

ADVANCED TECHNIQUES

FOR

MULTIPARTON LOOP CALCULATION

-- A mini Review

Z. Bern

L. Dixon

D. Dunbar

D. Kosower

QCD *is* the theory of the strong interactions.

The era of *testing QCD* is over.

We are in the era of using QCD, and in particular jet physics,

- To search for new physics,
- To measure non-perturbative hadronic quantities.

For these purposes, we need precise calculations in perturbative QCD; next-to-leading order ones are a bare minimum.

# IN MATRIX ELEMENT CALCULATIONS

- 1

Tree level! analytically up to 7 partons or 6+ vector boson  
 Berends, Giele, Kuifje, Tausk  
 Leading-order predictions for  $p\bar{p} \rightarrow \leq 5$  jets or  $V + \leq$   
 Kuifje BG+

all-n for certain helicities & external particles  
 Parke, Taylor, Mangano; Berends & Giele; DAK.

numerically recurrence relations for all-n with  $\leq 4$  external ge  
 Berends, Giele.

One loop analytically 4 partons (unpolarized) Ellis & Sexton  
 $V + 3$  partons (unpolarized) Ellis, Ross, Terning; Ali et al.  
 Gonsalves

$V + 3$  partons (polarized) Giele & Glover

4 partons (polarized) Bern & DAK; Kunszt, Signer, Trica

5 partons Bern, Dixon, DAK; Kunszt, Signer, Trica

Next-to-leading order predictions for  $p\bar{p} \rightarrow 2$  jets

Aversa, Chiappetta, Greco, Galletti; Ellis, Kunszt,  
 Giele, Glover, DAK

for  $e^+e^- \rightarrow \leq 3$  jets Kunszt & Nason; Giele & Glover;  
 Catani & Seymour

for  $e p \rightarrow 2$  jets + remnant Minkov & Zeppenfeld;  
 Catani & Seymour

$\pi^+ p \rightarrow \pi^+ p$ : Minkov, Zeppenfeld; Giele & Kilgore; Kunszt, Frixtum, Signer, Trica; et al.

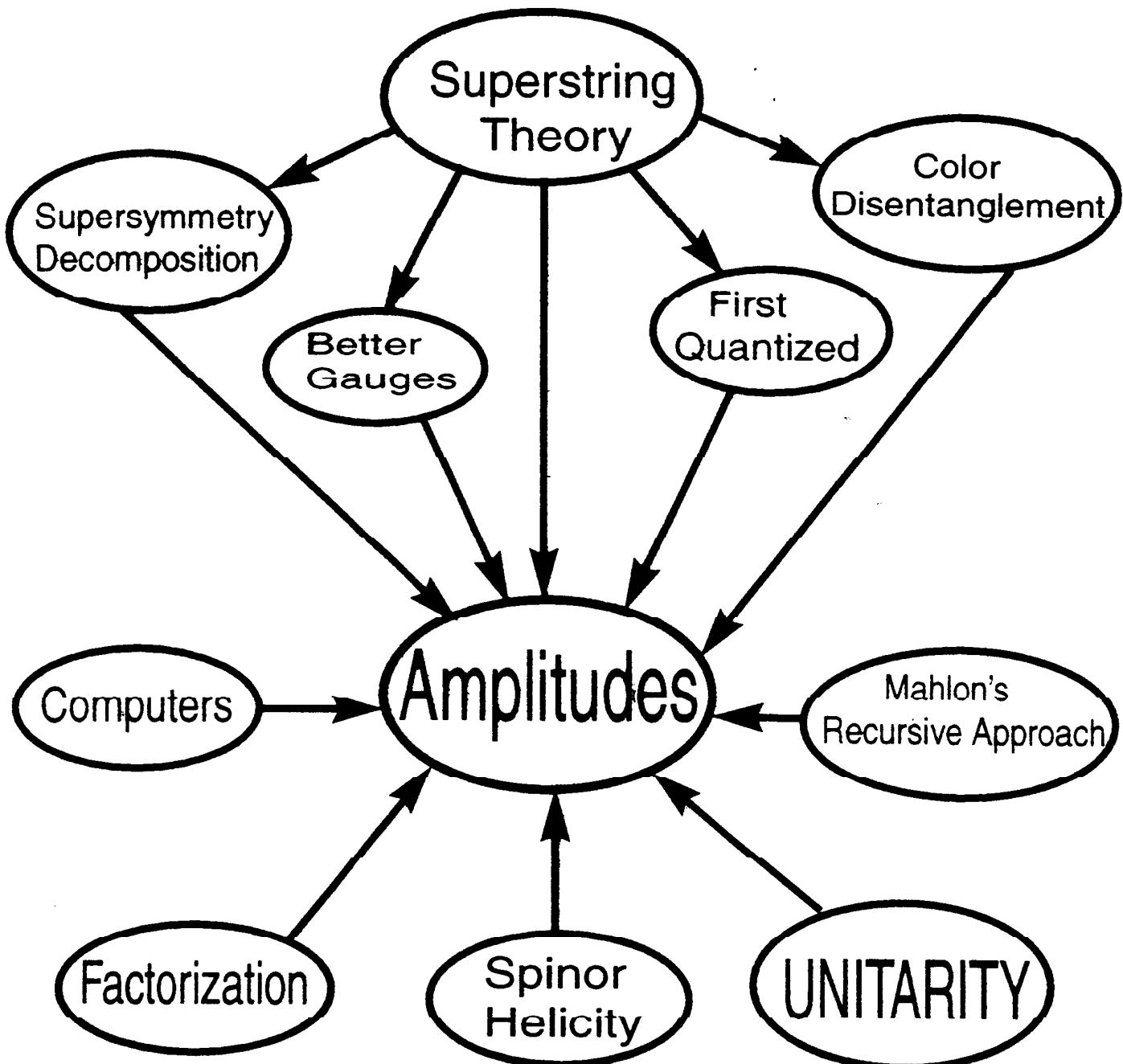
$\pi^+ \pi^- \pi^+ \pi^-$  Bern, Dixon, DAK, Weinzierl;  
 Glover & Miller

