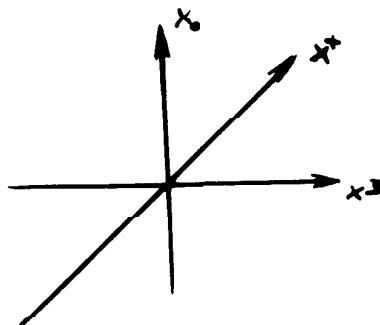
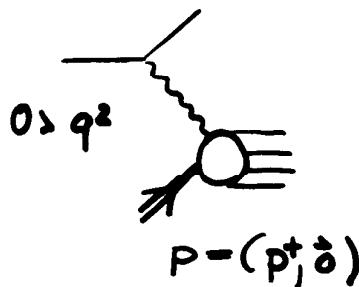


# Intrinsic Glue @ very small $x_{Bj}$

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- ▶ small  $x_{Bj}$  & high parton densities
- ▶ classical color charge & classical gluon fields  
resummation tools
- ▶ R.G. program: status report
- ▶ Summary

## ■ The problem



$$x_{Bj} = -\frac{q^2}{2p \cdot q}$$

$$\alpha_s \ln \frac{1}{x_{Bj}}$$

BFKL resummation  $\rightarrow x_{Bj} g(x_{Bj}, Q^2) \sim \frac{1}{x_{Bj}^\#} \alpha_s$

based on perturbatively  
leading  $\alpha_s \ln \frac{1}{x_{Bj}}$

unitarity bounds?

experimentally seen?

recombination effects @ high densities

## ■ A solution?

perturbative

► small  $\alpha_s$  (?)

► simple (naive?) model  
for 'valence' part:

classical color charge  $\rightarrow$  full classical field  
stat. weight

'nonperturbative'

► resum high density effects

## • Motivation:

► longitudinal properties

large  $x_{gj}$  as seen by small  $x_{gj}$

- classical colored particles

- couple to smaller  $x_{gj}$  cikonally

- growing numbers  $\downarrow x_{gj}$

- independent on time scales resolved

►

transverse properties

- large classical charge densities



classical gluon fields

source of  
color charge

Wilson lines

Wilson R.G.

statistical  
ensemble

high density  
renormalization

## ■ Effective action

$$S_{pt} = i \int d^4x F[g] - \int d^4x \frac{G[A]}{4} + \frac{i}{N_c} \int d^2x_1 d\bar{x} \delta(x^-) \text{tr}[g(x)] W_{\infty\infty}[A](x, x_\perp)$$

- $F[g]$  statistical weight

inv.  $\rho \rightarrow U(x_\mu) \rho U^\dagger(x_\mu)$

to be determined  
via R.G.

Gaussian approx.

$$\frac{\text{tr } \rho^2(x^-, x_1)}{2 M(x^-)}$$

- $\text{tr } g(x_\perp) W_{\infty\infty}[A]$

gauge inv. generalization of  $J^+ A^-$

see below

$$A^+ = \alpha \text{ gauge}$$

meaningful  $p^+$  cutoffs

(keep res. gauge inv.;  
consistency checks )

$\uparrow$   
R.G.

## Y.M. Equations

$$0 = \int_{P^+} D[S, A] \frac{\delta}{\delta A_\nu} e^{iS[S, A]}$$

$$- \int_{P^+} D[S, A] \left( [D_\mu[S], G^\mu[A]] - J^\mu[S, A] \right) e^{iS[S, A]}$$

$$J^+[\rho, A] = \frac{g}{N_c} \text{tr} \left( g(x; x_\perp) W_{-\infty x^+} [A] t_a W_{x^+ \infty} [A] \right) t_a$$

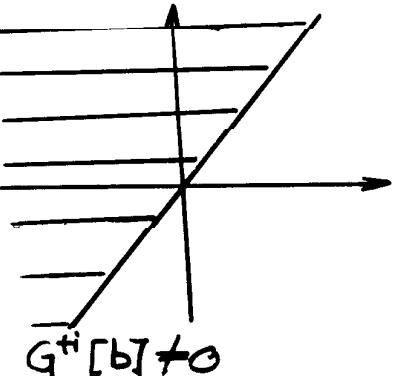
$$[D_\mu[A], J^\mu[\rho, A]] = 0$$

## Classical solution reparametrized

$A^+ = 0$ , static gauge

$$b_i = -\frac{i}{q} U_{-\infty x^-} [\Lambda] \partial_i U_{x^- \infty} [\Lambda]$$

$$U_{x^- \infty} [\Lambda] := T \exp -iq \int_{-\infty}^{x^-} dy^- \Lambda(y^-; x_\perp)$$



new Y.M.:

$$\nabla^2 \Lambda(x^-; x_\perp) = U_{x^- \infty} [\Lambda] g(x^-; x_\perp) U_{-\infty x^-} [\Lambda]$$

