

Experimental Prospects for Direct Measurements of Neutrino Mass

Outline

Directness?

The Various Techniques: for completeness

Astrophysics/Cosmology (very short)

Nuclear and Particle Physics: heart of the talk

beta decay

double-beta decay

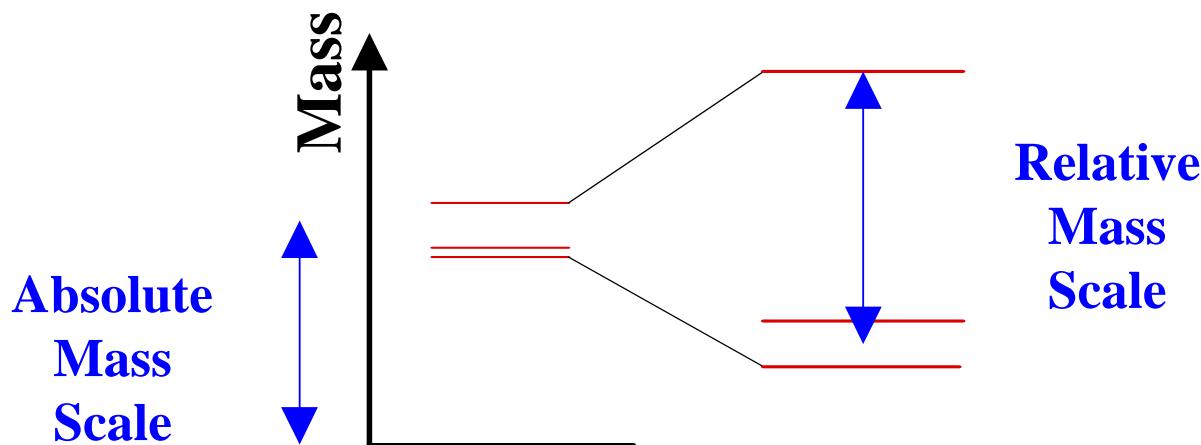
Steve Elliott



Direct vs. Indirect Absolute vs. Relative Kinematic vs. Interference

- No m_ν , experimental technique is direct in that it measures one of the m_i . All experiments measure parameters that depend on the mass eigenvalues.
- So, I took my assignment to be a discussion of laboratory experiments that indicate the “absolute scale of ν mass”.
- These are experiments that search for a kinematic effect due to mass.
- These are not experiments that search for an interference effects. Oscillation experiments provide data on the “relative mass scale”.

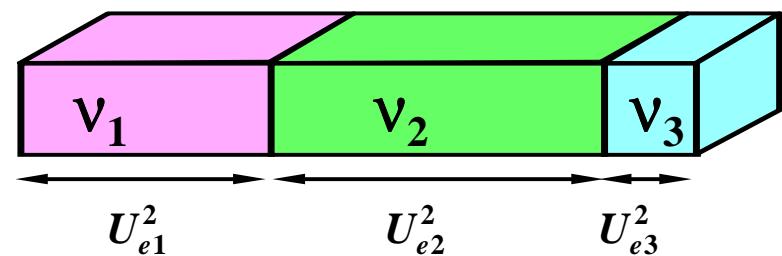
Neutrino Mass: What do we want to know?



Dirac or Majorana

$$\begin{pmatrix} \nu_{\uparrow} \\ \nu_{\downarrow} \\ \bar{\nu}_{\downarrow} \\ \bar{\nu}_{\uparrow} \end{pmatrix} \text{ or } \begin{pmatrix} \nu_{\uparrow} \\ \nu_{\downarrow} \end{pmatrix}$$

ν_e



Mixing

A List of Upcoming Techniques

Present focus in red.

Supernovas - previous speaker, relatively poor sensitivity

Cosmology - previous speaker, hope of reaching
below 100 meV

Nuclear/Particle Physics

τ decay - relatively poor sensitivity

μ decay - relatively poor sensitivity

β decay - hope to reach below 500 meV

$\beta\beta$ decay - hope to reach below 50 meV

oscillations - other speakers, great sensitivity

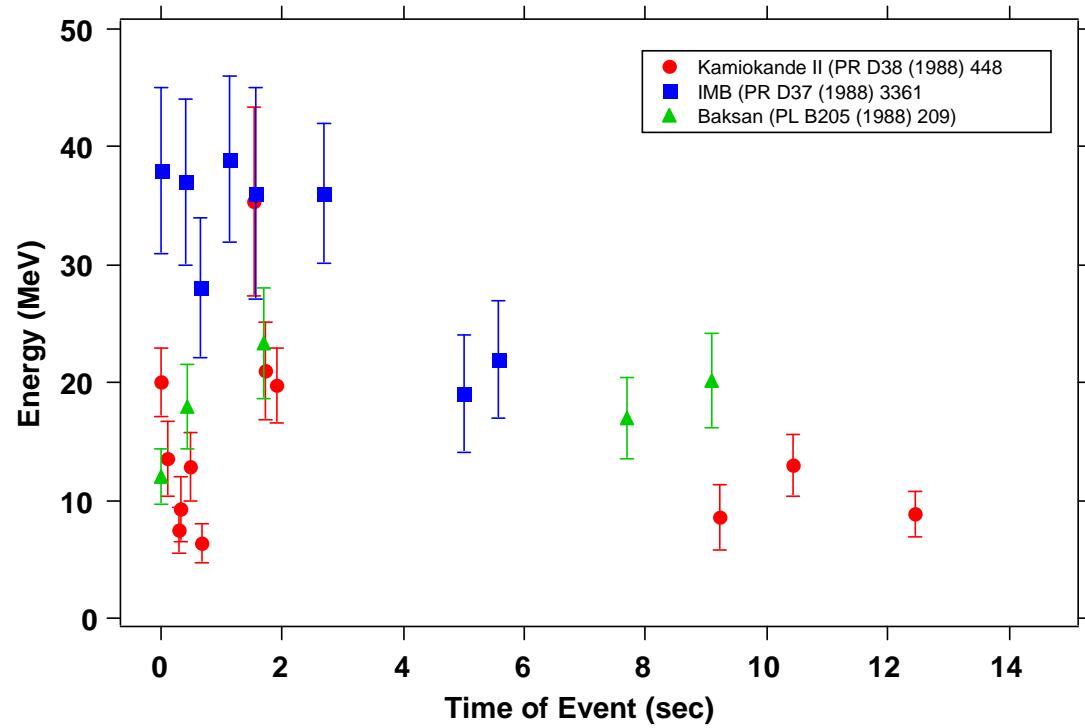
Supernova Tests

Spread of neutrino arrival times can give indication of mass.

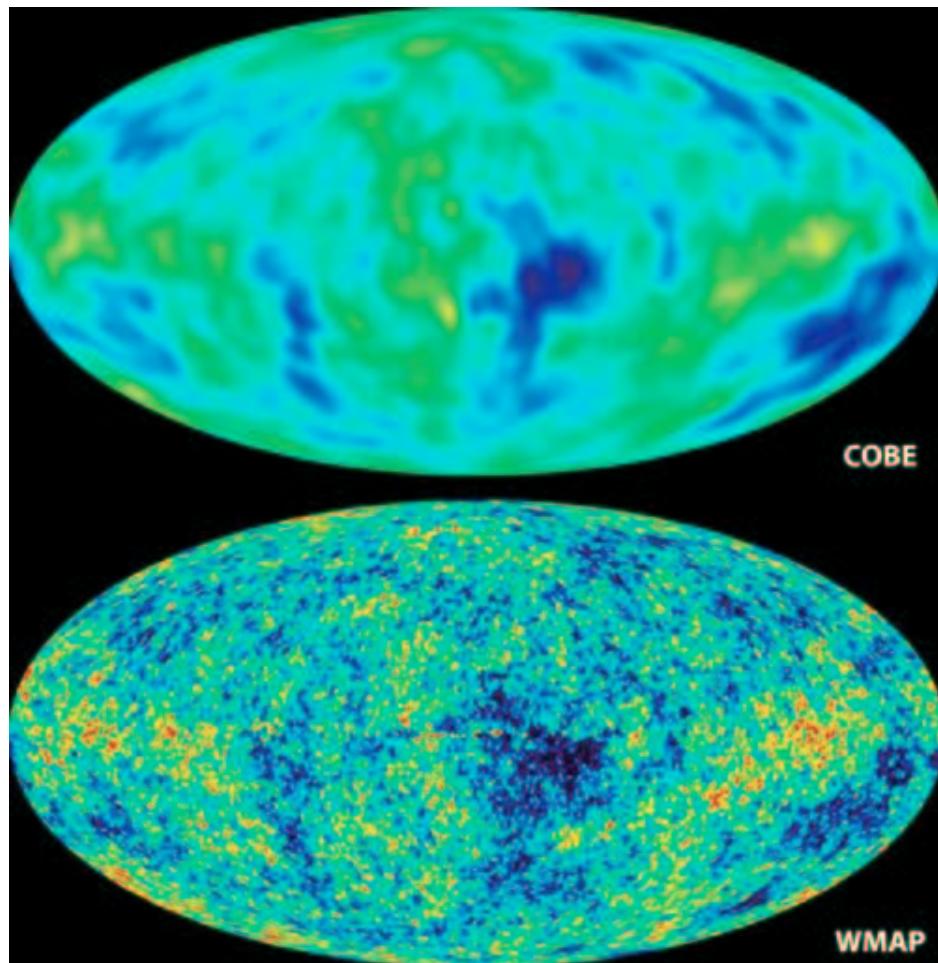
SN1987a: about 20 eV limit but conclusions varied.

SN dynamics makes for model dependencies.

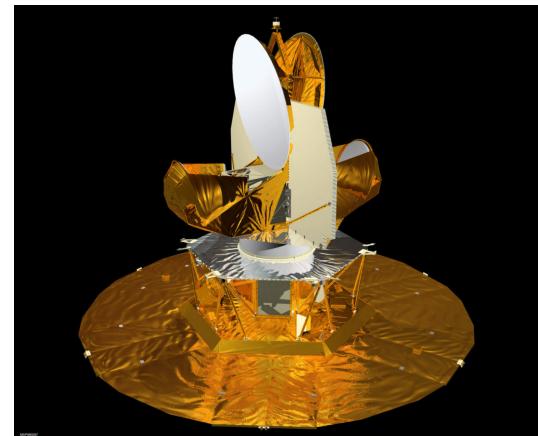
Future sensitivity might be a few eV.



Cosmology Measure $\Omega_v h^2$



Steve Elliott, ANL Neutrino Workshop 2003



- WMAP measured cosmological parameters very precisely. This allowed precise estimates of $\Omega_v h^2$ from LSS measurements.
- WMAP results indicate $\Sigma m_i <$ about 1 eV. A very competitive result.
- But, correlations between parameters result in “arguable” conclusions.
- Want laboratory experiments.

