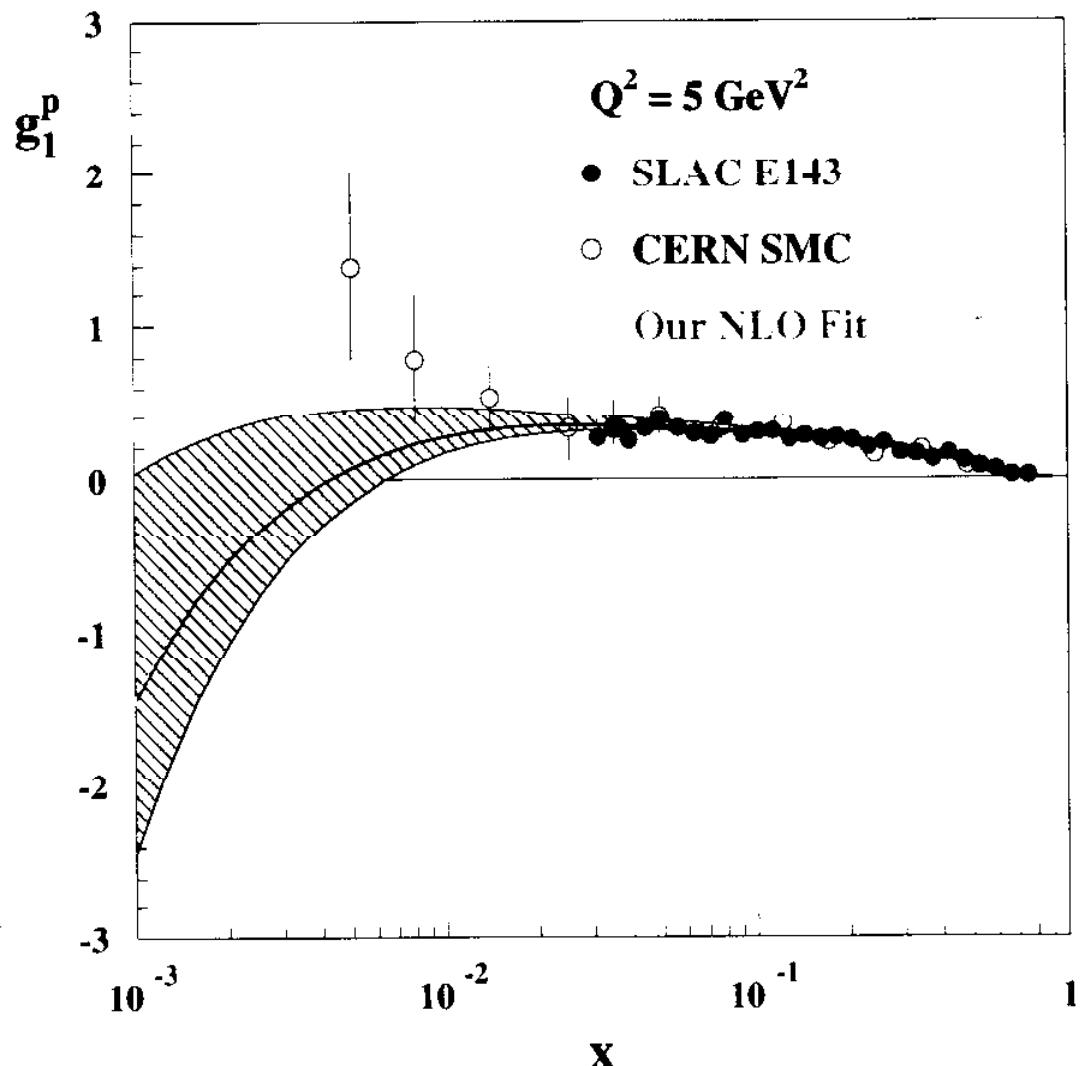
- Negative value of g_1^p at lower x and higher Q^2 would provide direct evidence of polarized sea is on the order of present day experimental precision

Sea and Valence



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Proton



- First moment of $g_1^n(x)$
 - * Using E154(n) + $\overline{\text{MS}}$ fit for low and high x extrapolations

$$\Gamma_1^n = -0.058 \pm 0.004 \pm 0.007 \pm 0.007$$

where the errors are $\pm \text{stat.} \pm \text{syst.} \pm \text{extrap.}$

- Bjorken sum rule $\Gamma_1^{p-n} = \int_0^1 dx (g_1^p(x) - g_1^n(x))$
 - * Using E143(p) + E154(n) + $\overline{\text{MS}}$ extrapolations

$$\Gamma_1^{p-n} = 0.171 \pm 0.004 \pm 0.010 \pm 0.006$$

- * Prediction up to $O(\alpha_s^3)$,

$$\Gamma_1^{p-n}(\text{theory}) = 0.186$$

Summary

- The first moments of polarized valence distributions well established
- Gluon and sea quarks distributions only qualitatively constrained
- Q^2 dependence of g_1/F_1 is sizable in comparison to present experimental uncertainties
- Low x behavior of the non-singlet distributions understood
 - * "Regge-like" at low Q_0^2
 - * Divergent at $Q^2 > 1 \text{ GeV}^2$ due to evolution
 - * Powers well constrained
- Bjorken sum rule confirmed within one σ