## Gluon content of the classical field

$$\rightarrow F[g] \xrightarrow{\text{Gaussian approx.}} \frac{\text{tr } g^2(x_1, x_2)}{2 \mu^2(x)}$$

$$\begin{aligned} \rightarrow D[g] & e^{- \int dy \int dx_2 \frac{\text{tr } g^2(y, x_2)}{2 \mu^2(y)}} \\ & - \int J[\Lambda] \exp - \int dy \int dx_2 \frac{\text{tr } (\nabla^2 \Lambda(y, x_2))^2}{2 \mu^2(y)} \end{aligned}$$

$$\rightarrow G_{ij}^{\text{class}}(y_1, x_1, y'_1, x'_1) = \langle b_i(y_1, x_1) b_j(y'_1, x'_1) \rangle$$

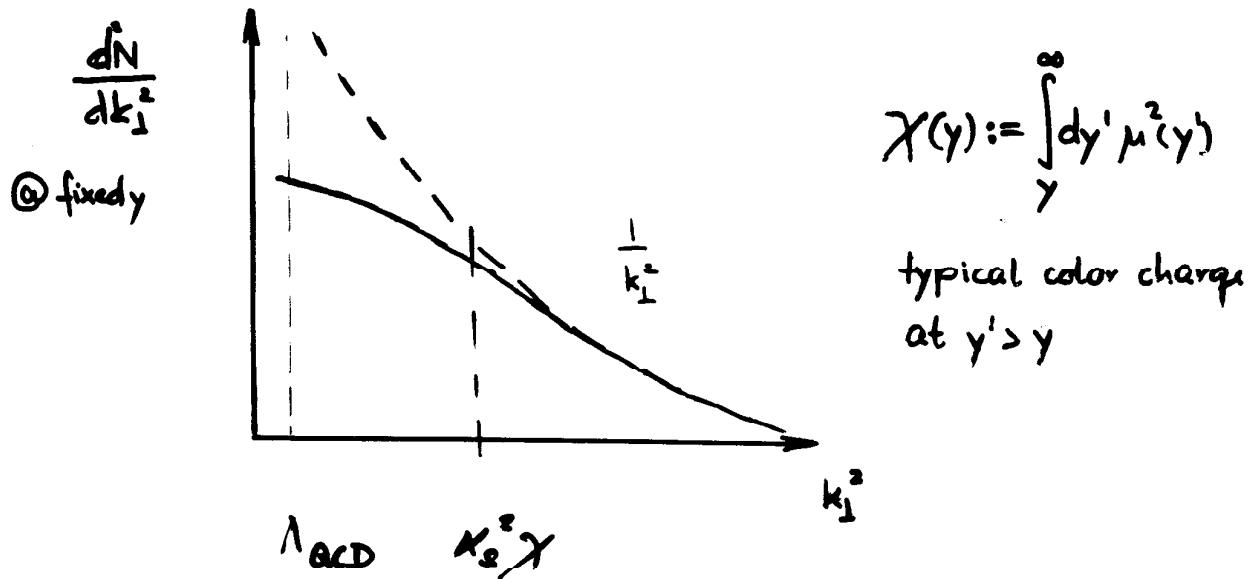
$$b_j = \int_y dy' U_{\alpha y} [\Lambda](x_1) (\nabla^i \Lambda(y'_1, x_1)) U_{y'_1 \alpha} [\Lambda](x_1)$$

$$\rightarrow \langle \Lambda(y_1, x_1) \Lambda(y'_1, x'_1) \rangle = \mu^2(y_1, Q^2) \delta(y-y') \frac{1}{4\pi} \gamma(x_1, x'_1)$$

$$\frac{1}{4\pi} x_1^2 \ln x_1^2 \Lambda_{QCD} + \gamma_1(0) =: \gamma(x_1)$$

$$G_{ij}^{ab}(y, x_\perp; y', x'_\perp) = -g^{ab} (\nabla_i \nabla'_j \delta_1(x_\perp - x'_\perp)).$$

$$\cdot \frac{1 - \exp^{-N_c \chi(y)} [\gamma_1(x_\perp - x'_\perp) - \gamma_1(\infty)]}{N_c [\gamma_1(x_\perp - x'_\perp) - \gamma_1(\infty)]}$$