all-n formulae for certain helicities  
 Bern, Dixon, Duhr, DAK; Mahlon

Two loops

2 partons + V  
 van Neerven et al.

## Non-Traditional Approaches



Power of methods magnified when ideas are combined.

# SPINOR PRODUCTS

Kleiss, Stirling; Xu, Zhang, Chang.

Ultimately, we need the square of the amplitude. If we square it analytically, using  $\sum \epsilon_\mu^{*(i)} \epsilon_\nu^{(i)} = g_{\mu\nu}$ , we get a huge number of terms (at least in intermediate stages), and a less than optimally compact final answer. Ideally, we would square numerically; but how can we evaluate the amplitude numerically?

Introduce spinor products  $\bar{u}(k_1) u(k_2) = \langle 12 \rangle = \langle 1^- | 2^+ \rangle$   
 $\downarrow$   
 + massless  $k_1^z = 0$

These can be evaluated numerically,

$$\langle 12 \rangle = \frac{\overline{k_{1+}}}{\sqrt{k_{2+}}} (k_{2x} + i k_{2y}) - \frac{\overline{k_{2+}}}{\sqrt{k_{1+}}} (k_{1x} + i k_{1y})$$

$$[12] = -\text{sign}(k_1^0 k_2^0) \langle 12 \rangle^*$$

$$k_{1+} = k_1^0 + k_1^z$$

are antisymmetric & the  $\sqrt{-}$  of a Lorentz product:

$$\langle 12 \rangle [21] = 2 k_1 \cdot k_2$$

Define representation for polarization vectors,

$$\epsilon_\mu^{(+) \dagger} = \frac{\langle g^- | \gamma_\mu | k^- \rangle}{\sqrt{2} \langle g | k \rangle}, \quad \epsilon_\mu^{(-) \dagger} = \frac{\langle g^+ | \gamma_\mu | k^+ \rangle}{\sqrt{2} \langle g | k \rangle}$$

where  $g$  is an arbitrary null reference momentum.  
 Changing  $g$  shifts  $\epsilon$  by  $\lambda k$  — i.e. it is a gauge transform



$$A_5(+-+-+)=0$$

$$A_5(-++++)=0$$

$$A_5(--+-+)=\frac{i \langle 12 \rangle^4}{\langle 12 \rangle \langle 23 \rangle \langle 34 \rangle \langle 45 \rangle \langle 51 \rangle}$$

$$A_5(-+-+)=\frac{i \langle 13 \rangle^4}{\langle 12 \rangle \langle 23 \rangle \langle 34 \rangle \langle 45 \rangle \langle 51 \rangle}$$

# THE TREE-LEVEL STORY

- Color decomposition

$$A_{n+1}^{\text{tree}} = g^{n-2} \sum_{\sigma \in S_{n-2}} (T^{a_0 \dots T^{a_{n-2}}})^{\bar{i}_1}_{\bar{i}_2} A_{n+1}^{\text{tree}}(1_{\bar{i}_1}, 2_{\bar{i}_2}; \sigma(3 \dots n); l_i)$$

- Helicity amplitudes, using spinor products  $\langle ij \rangle, [ij]$

$$\langle ij \rangle [ij] = 2 k_i \cdot k_j$$

Partial or color-ordered sub-amplitudes  $A_n^{\text{tree}}$

- are gauge invariant
- satisfy Faddeev identities
- can split up into scattering functions
- their decomposition (and combination) is very useful for building universal functions summarizing [over soft & collinear phase space]