Cosmology - Future Measurements

MAP/PLANCK CMB measurements with high precision galaxy surveys (Sloan Digital Sky Survey): $\sum m_i < \sim 300 \text{ meV}$

If weak lensing by LSS is also considered:

$$\sum m_i < \sim 40 \text{ meV}$$

Even with the correlations, cosmology will play an important role in the interpretation of neutrino mass.

Z-burst and high-energy ν

- Requires, as yet unknown (and unneeded?) flux of UHE ν with energy > Greisen-Zatsepin-Kuzmin (GZK) energy (5×10^{19} eV).
- Although could explain existence of cosmic rays with $E > E_{\text{GZK}}$, it doesn't explain source of proposed UHE ν.
- UHE ν + cosmic relic ν $\rightarrow Z \rightarrow$ hadrons, γs
- m_i near 500 meV could produce p or γ just above GZK cutoff. Model can be “tweaked” to get lower masses.
- Detect p or γ: their multiplicity and energies relate to E^R and hence m_i .

Hope you saw
Sarcevic talk.

$$E_{\nu_i}^R \approx 4.2 \left(\frac{eV}{m_i} \right) \times 10^{21} \text{ eV}$$

Nuclear and Particle Physics Techniques

τ decay: decays into 5 or 6 π most sensitive because of restricted phase space for ν .

$$\tau \rightarrow n\pi + \nu_\tau$$

$$E^* = \frac{m_\tau^2 + m_h^2 - m_i^2}{2m_\tau}$$

m_h is invariant mass of $n\pi$.

E^* is total π energy in τ rest frame.

Obtain data on degenerate scale m_1 .

$$m_{\nu_1} < 18.2 \text{ MeV}$$

Nuclear and Particle Physics Techniques

μ decay: $\pi \rightarrow \mu + \nu$

$$m_{\nu_\mu}^2 = m_\pi^2 + m_\mu^2 - 2m_\pi \sqrt{m_\mu^2 + p_\mu^2}$$
$$m_{\nu_\mu}^2 = \sum |U_{\mu i}^2| m_i^2$$

$$m_{\nu_\mu} < 190 \text{ keV}$$

PR D53 (1996) 6065

^{187}Re β Decay Experiments

- Low Q-value: 2.6 eV
- Long half-life = low specific activity
- Bolometric techniques = measure whole spectrum
- Measure 10^{-10} of decays near endpoint, but bolometer response time is 100s μs .
- Future sensitivity should be about 10 eV.

Milano: Ag ReO_4 NIM A444 (2000) 77

Genoa: Metallic Re, $m_\nu < 26$ eV NP B (proc. Suppl.) 91 (2001) 293

The β Spectrum (bare nucleus)

$$\frac{dN}{dE} = \frac{G_F^2}{2\pi^3 h^7} \cos^2 \theta_C |M|^2 F(E, Z+1) p(E + m_e) \epsilon \sqrt{\epsilon^2 - m_\nu^2} \Theta(\epsilon - m_\nu)$$

Fund. Constants **Energy-independent Matrix Element** **Fermi Function**

Electron momentum, Energy, mass **$\epsilon = E_0 - E$** **Step Function**

β Spectrum (atom or molecule) plus neutrinos mix

Many possible final states. ε_j for state j

$$\frac{dN}{dE} = A|M|^2 F(E, Z+1) p(E + m_e)$$

$$\sum_j W_j \varepsilon_j \sum_i |U_{ei}|^2 \sqrt{\varepsilon_j^2 - m_i^2} \Theta(\varepsilon_j - m_i)$$

W_j - probability for
transition to state j

U_{ei} - mixing
matrix elements

m_i - ν mass
eigenstate

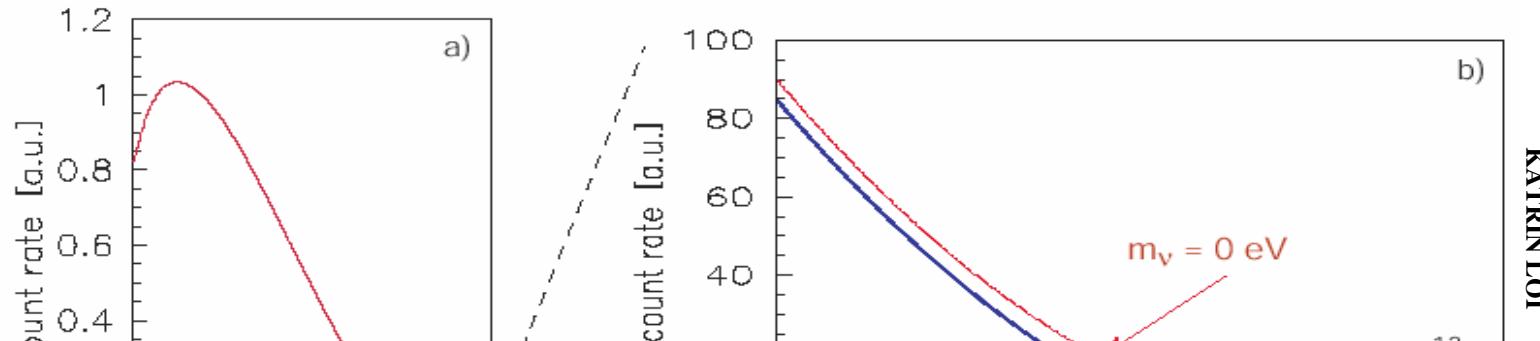
β Spectrum (atom or molecule)

When convolved with resolution function
with width $> m_i$, we can analyze the
spectrum with one mass parameter: $\langle m_\beta \rangle$.

$$\frac{dN}{dE} = A|M|^2 F(E, Z+1) p(E + m_e)$$
$$\sum_j W_j \epsilon_j \sqrt{\epsilon_j^2 - m_\beta^2} \Theta(\epsilon - m_\beta)$$

The Neutrino Mass from β decay

The shape of the β energy spectrum near the endpoint depends on m_ν .



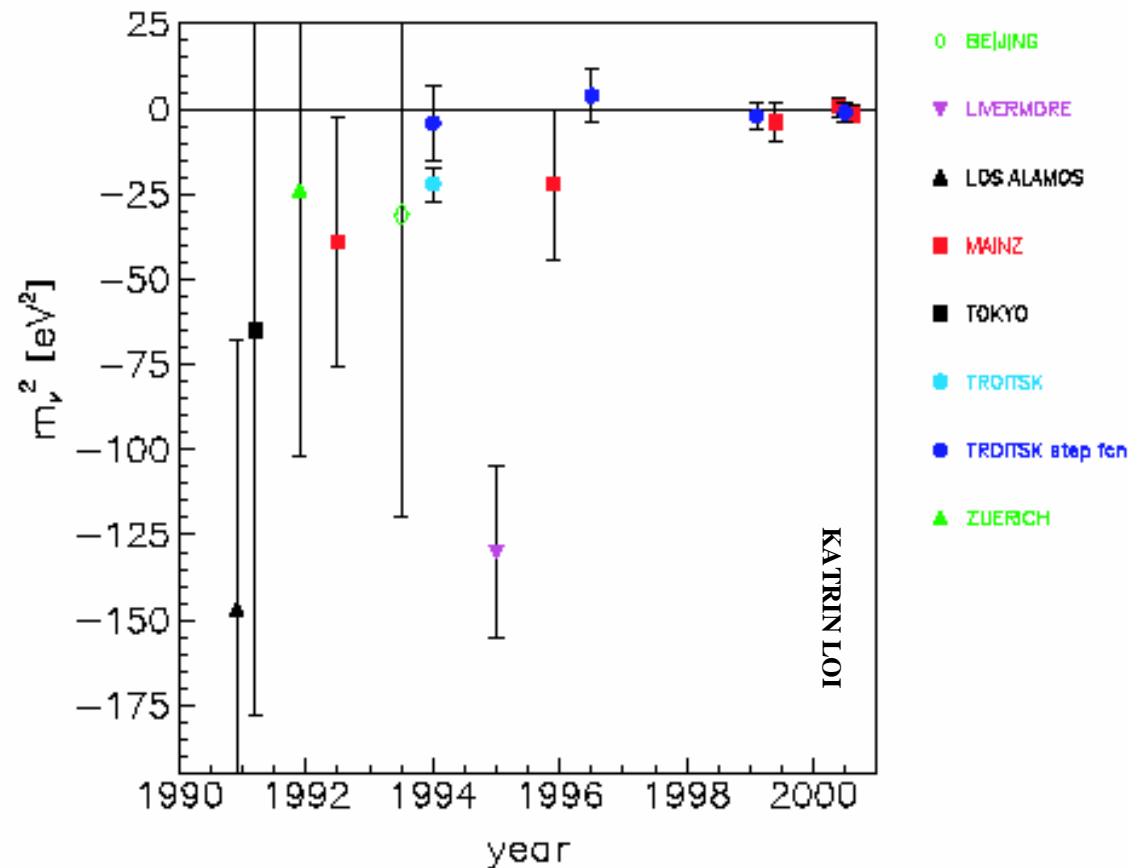
$$\langle m_\beta \rangle = \sqrt{\sum_{i=1}^3 |U_{ei}|^2 m_i^2} < 2.2 \text{ eV}$$

NP B (Proc. Suppl.) 91 (2001), 273

Primary Systematic Uncertainties

Resolution error
leads to m_ν error

$$\frac{dN}{dE} \propto \epsilon^2 - \frac{\langle m_\beta \rangle^2}{2}$$
$$\frac{dN}{dE} \propto \epsilon^2 + \sigma^2$$
$$\delta \langle m_\beta \rangle^2 \approx -2\sigma^2$$



Identifying the Systematics

Mainz: tritium as thin film on flat surface. Temperature activated roughening led to microcrystals in turn leading to large inelastic scattering.

Troisk: Gaseous tritium. Large angle scattering of electrons trapped in source.

Addressing these issues removed the “negative mass squared” in those experiments.