Lessons to be learned

- classical fields  $\longleftrightarrow$  new scale  $\alpha_s^2 \chi(y)$
- pert. limit @ short distances
- softening at small  $k_\perp$ :  $G_{ii}^{aa} \sim \ln \frac{k_\perp^2}{\delta^2 \chi(y)}$
- $\lambda_{\text{QCD}}$  dep. weak unless  $k_\perp^2 \ll \alpha_s^2 \chi(y)$

## ■ R.G. improvement

$$0 - \int_{p^+} D_{\mu} [g, A] \left( [D_\mu [A], G'''[A]] - J''[g, A] \right) e^{i S[g, A]}$$

$$A = b + a + SA \xrightarrow{P^+ < k^+ < P^+}$$

class. sol.

$$\begin{aligned}
 & \left( D[s, A] \int_{\mathbb{R}^3 k + \epsilon p} D[\delta A] \left[ J_\mu[b+a], G^\mu[b+a] \right] - J^\nu[b+a] \right. \\
 & \quad \left. - i q [s_A, G^\mu[b+a]] - i q [s_A, G^\mu[\delta A]] \right. \\
 & \quad \left. + \dots \right) e^{is[s, b+a + \delta A]} \\
 & \xrightarrow{\delta s} \delta J^\nu[b+a] \rightarrow \delta \rho
 \end{aligned}$$

induced current

perturbatively

$$g = f^{abc} \left( g^{+b} s a^c + g^{-b} \int dy^+ \theta^+ s a^{-c} \right) +$$

$$+ f^{abc} \delta A^b \partial^+ \delta A^c$$

$$+ \frac{q^2}{N_c} g^6 " \int dy^+ dz^+ \Theta\Theta h tttt " \delta A^{-c} \delta A^{-d}$$

■ procedure

$$\int D[\delta A] @ \text{fixed } g \& \delta g$$

$$\int D[g] @ \text{fixed } g' := g + \delta g$$

■ new effective action

$$S[g', a] = -F'[g'] - \frac{i}{4} G^2[a] + i/N_c + g' W[a]$$

with

$$\exp(-F'[g']) = \int D[g, \delta A] \delta(g' - g - \delta g[\delta A]) \exp\left(-F[g] - \frac{i}{2} \delta A D^2[g] \delta A\right)$$

■ to leading  $\alpha_s \ln \frac{1}{x_{\text{cut}}}$

$$\frac{d}{d \ln \frac{1}{x_{\text{cut}}}} F[g] = \alpha_s \Delta(F[g])$$

■ tools needed

$$\Rightarrow - \sum \text{---} \overline{\text{---}} \text{---} \text{---} \text{---}$$

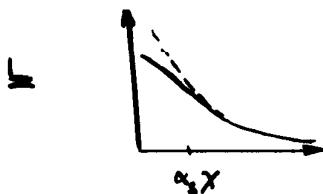
gauge inv. def of  $\delta g$

■ note:  $b_\mu$  retains structure

■ double log approx:

$$\frac{d}{d \ln \frac{1}{x_{\text{bj}}}} F[s] \approx < \overbrace{\text{L}}^{\times} + 2 \overbrace{\text{L}}^{\times} + \overbrace{\text{H}}^{\times} >_s$$

$$= - \sum \text{wavy lines}$$



+  
qualitatively  
similar  
contributions

## ■ Summary

- classical fields resum density effects
- new scale ( $\alpha_s^2 X$  in Gaussian approx)  
at large distances
- perturbative at small distances
- expansion in  $g$  (small densities)  $\rightarrow$  BFKL

## ■ perspectives: To Do list has become shorter

- BG field propagator
- BFKL derivation written up
- gauge 'invariant' treatment of  $fg$
- ...