Collinear limits

$$A_n^{\text{tree}}(1, \dots, n-2, a^{\lambda_1}, b^{\lambda_2}) \xrightarrow{a \parallel b} \sum_{\lambda_1 + \lambda_2 = \lambda} C_{\lambda_1 \lambda_2}^{\text{tree}}(a^\lambda, b^\lambda) A_{n-1}^{\text{tree}}(1, \dots, n-2, z)$$

$$k_a = z k_\Sigma, \quad k_b = (1-z) k_\Sigma, \quad C \sim \sqrt{\text{AP fons}}$$

e.g. gluons  $C_+^{\text{tree}}(a^+, b^+) = 1 = C_{-+}^{\text{tree}}(a^-, b^+),$

$$C_{+-}^{\text{tree}}(a^+, b^-) = \frac{1}{\sqrt{2}(1+\delta)}, \quad -C_{+-}^{\text{tree}}(a^-, b^-) = \frac{1}{\sqrt{2}(1-\delta)}$$

$$C_{--}^{\text{tree}}(a^-, b^-) = \frac{1}{\sqrt{2}(1-\delta)}, \quad -C_{--}^{\text{tree}}(a^+, b^+) = \frac{1}{\sqrt{2}(1+\delta)}$$



# STRING-BASED RULES FOR PERT. QCD

All- $n$  master formula transformed according to  $\phi^3$  diagrams via two rules ('pinch' and 'loop expansion') into a Feynman parameter polynomial. ( $\oplus$  loop integrals  $\rightarrow$  amplitude)

String theory important to derivation : path is straightforward  
Use of rules does not require knowledge of string theory  
Can be understood in ordinary field theory language  
(Bern + Dunbar ;  $\rightarrow$  Lessons for variety of gauge-theory calculations)

- Minimizes number of terms
- Meshes well with the spinor-helicity method
- Cancellations squeezed out at an early stage
- Systematic structure

Reverses connections between contributions of internal states with different spins

Allows required combination of tensor integrals to be performed in one fell swoop

Suggests formula for subleading-color amplitudes in terms of leading one.

Superseded by unitarity-based rules for nitty-gritty of calculations, but remains useful for hints & as organizing principle.

- Color decomposition

$$A_{n+1}^{\text{tree}} = g^{n-2} \sum_{\sigma \in S_{n-2}} (T^{a_{\sigma(3)}} \dots T^{a_{\sigma(n)}})_{i_2}^{i_1} A_n^{\text{tree}}(1_g, 2_g; \sigma(3..n); l_1 l_2)$$

$$A_{n+1}^{\text{1-loop}} = g^n \sum_{j=1}^{n-1} \sum_{\sigma \in S_{n-2}/S_{n-j}} G_{n;j}^{(\bar{q}q)} (\sigma(3..n)) A_{n;j}^{\text{1-loop}}(1_g, 2_g; \sigma(3..n); l_1 l_2)$$

$$G_{n;1}^{(\bar{q}q)}(3..n) = N_c (T^{a_3} \dots T^{a_n})_{i_2}^{i_1} \quad \text{leading color}$$

$$G_{n;j+1}^{(\bar{q}q)}(3..n) = \text{Tr} (T^{a_3} \dots T^{a_{j+1}}) (T^{a_{j+2}} \dots T^{a_n})_{i_2}^{i_1} \quad \text{subleading color}$$

- Spinor products  $\langle ij \rangle, [ij] \quad \langle j i \rangle [ij] = 2k_i \cdot k_j$   
 $\bar{u}(k_i) u(k_j)$

- Unitarity

Cut-containing parts (the bulk of the amplitude) can be determined uniquely by sewing on-shell tree amplitude, rendering the calculation vastly simpler than with traditional methods. Rational pieces can be calculated using cuts to  $\Theta(z)$ .

- Factorization in collinear limits  
Determine rational pieces.

- String-based ideas as organizational principle  
Leading-color  $\Rightarrow$  subleading-color; use of supersymmetric decomposition & supersymmetry identities

# COLLINEAR LIMITS AT ONE LOOP

$$A_{n;1}^{[J]}(1, \dots, n-2, a^{\lambda_a} b^{\lambda_b}) \xrightarrow{\text{coll}} \quad \text{from } 3\sigma_n \neq \omega_n$$

$$\sum_{\lambda_\Sigma = \pm} C_{-\lambda_\Sigma}^{\text{tree}}(a^{\lambda_a} b^{\lambda_b}) A_{n-1;1}^{[J]}(1, \dots, n-2, \Sigma^{\lambda_\Sigma}) + C_{-\lambda_\Sigma}^{[J]}(a^{\lambda_a} b^{\lambda_b}) A_{n-1}^{\text{tree}}(1, \dots, n-2, \Sigma^{\lambda_\Sigma})$$

i.e. tree  $\times$  loop + loop  $\times$  tree

$$C_+^{[1]}(++) = - \frac{[ab]}{\langle ab \rangle^2} \sqrt{\frac{z(1-z)}{3}}$$

$$C_+^{[1]}(++) = - C_+^{[1/2]}(++) = C_+^{[0]}(++)$$

others:  $C_\lambda^{[J]}(\lambda_a, \lambda_b) = c_r r_\lambda^{[J]}(\lambda_a, \lambda_b) C_\lambda^{\text{tree}}(\lambda_a, \lambda_b)$