Tritium β decay Experiments

KATRIN

**Very big spectrometer using
gaseous and thin sources. A
big step forward.**

**Univ. of Texas-Austin
 t_2 source in magnetic
free environment.**

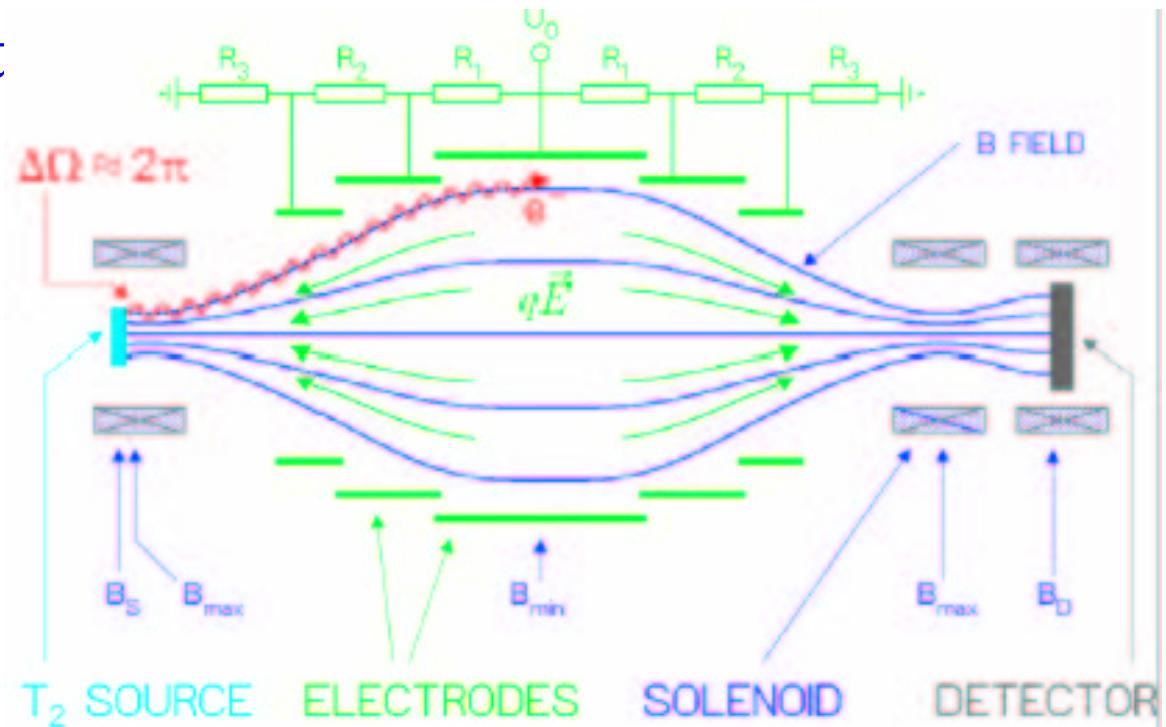
The MAC-E Filter

Magnetic Adiabatic
Collimation followed
by an Electrostat
Filter

High luminosity
Low background
Good energy resolution

Integrating
high-pass filter

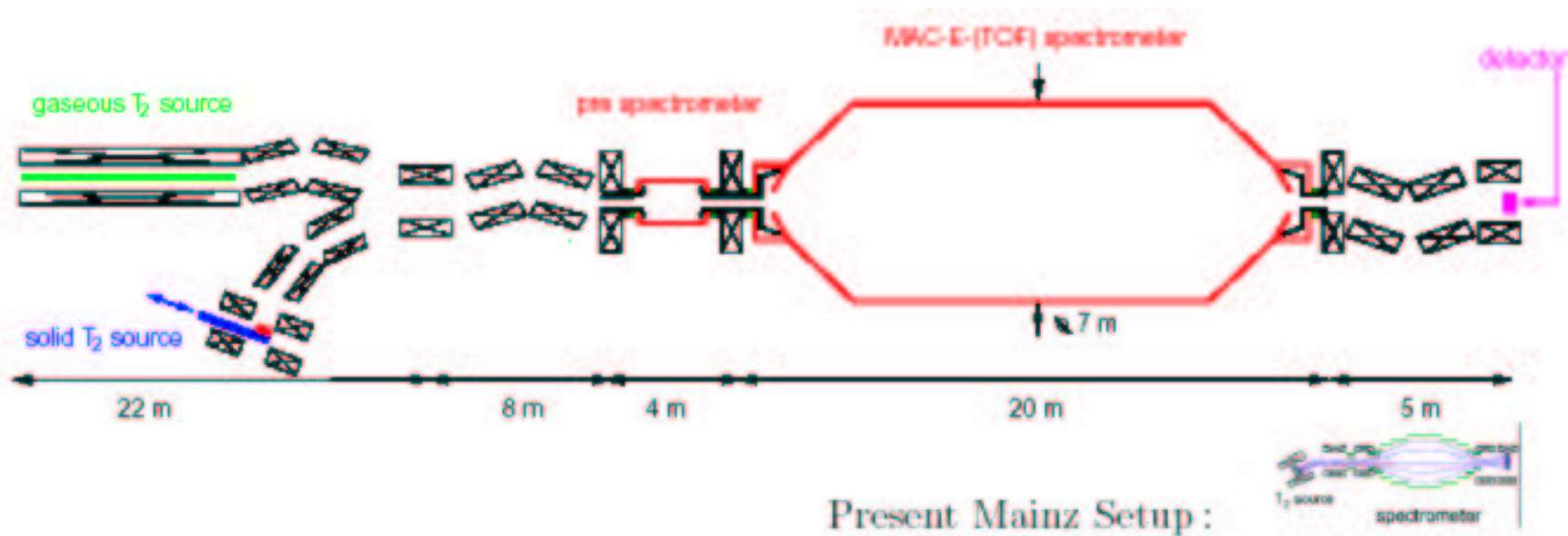
$$\frac{\delta E}{E} = \frac{B_{\min}}{B_{\max}}$$



KATRIN

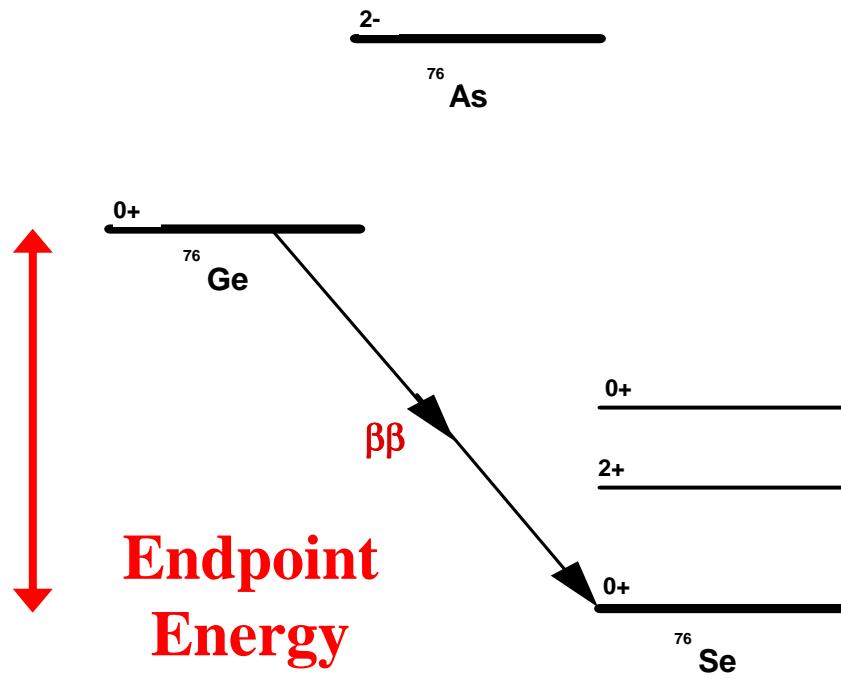
- **1-eV resolution (x4 better)**
- **Electron transport system guides β to pre-spectrometer while preventing tritium flow to spectrometers.**
- **Pre-spectrometer filters out all β except those near endpoint. Keeps residual ionization minimized hence reducing background.**
- **Large analyzing plane increases signal rate**
- **Si drift detectors. 600 eV resolution for 18.6 keV β . Good electron sensitivity but low efficiency for γ .**

KATRIN (LOI version)



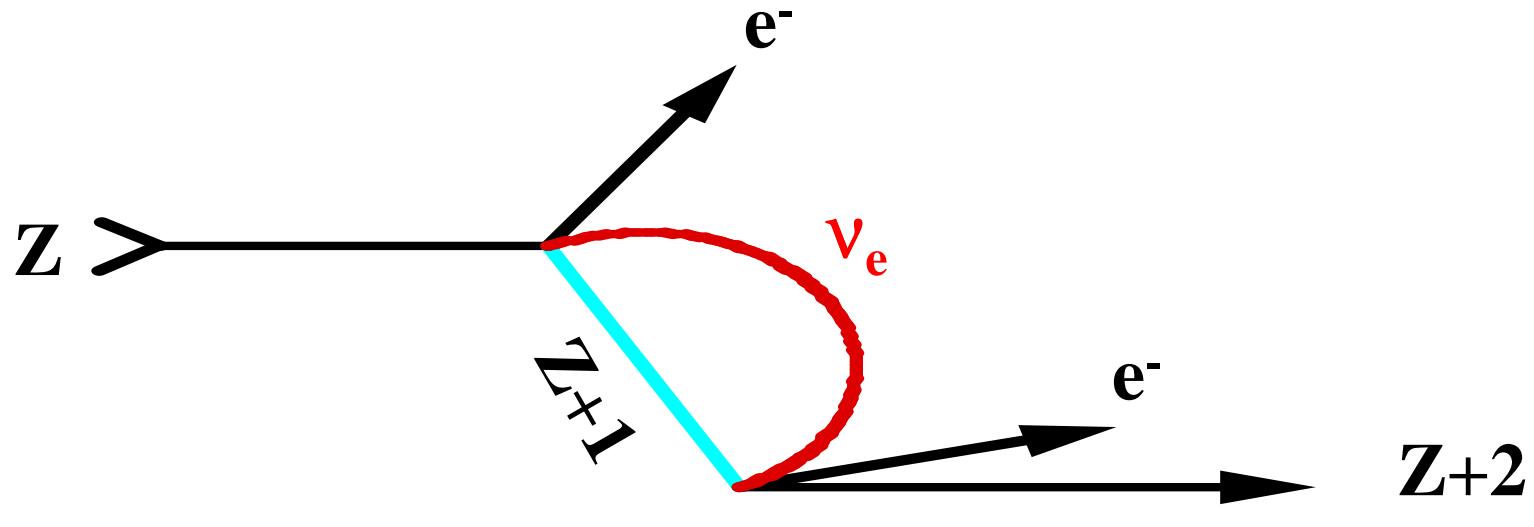
KATRIN will be sensitive to about 350 meV. Thus if the m_i follow a degenerate pattern and m_1 is within the sensitivity, the experiment may see $\langle m_\beta \rangle = m_1$.

Example $\beta\beta$ Decay Scheme



In many even-even nuclei, β decay is energetically forbidden. This leaves $\beta\beta$ as the allowed decay mode.

$\beta\beta(0\nu)$: requires massive Majorana ν



$$n \Rightarrow p + e^- + \bar{\nu}_e$$

$(RH \bar{\nu}_e) \rightarrow (LH \nu_e)$

$$\nu_e + n \Rightarrow p + e^-$$

$\beta\beta$ Decay Rates

$$\Gamma_{2\nu} = G_{2\nu} |M_{2\nu}|^2$$

$$\Gamma_{0\nu} = G_{0\nu} |M_{0\nu}|^2 m_\nu^2$$

G are calculable phase space factors.

$$G_{0\nu} \sim Q^5$$

|M| are nuclear physics matrix elements.

Hard to calculate.

m_ν is where the interesting physics lies.

What about mixing, m_ν & $\beta\beta(0\nu)$?

No mixing:

$$\langle m_{\beta\beta} \rangle = m_{\nu_e} = m_1$$

$$\langle m_{\beta\beta} \rangle = \sum_{i=1}^3 |U_{ei}|^2 m_i \epsilon_i$$

virtual ν exchange
 $\epsilon = \pm 1$, CP cons.

Compare to β decay result: real ν emission

$$\langle m_\beta \rangle = \sqrt{\sum_{i=1}^3 |U_{ei}|^2 m_i^2}$$

Min. $\langle m_{\beta\beta} \rangle$ as a vector sum



$\langle m_{\beta\beta} \rangle$ is the modulus of the resultant vector in the complex plane.
(In this example,
 $\langle m_{\beta\beta} \rangle$ has a min. It
cannot be 0.)

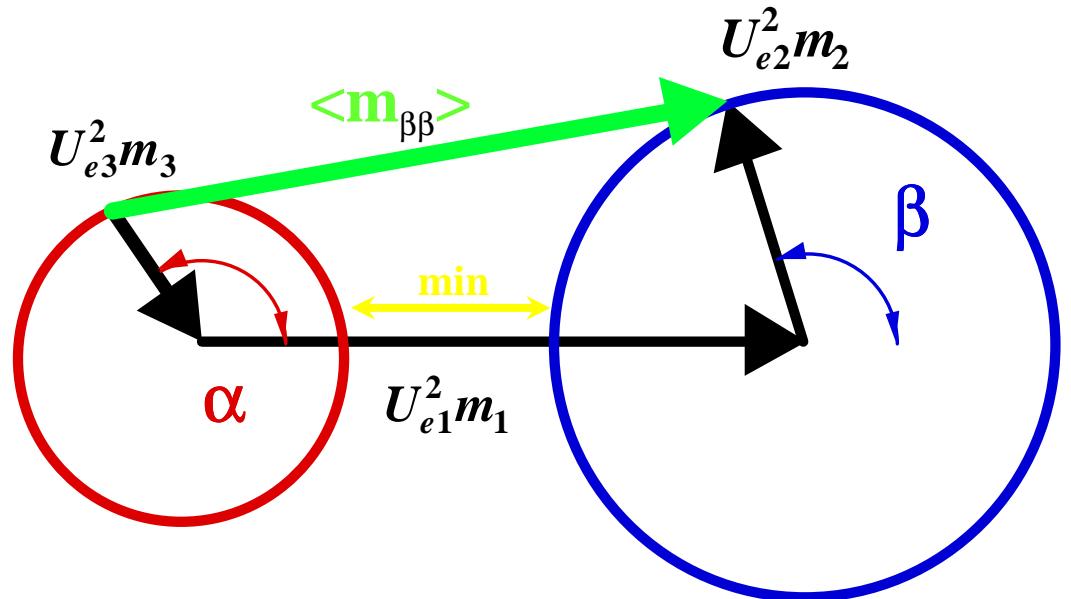
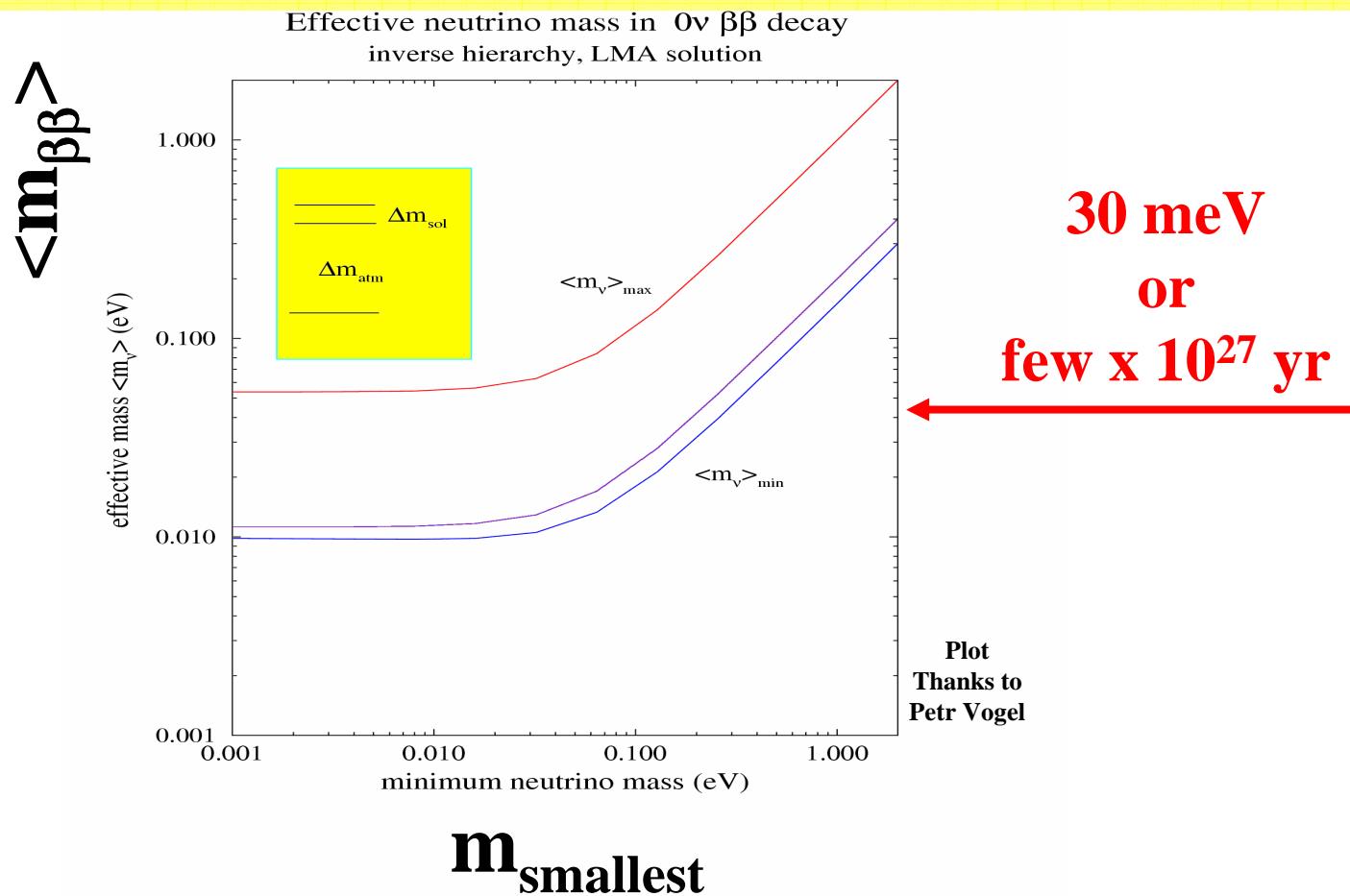
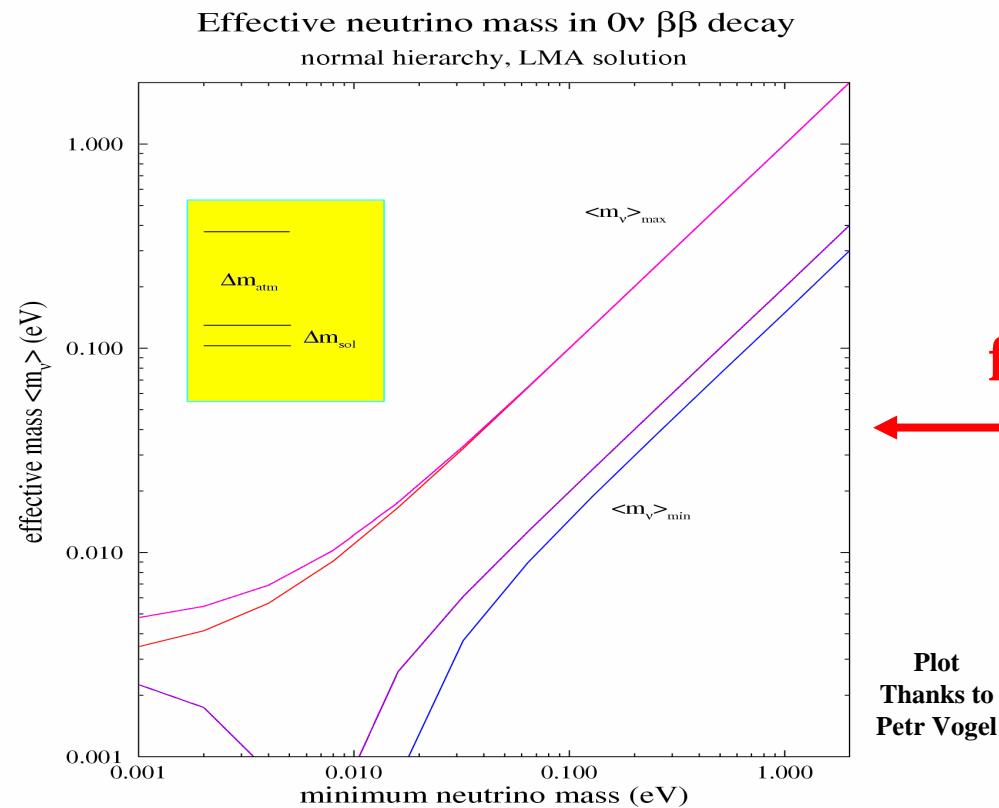


Figure from: PR D63, 073005

More General: 3 ν (Inverted Heir.)



More General: (Normal Heir.)



30 meV
or
few $\times 10^{27}$ yr

Plot
Thanks to
Petr Vogel

Summary of Physics Reach

Even null results have implications!

Normal Hierarchy	Inverted Hierarchy	Degenerate
$m_1 \sim 0$ meV	~ 55 meV	$= M >$ about 100 meV
$m_2 \sim 7$ meV	~ 55 meV	M
$m_3 \sim 55$ meV	~ 0 meV	M
$\langle m_{\beta\beta} \rangle \sim 5$ meV	28 or 55 meV	$M/2$ or M

Solar + KamLAND + Atmospheric ($U_{e3} \sim 0$)

$$\langle m_{\beta\beta} \rangle \approx (0.5)^2 m_1 + \varepsilon_{21} (0.866)^2 \sqrt{m_1^2 + \delta m_{21}^2}$$

An Ideal Experiment

Maximize Rate/Minimize Background

$$\langle m_{\beta\beta} \rangle \propto \left(\frac{b\Delta E}{Mt_{live}} \right)^{\frac{1}{4}}$$

- Large Mass (~ 1 ton)
- Good source radiopurity
- Demonstrated technology
- Natural isotope
- Small volume, source = detector
- Good energy resolution
- Ease of operation
- Large Q value, fast $\beta\beta(0\nu)$
- Slow $\beta\beta(2\nu)$ rate
- Identify daughter
- Event reconstruction
- Nuclear theory

Classes of Background

$\beta\beta(2\nu)$ tail

Need good energy resolution.