$$c_r = \frac{\Gamma^2(1-\epsilon) \Gamma(4+\epsilon)}{16\pi^2 \Gamma(1-2\epsilon)}$$

unrenormalized

$$r_-^{[1]}(++) = - \frac{1}{\epsilon^2} \left[ \frac{\mu^2}{z(1-z)(-s_{ab})} \right]^\epsilon + 2 \ln z \ln(1-z) + \frac{z(1-z)}{3} - \frac{\pi^2}{6}$$

$$r_-^{[0]}(++) = \frac{z(1-z)}{3} = - r_-^{[1/2]}(++)$$

$$r_+^{[1]}(\pm \mp) = - \frac{1}{\epsilon^2} \left[ \frac{\mu^2}{z(1-z)(-s_{ab})} \right]^\epsilon + 2 \ln z \ln(1-z) - \frac{\pi^2}{6}$$

$$r_+^{[0]}(\pm \mp) = 0 = r_+^{[1/2]}(\pm \mp)$$

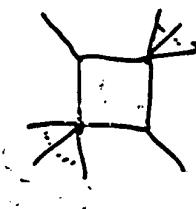
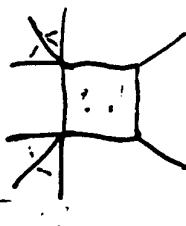
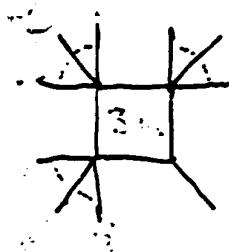
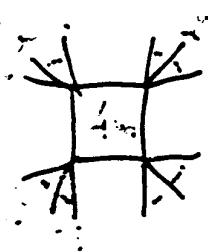
1.2.1.  $r^{\text{SUSY}}$  is independent of helicities

# UNIQUE RECONSTRUCTION OF LOOPS

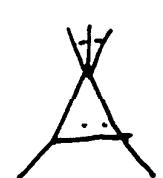
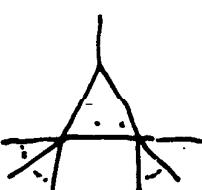
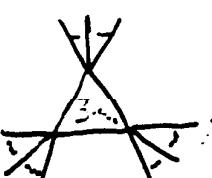
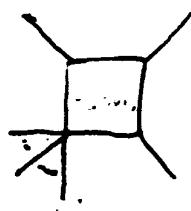
- For  $n > 4$ , can reduce any tensor integral to a sum of scalar boxes, tensor triangle, and tensor bubble integrals using Passarino-Veltman & van Neerven-Vermaseren reductions or equivalent techniques.



- All  $D=4-2\epsilon$   $n$ -tensor  $n$ -point integrals ... are expressed in terms of



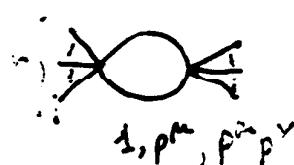
$$\mathcal{F}_n = \{$$



$$1, p^{\mu}, p^{\mu} p^{\nu}, p^{\mu} p^{\nu} p^{\lambda}$$

$$1, p^{\mu}, p^{\bar{\mu}} p^{\nu}, p^{\mu} p^{\nu} p^{\lambda}$$

$$1, p^{\mu}, p^{\bar{\mu}}$$



$$1, p^{\mu}, p^{\bar{\mu}} p^{\nu}$$

$$A_n = \sum_{\{I_1, I_2, \dots, I_n\}} \quad :I:$$



rational function of invariants  
& spinor products

If, for every loop integral

$$\int \frac{k^2}{\dots} \frac{P(k^\mu)}{k^2 (k-k_1)^2 (k-k_2)^2 \dots (k-k_n)^2}$$

$$\begin{aligned} \text{the degree of } P &\leq n-2 & n > 2 \\ &\leq 1 & n = 2 \end{aligned}$$

then  $A_n$  is determined uniquely by its cuts in  $\mathbb{R}^4$ .

Explanation:

$$\sum_{i|I_i \in \mathcal{F}_n} c_i I_i \Big|_{\text{cuts}} = \sum_{i|I_i \in \mathcal{F}_n} c'_i I_i \Big|_{\text{cuts}}$$

Then,

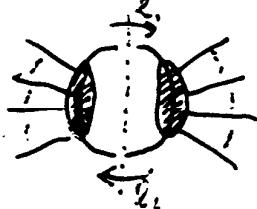
$$\sum_{i|I_i \in \mathcal{F}_n} (c_i - c'_i) I_i \quad (= \text{rational}) = 0$$

Proof via consideration of large-invariant regions & demonstration that cuts don't overlap.

## METHOD OF CALCULATION

Basic idea: sew tree amplitudes not tree diagrams;  
extract coefficients  $c$ :

- Look at a given channel carrying invariant  $(k_i + \dots + k_{i+n-1})^2$
- Form the cut in that channel, summing over all intermediate states



$$\int d^D k_1 \delta(k_1 \cdot k_2) A^{\text{tree}}(-k_1, -k_2) A^{\text{tree}}(-k_2, -k_3)$$

- "Promote" it to a loop integral by writing as an imaginary part,

$$\left( \int \frac{d^D k_1}{(2\pi)^D} A^{\text{tree}}(-k_1, -k_2) \frac{1}{k_1} \cdot A^{\text{tree}}(-k_2, -k_3) \frac{1}{k_2} \right) \text{-cut}$$

- Perform reductions & extract coefficient of function containing cut in this channel
- Re-assemble answer by summing over channels.

(In  $N=4$  case, linear independence gives some cross-channels "for free")

In collaboration with

Zvi Bern (UCLA)

Lance Dixon (SLAC)

Dave Dunbar (Swansea)

recently,

Stefan Weinzierl (student)  $Z \rightarrow g\bar{g}g'\bar{g}$

RESULTS (ONE LOOP)

Separate helicity amplitudes for  $0 \rightarrow gggg$

$0 \rightarrow ggggg$

$0 \rightarrow g\bar{g}ggg$

all-n  $0 \rightarrow g^+g^+\dots g^+$

all-n  $0 \rightarrow g^-g^+\dots g^+g^-g^+g^-\dots g^+$  (MHV) in N=4 SUSY

all-n  $0 \rightarrow g^-g^+\dots g^+g^-g^+g^-\dots g^+$  (MHV) in N=1 SUSY

$0 \rightarrow gggggg$  in N=4 SUSY

In final stages

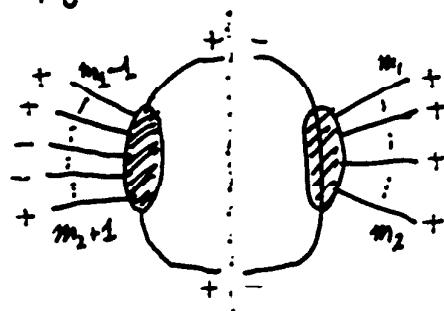
$0 \rightarrow Vg\bar{g}gg$

$0 \rightarrow Vg\bar{g}g'\bar{g}'$  (done & accepted in NPB)

# SAMPLE CALCULATION: $N=4$ MHV

Two configurations

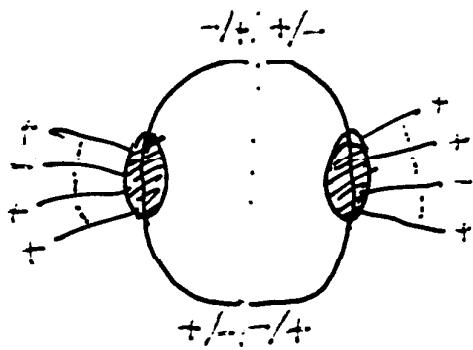
(a)



gluons only;

MHV tree  $\Rightarrow$  simple all-n  
formulae available

(b)



all intermediate states;  
simple formulae for  
SUSY Ward identities

After summing over  $N=4$   
multiplet, integrand is identical to  
case (a).

$$A^{tree} A^{\text{loop}} A^{\text{loop}} = -i A^{\text{loop}} \int d\tau \cos \frac{\sum_{j=1}^4 m_j \sin \theta_j \dot{\theta}_j}{\sum_{j=1}^4 k_j \sin \theta_j \dot{\theta}_j}$$

$$= A^{\text{tree}} \sum_{1m, 2mc} \text{Hatched denominator} \equiv A^{\text{tree}} V_n^g$$

Define for each box integral  $I_{4;r-i}^{nm}$ :

$$F_{4;r-i}^{nm} = -2\text{denominator}(I) \cdot I$$

Then

$$V_{2m+1}^3 = (\mu^2)^{\frac{1}{2}} \left[ \sum_{r=1}^{m+1} \sum_{i=1}^n F_{m+r,i}^{3nm} + \sum_{i=1}^n F_{m+1,i}^{3nm} \right]$$

$$V_{2m}^3 = (\mu^2)^{\frac{1}{2}} \left[ \sum_{r=1}^{m+1} \sum_{i=1}^n F_{m+r,i}^{3nm} + \sum_{i=1}^n F_{m+1,i}^{3nm} + \sum_{i=1}^n F_{m+2,i}^{3nm} \right]$$