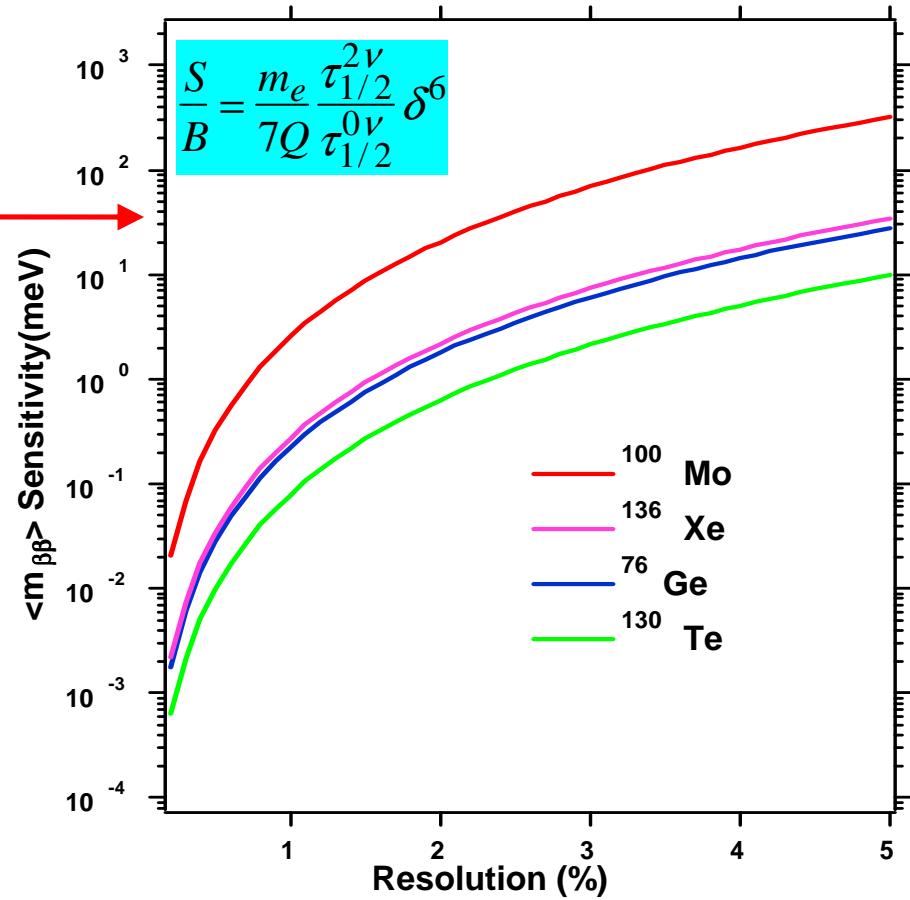
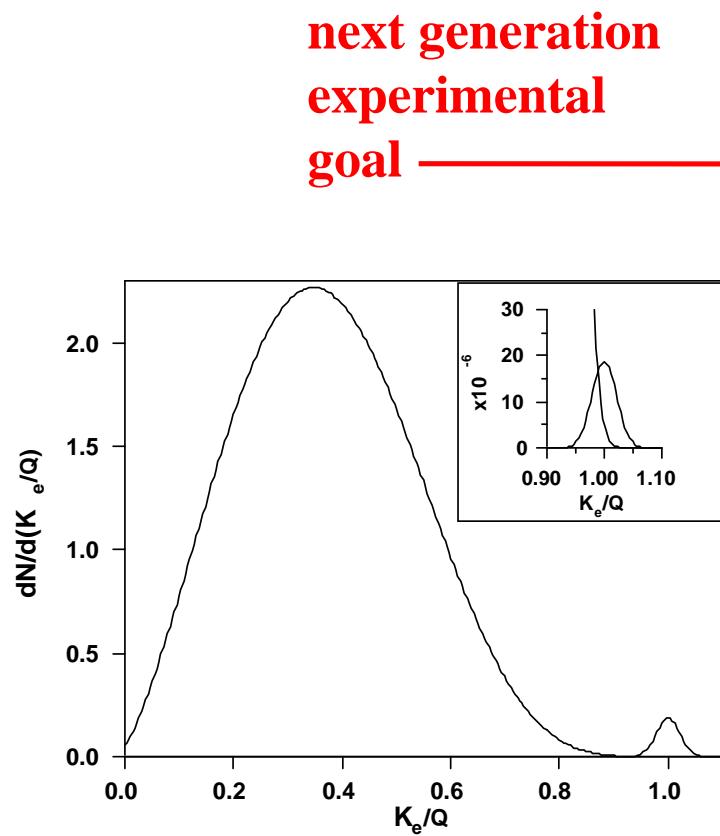
Natural U, Th in source and shielding

Pure materials, identify $\beta\beta$ daughter, pulse shape, timing, position.

Cosmic ray activation

Store and prepare materials underground.

$\beta\beta(2\nu)$ as a Background. Sum Energy Cut Only



Evidence of ^{68}Ge

From: NIM A292 (1990) 337-342.

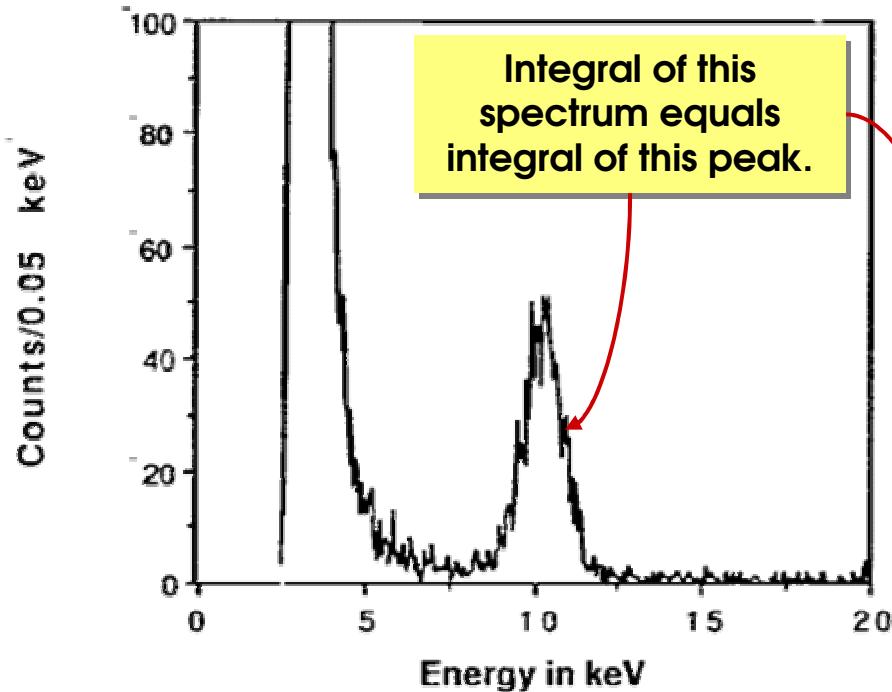


Fig. 7. 10.367 keV gallium K-shell binding energy in one of our current detectors.

Steve Elliott

This peak decays with the right half life.

Experimental data from two 1.05 kg natural detectors

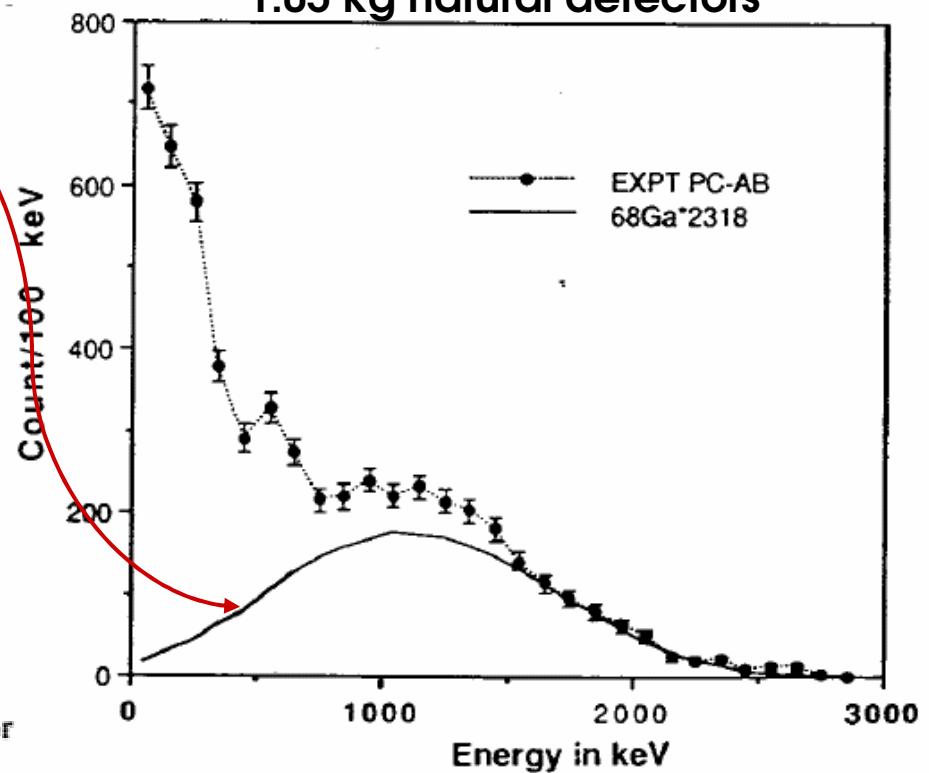


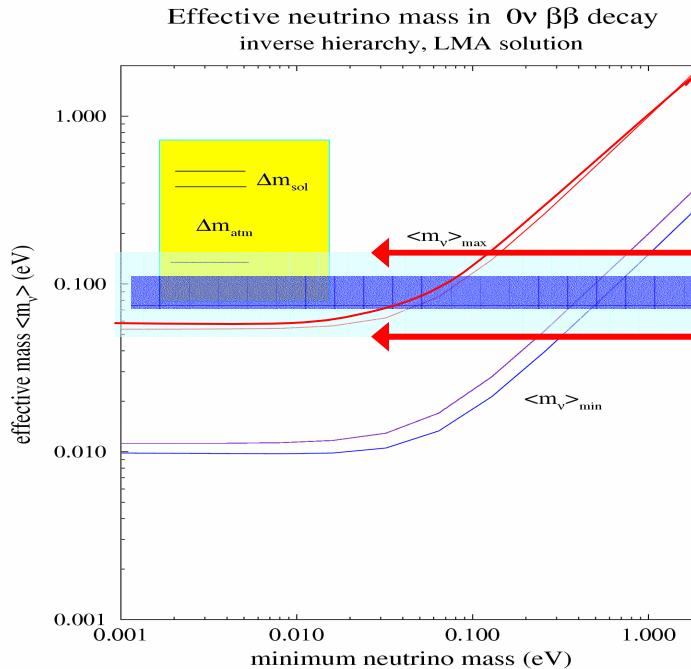
Fig. 8. Current detector data summed into 100-keV bins and compared to a Monte Carlo calculated ^{68}Ga spectrum normalized to ^{68}Ga X-rays.