Define

$$t_i^{[r]} = (k_i + \cdots + k_{i+r-1})^2$$

for all  $n \geq 5$ :

$$V_n^g = \sum_{i=1}^n -\frac{1}{\epsilon^2} \left( \frac{\mu^2}{-t_i^{[2]}} \right)^\epsilon - \sum_{r=2}^{\lfloor n/2 \rfloor - 1} \sum_{i=1}^n \ln \left( \frac{-t_i^{[r]}}{-t_i^{[r+1]}} \right) \ln \left( \frac{-t_{i+1}^{[r]}}{-t_i^{[r+1]}} \right) + D_n + L_n + \frac{n\pi^2}{6},$$

where for  $n = 2m + 1$ ,

$$D_{2m+1} = - \sum_{r=2}^{m-1} \left( \sum_{i=1}^n \text{Li}_2 \left[ 1 - \frac{t_i^{[r]} t_{i-1}^{[r+2]}}{t_i^{[r+1]} t_{i-1}^{[r+1]}} \right] \right),$$

$$L_{2m+1} = -\frac{1}{2} \sum_{i=1}^n \ln \left( \frac{-t_i^{[m]}}{-t_{i+m+1}^{[m]}} \right) \ln \left( \frac{-t_{i+1}^{[m]}}{-t_{i+m}^{[m]}} \right),$$

whereas for  $n = 2m$ ,

$$D_{2m} = - \sum_{r=2}^{m-2} \left( \sum_{i=1}^n \text{Li}_2 \left[ 1 - \frac{t_i^{[r]} t_{i-1}^{[r+2]}}{t_i^{[r+1]} t_{i-1}^{[r+1]}} \right] \right) - \sum_{i=1}^{n/2} \text{Li}_2 \left[ 1 - \frac{t_i^{[m-1]} t_{i-1}^{[m+1]}}{t_i^{[m]} t_{i-1}^{[m]}} \right],$$

$$L_{2m} = -\frac{1}{4} \sum_{i=1}^n \ln \left( \frac{-t_i^{[m]}}{-t_{i+m+1}^{[m]}} \right) \ln \left( \frac{-t_{i+1}^{[m]}}{-t_{i+m}^{[m]}} \right).$$

# WHAT ABOUT PURE RATIONAL PIECES?

Two approaches:

- a) Deduce them from collinear limits,

$$\lim_{\epsilon \rightarrow 0} A_n|_{\text{cuts}} + A_n|_{\text{rational}} \xrightarrow{\text{deduced}} \text{Split} \cdot A_{n-1}$$

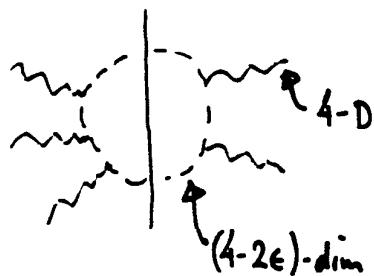
- Ambiguity at 5-pt :  $\frac{E(1234)}{\langle 12 \rangle \langle 23 \rangle \langle 34 \rangle \langle 45 \rangle}$  is collinear-finite
- No proof of uniqueness beyond 5-pt, but believed true.

- b) Cuts at  $\mathcal{O}(\epsilon)$ .

$A_n$  has an overall  $g^{-\epsilon}$ , so at  $\mathcal{O}(\epsilon)$  the pure rational pieces do involve cuts. (logs appear).

Need tree amplitudes & loop integrals to  $\mathcal{O}(\epsilon)$

↑  
explicitly    as ↑ formal objects



At one loop, momentum in  $D=4-\epsilon$  is equivalent to integrating over masses,

$$\int d^{4-2\epsilon} p \frac{1}{p^2 (p-g_1)^2 (p-g_2)^2 \dots (p-g_{n+1})^2}$$

$$= \int d^4 p (\mu^2)^{-1-\epsilon} d\mu^2 \frac{1}{(p^2 - \mu^2) [(p-g_1)^2 - \mu^2] \dots [(p-g_{n+1})^2 - \mu^2]}$$

so we need amplitudes with one massive scalar pair, remaining legs massless.

Example.  $A_4(++++)$

Need  $A_4^{\text{tree}}(1_s, 2^+, 3^+, 4_s) = -i \frac{\mu^2 [23]}{\langle 23 \rangle (2k_3 \cdot k_2)}$

$$\text{Abspl}_{S_{12}} [A_4^{1\text{-loop}} (1^+, 2^+, 3^+, 4^+)] =$$

$$\text{Abspl}_{S_{12}} \left[ 2 \frac{[12][34]}{\langle 12 \rangle \langle 34 \rangle} \int (\mu^2)^{1-\epsilon} d\mu^2 \quad \mu^4 \times \begin{array}{c} \text{square loop diagram} \\ \text{with internal lines } 1, 2, 3, 4 \end{array} \right]$$

$$= \left[ 2 \frac{[12][34]}{\langle 12 \rangle \langle 34 \rangle} \frac{1}{16\pi^2} \frac{\epsilon (4\pi)^6}{\Gamma(1-\epsilon)} \int da_i \delta(\sum a_i - 1) \int_0^\infty d\mu^2 \frac{(\mu^2)^{1-\epsilon}}{[sa_1a_3 + ta_2a_4 + \dots]} \right]$$