Matrix Elements

**There are many calculations.
Most authors quote mass limits derived
from all or at least representatives of the
whole range.**

**How do we interpret the uncertainty
associated with the mass due to the
nuclear physics?**

Consider a 100 meV result.



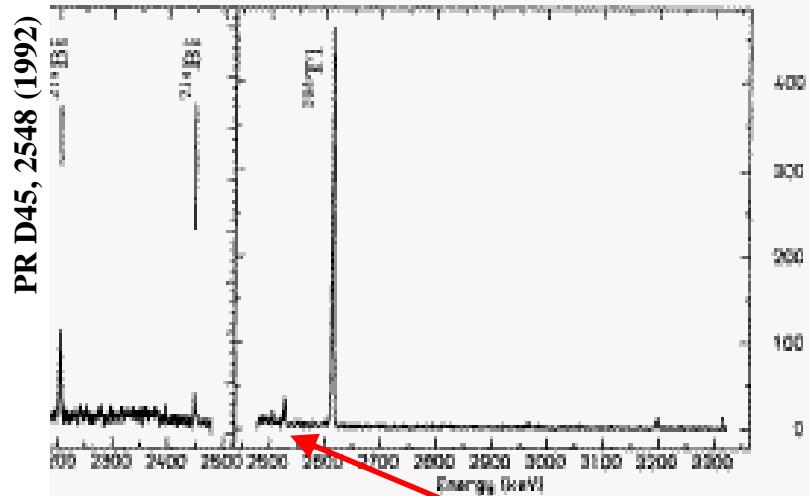
Statistical contribution
to uncertainty.

Matrix Element
contribution to uncertainty.

Would this exclude the inverted hierarchy with small m_{smallest} ?

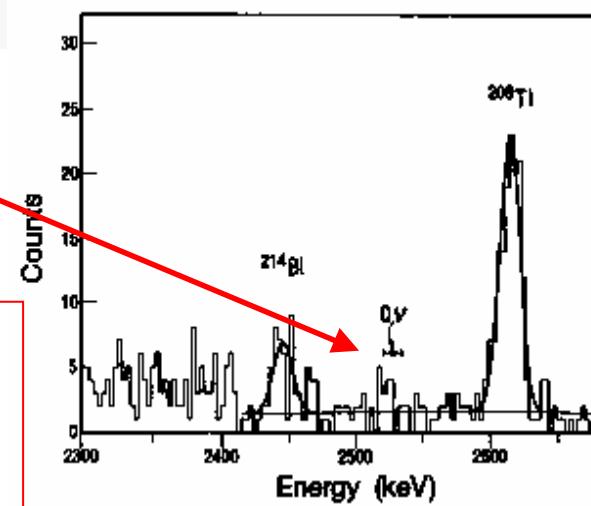
Need improvement in the Theory.

“Found” Peaks



A 2527-keV Ge-det.
peak that was an
electronic artifact.

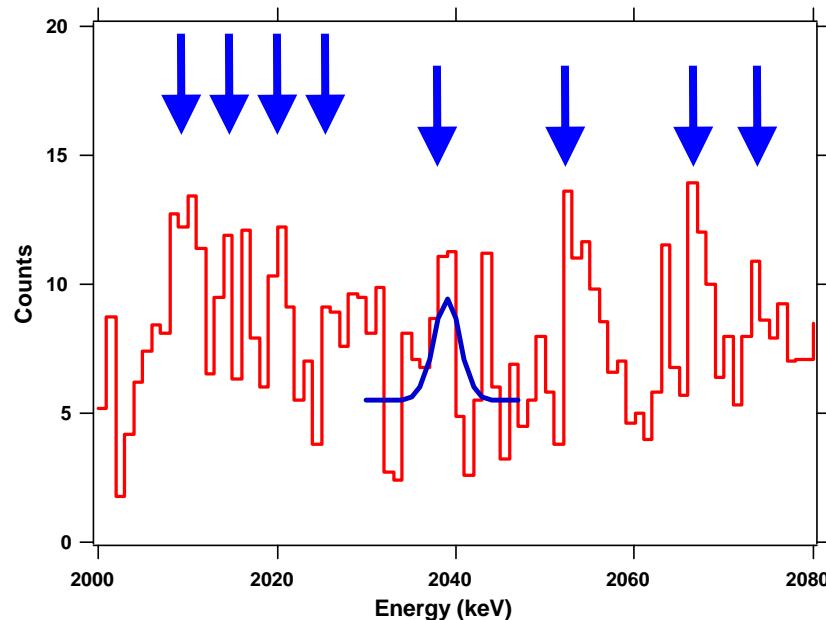
Need more than
one experiment



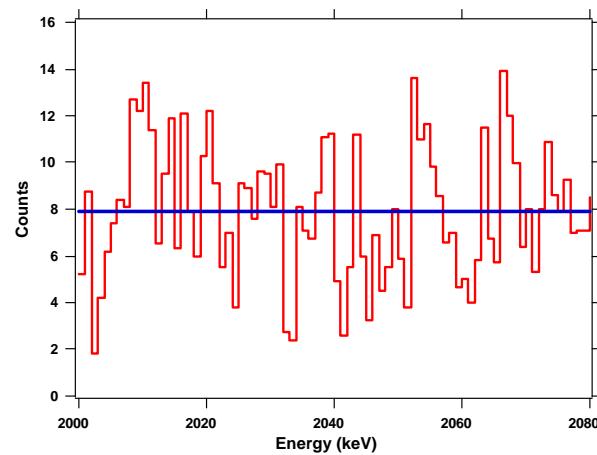
NP B35 (Proc. Supp.), 366 (1994).

A ~2528-keV Te-det.
peak that was a 2σ
Statisticalflucuation.

The Controversy.



Locations of
claimed peaks



Mod. Phys. Lett. A16, 2409 (2001)

If one had to summarize the controversy in a short statement:
Consider two extreme background models:

1. Entirely flat in 2000-2080 keV region.
2. Many peaks in larger region, only $\beta\beta$ peak in small region.

These 2 extremes give very different significances for peak at 2039 keV.
KDHK chose Model 2 but did not consider a systematic uncertainty associated with that choice.

A Great Number of Proposed Experiments

COBRA	Te-130	10 kg CdTe semiconductors
DCBA	Nd-150	20 kg Nd layers between tracking chambers
NEMO	Mo-100, Various	10 kg of $\beta\beta$ isotopes (7 kg of Mo)
CAMEO	Cd-114	1 t CdWO ₄ crystals
CANDLES	Ca-48	Several tons CaF ₂ crystals in liquid scint.
CUORE	Te-130	750 kg TeO ₂ bolometers
EXO	Xe-136	1 ton Xe TPC (gas or liquid)
GEM	Ge-76	1 ton Ge diodes in liquid nitrogen
GENIUS	Ge-76	1 ton Ge diodes in liquid nitrogen
GSO	Gd-160	2 t Gd ₂ SiO ₅ :Ce crystal scint. in liquid scint.
Majorana	Ge-76	500 kg Ge diodes
MOON	Mo-100	Mo sheets between plastic scint., or liq. scint.
Xe	Xe-136	1.56 t of Xe in liq. Scint.
XMASS	Xe-136	10 t of liquid Xe

The Majorana Project

Duke U.

North Carolina State U.

TUNL

Argonne Nat. Lab.

JINR, Dubna

ITEP, Moscow

New Mexico State U.

ORNL

Pacific Northwest Nat. Lab.

U. of Washington

LANL

U. of South Carolina

Brown

Univ. of Chicago

RCNP, Osaka Univ.

Univ. of Tenn.



Ge Basics

Large Mass (~ 1 ton):

Good source radiopurity:

Demonstrated technology:

Natural isotope

Small volume, source = det:

Good energy resolution:

Ease of operation

Large Q value, fast $\beta\beta(0\nu)$

Slow $\beta\beta(2\nu)$ rate

Identify daughter

Event reconstruction

Nuclear theory

500 kg of ^{enr}Ge

Intrinsic Ge

“Ready to Go”