$$= \left[ 2 \frac{[12][34]}{\langle 12 \rangle \langle 34 \rangle} \frac{1}{16\pi^2} \epsilon(1-\epsilon)(4\pi)^6 I_4^{D=8-2\epsilon} \right] |_{S_{12} \text{ cut}}$$

$$I_4^{D=8-2\epsilon} \sim \frac{r_p}{6\epsilon} + O(\epsilon^0)$$

so we obtain

$$A_4^{1\text{-loop}}(1^+, 2^+, 3^+, 4^+) = \frac{i}{48\pi^2} \frac{[12][34]}{\langle 12 \rangle \langle 34 \rangle} = \frac{i}{48\pi^2} \frac{[23][41]}{\langle 23 \rangle \langle 41 \rangle}$$

in agreement with previous explicit computations

# $Z \rightarrow \text{Jets}$

3-jet    used in measurement of  $\alpha_S$   
2-jet

Theoretical uncertainties dominate in this extraction  
 → need NNLO calculation

## Ingredients

- $Z \rightarrow g\bar{g}g$  at 2 loops      Hard problem
- $Z \rightarrow g\bar{g}g$  at 1 loop      Glover & Miller      in progress  
 $Z \rightarrow g\bar{g}f\bar{f}'$
- $Z \rightarrow f\bar{f}g\bar{g}$  at tree level      Berends, Giele & Kuijf NPB324:  
 $Z \rightarrow f\bar{f}f\bar{f}'$
- Technology for IR cancellations at NNLO in a general-purpose numerical program — generalization of Giele & Glover (PRD46:1980)      Non trivial but feasible

## $Z \rightarrow 4 \text{ jets at NLO}$

Uses 2nd ingredient above; lowest order process in which gluon & quark color charges can be measured experimentally!!

Can exclude light neutral colored fermions or scalars.

## A Leading-Color Amplitude

$$A_{4+V}^{\text{tree}}(1_q^+, 2^+, 3^-, 4_{\bar{q}}^-; 5_i^-, 6_i^+) = \\ i \left[ - \frac{[1|2] \langle 1|3\rangle \langle 4|5\rangle \langle 6^+|(1+2)|3^+\rangle}{\langle 1|2\rangle s_{23} t_{123} s_{56}} + \frac{\langle 3|4\rangle [2|4] [1|6] \langle 5^-|(3+4)|2^-\rangle}{[3|4] s_{23} t_{234} s_{56}} - \frac{\langle 5^-|(3+4)|2^-\rangle \langle 6^+|(1+2)|3^+\rangle}{\langle 1|2\rangle [3|4] s_{13} s_{56}} \right]$$

$$A_{4+V,V}(1_q^+, 2^+, 3^-, 4_{\bar{q}}^-; 5_i^-, 6_i^-) = V A_{4+V}^{\text{tree}} + F_q + F_s$$

$$V = -\frac{1}{\epsilon^2} \left( \left( \frac{\mu^2}{-s_{12}} \right)' + \left( \frac{\mu^2}{-s_{23}} \right)' - \left( \frac{\mu^2}{-s_{34}} \right)' \right) - \frac{3}{2\epsilon} \left( \frac{\mu^2}{-s_{56}} \right)' - \frac{7}{2}$$

$$\text{flip} = \{1 \rightarrow 4, \quad 2 \rightarrow 3, \quad 5 \rightarrow 6; \quad (\cdot) \rightarrow [\cdot]\}\,.$$

Define some auxiliary functions.

$$\begin{aligned} L_0(r) &= \frac{\ln(r)}{1-r}, & L_1(r) &= \frac{L_0(r)+1}{1-r}, & L_2(r) &= \frac{\ln(r)-\frac{1}{2}(r-1/r)}{(1-r)^3}. \\ L_{S-1}(r_1, r_2) &= \text{Li}_2(1-r_1) + \text{Li}_2(1-r_2) + \ln r_1 \ln r_2 - \frac{\pi^2}{6}, \\ Ls_{-1}^{2m}(s, t; m_1^2, m_2^2) &= -\text{Li}_2\left(1 - \frac{m_1^2}{t}\right) - \text{Li}_2\left(1 - \frac{m_2^2}{t}\right) - \frac{1}{2} \ln^2\left(\frac{-s}{-t}\right) + \frac{1}{2} \ln\left(\frac{-s}{-m_1^2}\right) \ln\left(\frac{-s}{-m_2^2}\right) \\ &\quad + \left[\frac{1}{2}(s - m_1^2 - m_2^2) + \frac{m_1^2 m_2^2}{t}\right] I_3^{3m}(s, m_1^2, m_2^2). \end{aligned}$$