Fiorini “internal source method

3-4 keV at 2039 keV, 0.2%

Ge detectors are relatively simple

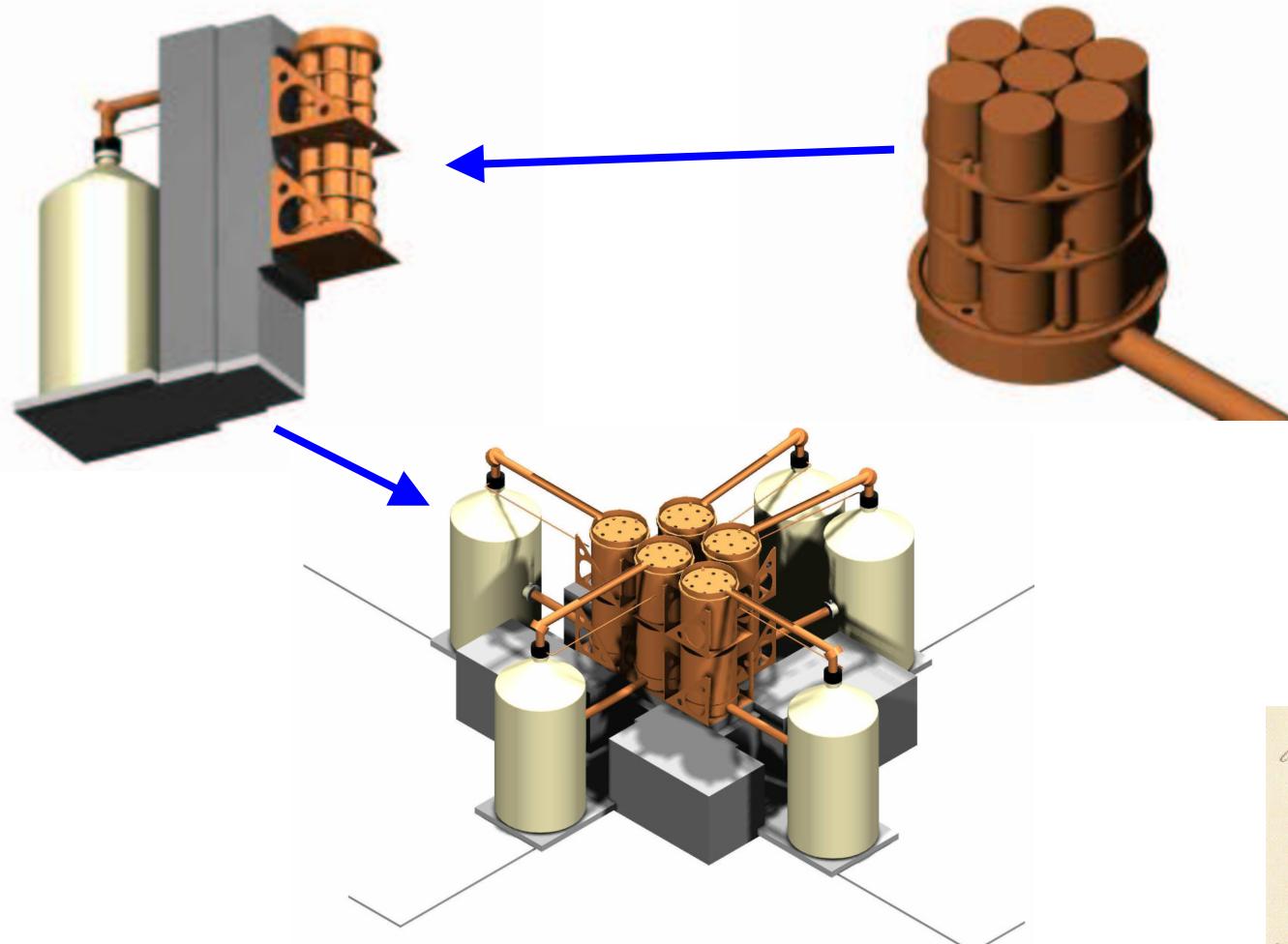
2039 keV, above most radioactivities

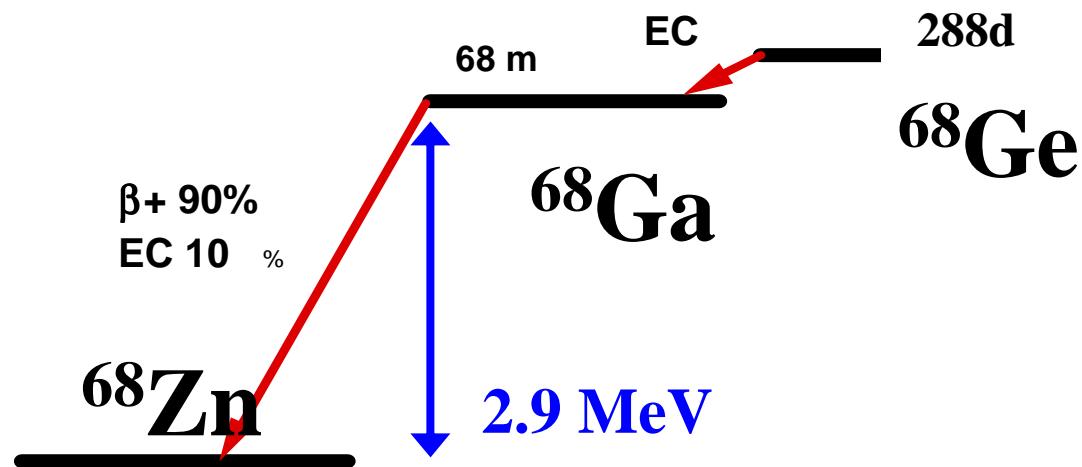
10^{21} yrs

segmentation, modularity, PSD

Low A - Shell Model and QRPA

Majorana Layout



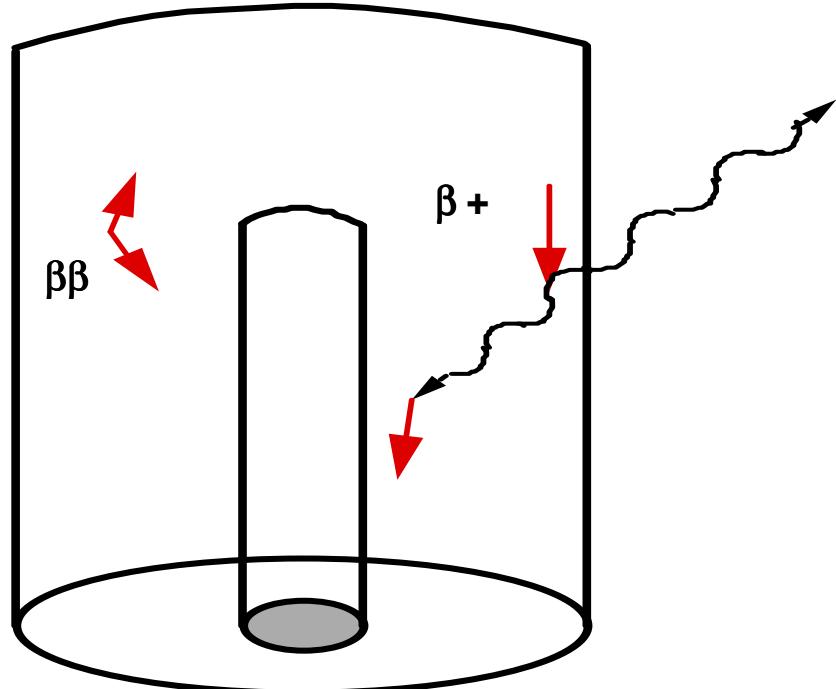


${}^{68}\text{Ge}$ is the dominate background.

For 500-kg enriched detector, initially expect ~ 500 ${}^{68}\text{Ge}$ decays/day. $\tau_{1/2} = 288 \text{ d}$

The naturally occurring ${}^{40}\text{K}$ in the human body decays at a rate of 12000 decays/second.

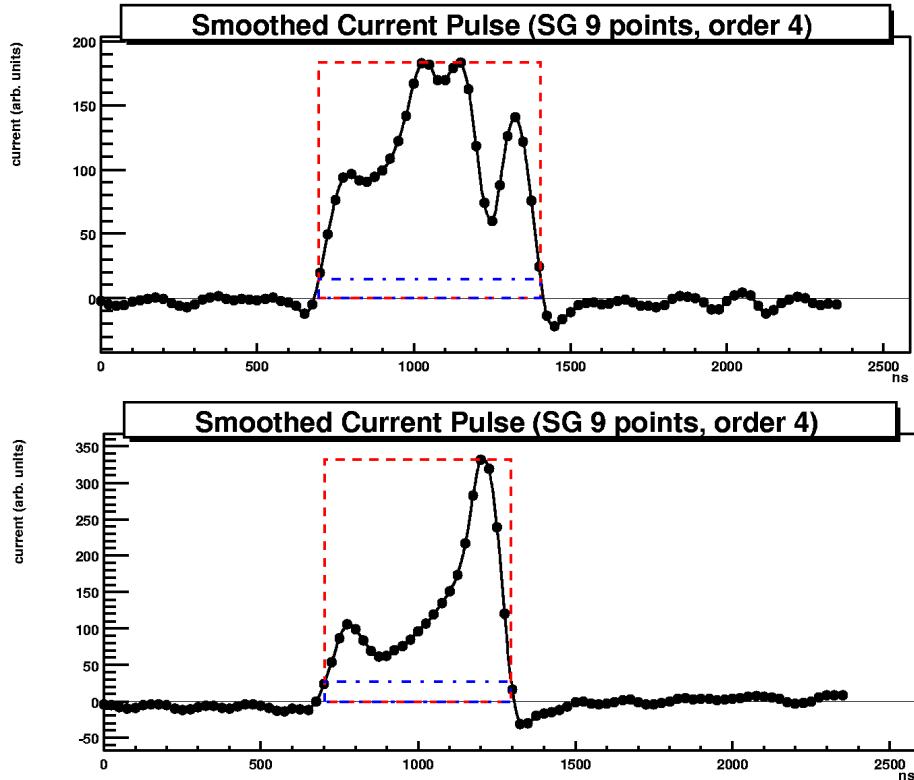
Why Segmentation and Pulse Shape Analysis?



$\beta\beta$ is pointlike.
 β^+ or Compton scattered
 γ rays deposit energy in
multiple locations.

Segmentation and PSD
help reduce these
Backgrounds.

Experimental Examples



Full-energy 1621-keV γ (top) and 1592-keV DEP (bottom) reconstructed current pulses from 120% P-type Ortec HPGe detector (experimental data)

Commercial digital spectroscopy hardware is available with fast (40 MHz), high-resolution (14-bit) digitization
Significant developments in pulse-shape discrimination techniques for HPGe have been made in the past 10 years and are ready to apply to new hardware

Detector Segmentation

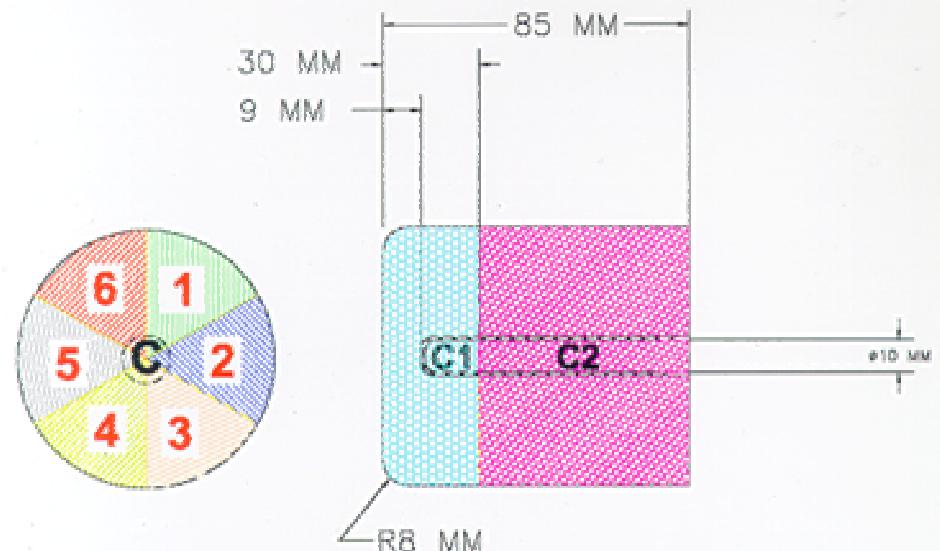
Sensitive to **axial and azimuthal** separation of depositions