Then

$$\begin{aligned}
 F_g = & b_1 \text{Ls}_{-1} \left( \frac{-s_{23}}{-t_{234}}, \frac{-s_{34}}{-t_{234}} \right) + b_2 \text{Ls}_{-1}^{2m^h}(s_{34}, t_{123}; s_{56}, s_{12}) + b_4 \text{Ls}_{-1}^{2m^h}(s_{12}, t_{234}; s_{56}, s_{34}) \\
 & + b_5 \text{Ls}_{-1} \left( \frac{-s_{12}}{-t_{123}}, \frac{-s_{23}}{-t_{123}} \right) + t I_3^{3m}(s_{12}, s_{34}, s_{56}) \\
 & - 2 \frac{\langle 1 3 \rangle \langle 6^+ | (1+2) | 3^+ \rangle}{\langle 1 2 \rangle [5 6] \langle 4^+ | (2+3) | 1^+ \rangle} \times \left[ \frac{\langle 6^+ | (2+3) | 2^- \rangle}{t_{123}} \frac{\text{Lo} \left( \frac{-s_{23}}{-t_{123}} \right)}{t_{123}} + \frac{[6 4] \langle 4 3 \rangle}{\langle 2 3 \rangle} \frac{\text{Lo} \left( \frac{-s_{56}}{-t_{123}} \right)}{t_{123}} \right] \\
 & - 2 \frac{[4 2] \langle 5^- | (3+4) | 2^- \rangle}{[4 3] (6 5) \langle 1^- | (2+3) | 4^- \rangle} \times \left[ \frac{\langle 5^- | (2+3) | 4 | 3^+ \rangle}{t_{234}} \frac{\text{Lo} \left( \frac{-s_{23}}{-t_{234}} \right)}{t_{234}} + \frac{\langle 5 1 \rangle [1 2]}{[3 2]} \frac{\text{Lo} \left( \frac{-s_{56}}{-t_{234}} \right)}{t_{234}} \right], \\
 F_{s1} = & - \frac{\langle 1 3 \rangle}{2 \langle 1 2 \rangle \langle 2 3 \rangle [5 6] t_{123} \langle 4^+ | (2+3) | 1^+ \rangle} \left[ (\langle 6^+ | (2+3) | 2 | 3^+ \rangle)^2 \frac{\text{Lo} \left( \frac{-s_{23}}{-t_{123}} \right)}{s_{23}^2} + ([6 4] \langle 4 3 \rangle t_{123}) \frac{(\langle 6^+ | (2+3) | 2 | 3^+ \rangle)^2}{s_{23}^2} \right] \\
 & + \frac{1}{2} \frac{[6 2]^2}{\langle 1 2 \rangle [2 3] [3 4] [5 6]}
 \end{aligned}$$

$$F_s = F_{s1} + \text{flip}(F_{s1}).$$

where

$$\begin{aligned}
 b_1 &= b_5|_{\text{flip}}, \\
 b_2 &= \frac{\langle 1 3 \rangle (\langle 6^+ | (1+2) | 3^+ \rangle)^2}{\langle 1 2 \rangle \langle 2 3 \rangle [5 6] t_{123} \langle 4^+ | (2+3) | 1^+ \rangle} + \frac{[1 2]^2 (4 5)^2 \langle 2^+ | (1+3) | 4^+ \rangle}{[2 3] \langle 5 6 \rangle t_{123} \langle 1^+ | (2+3) | 4^+ \rangle \langle 3^+ | (1+2) | 4^+ \rangle}, \\
 b_4 &= b_2|_{\text{flip}}, \\
 b_5 &= \frac{\langle 1 3 \rangle (\langle 6^+ | (1+2) | 3^+ \rangle)^2}{\langle 1 2 \rangle \langle 2 3 \rangle [5 6] t_{123} \langle 4^+ | (2+3) | 1^+ \rangle} + \frac{[1 2]^3 (4 5)^2}{[2 3] [1 3] \langle 5 6 \rangle t_{123} \langle 1^+ | (2+3) | 4^+ \rangle}.
 \end{aligned}$$

$$\begin{aligned}
 t = & \frac{t_{123} (s_{56} + s_{12} - s_{34}) - 2s_{12}s_{56}}{2t_{123}} b_2 + \frac{1}{2} \frac{[1 2] \left[ (\langle 4^- | (1+2)(3+4) | 5^+ \rangle)^2 - s_{12}s_{34} \langle 4 5 \rangle^2 \right]}{\langle 1 2 \rangle [3 4] \langle 5 6 \rangle \langle 1^+ | (2+3) | 4^+ \rangle \langle 3^+ | (1+2) | 4^+ \rangle} \\
 & + (\text{flip of above terms}).
 \end{aligned}$$

## SUMMARY

Unitarity-based rules are a powerful and efficient way of performing one-loop calculations in QCD.

Collinear limits give strong constraints on amplitudes, and simple forms for rational pieces.

They've been used to calculate a variety of processes:

$\ell \rightarrow g\bar{g}ggg$ ,  $O \rightarrow Zg\bar{g}g'\bar{g}'$ ,  $O \rightarrow Zg\bar{g}gg$ , and all-n equations for N=4 + N=1 SUSY MHV.

Techniques are starting to be used for two-loop calculations – listen to the next talk!