Reference design with six azimuthal and two axial contacts is low risk

This level of segmentation gives good background rejection

This segmentation gives us ~2500 segments of 200 g or 40 cc

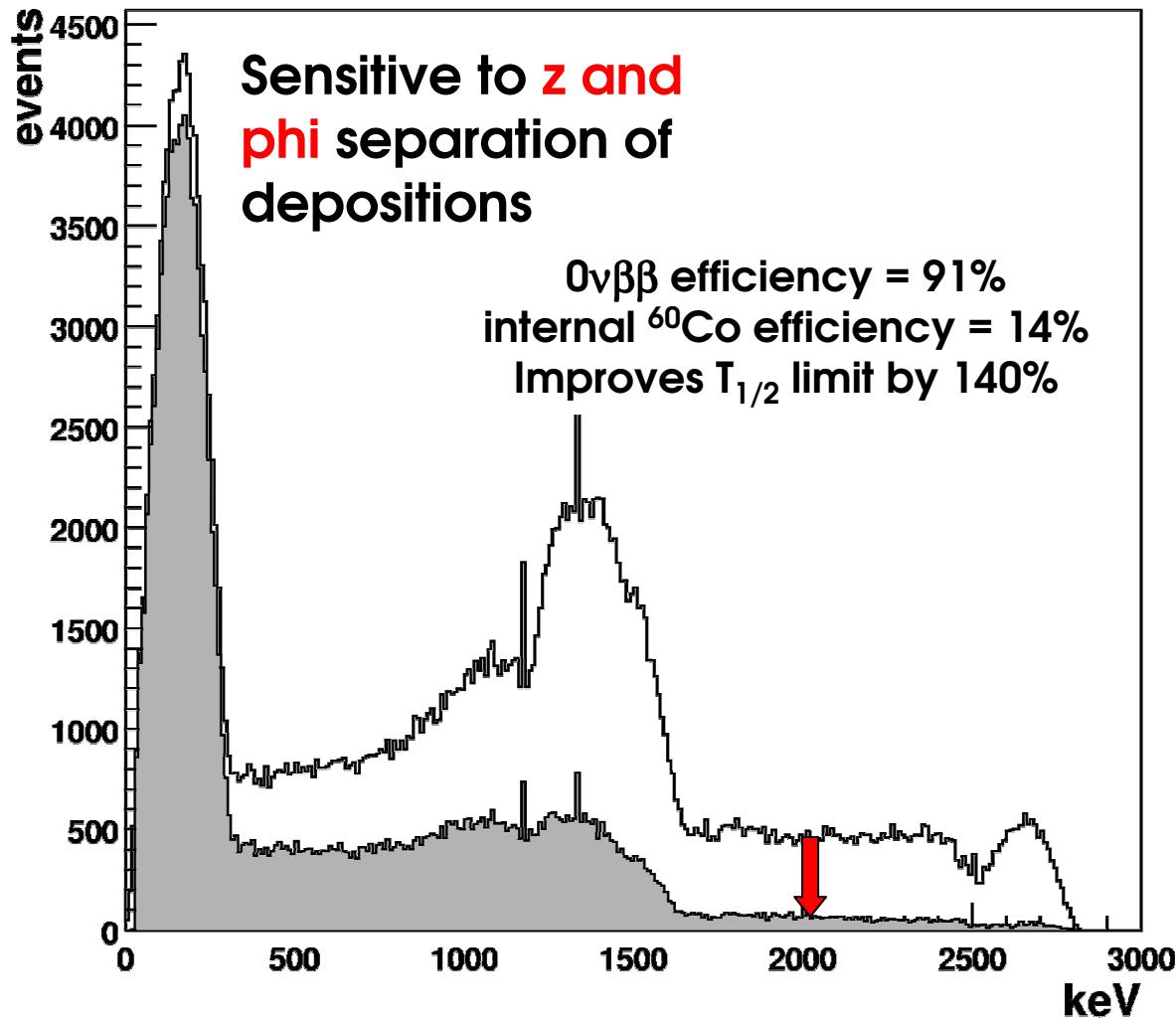
PT6X2
12-SEGMENTS
SEGMENTED DETECTOR
(6-EXTERNAL X 2-INTERNAL)



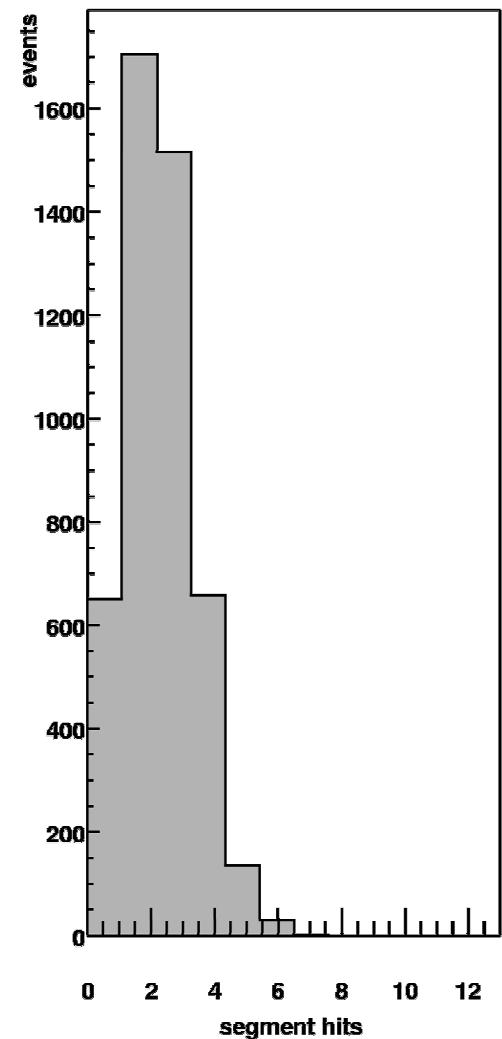
**6 SIDE CHANNELS
2 CENTER CHANNELS
TOTAL = 8 PREAMPLIFIERS**

Monte-Carlo Example

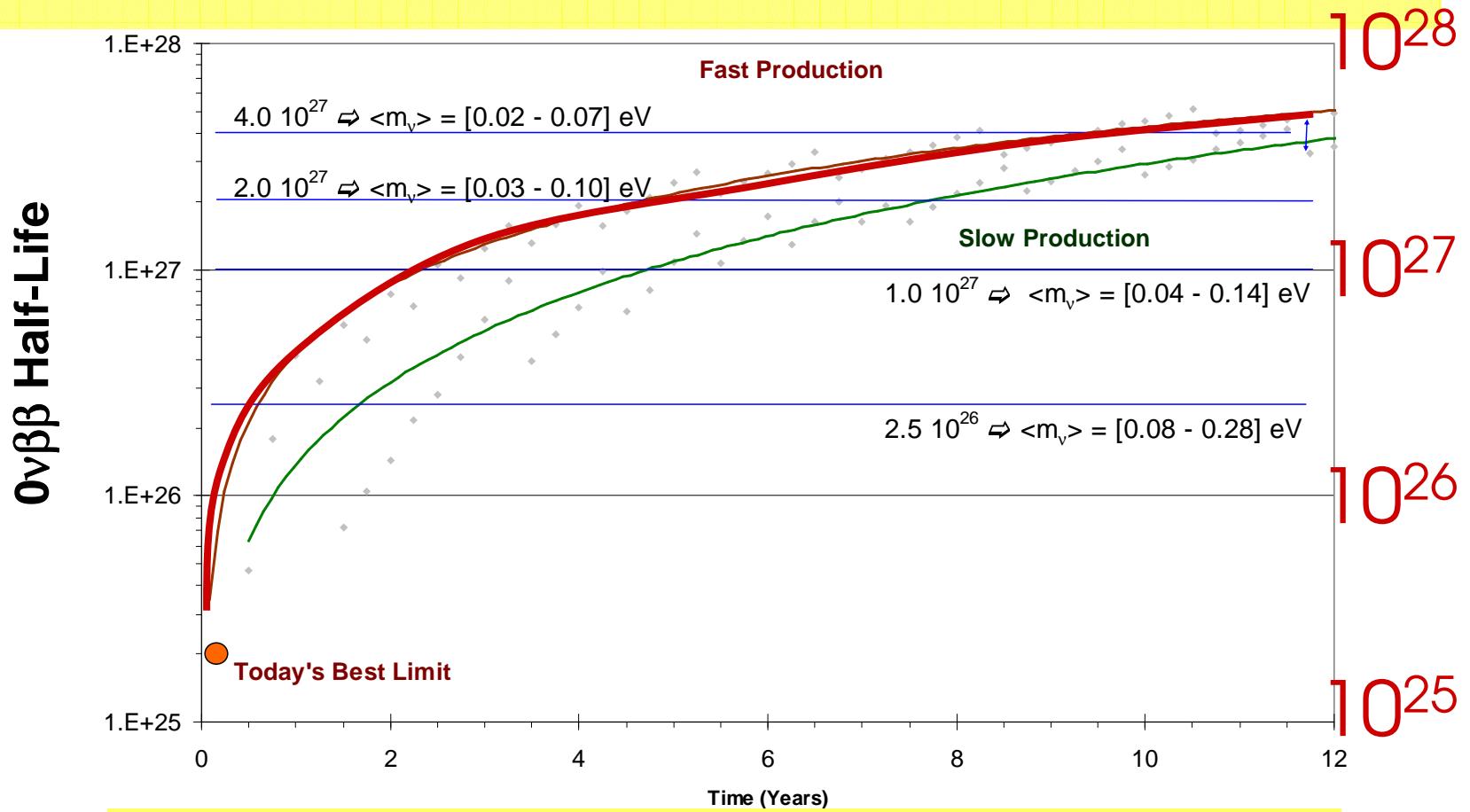
Internal ^{60}Co before and after one-segment cut



Segment multiplicity at 2039 keV

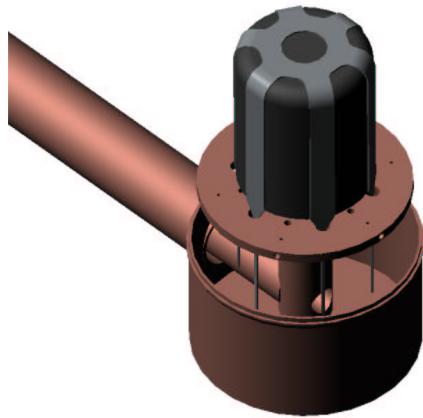


Sensitivity vs. Time



Based on early IGEX background levels with reasonable
background reduction and cutting methods applied

Collaboration Progress: Optimization and Prep for Full Experiment



**Segmented
Enriched
Germanium Array
(SEGА):
Segmented Ge**

1 to 5 Crystals

High energy n bkg

Pulse analysis test

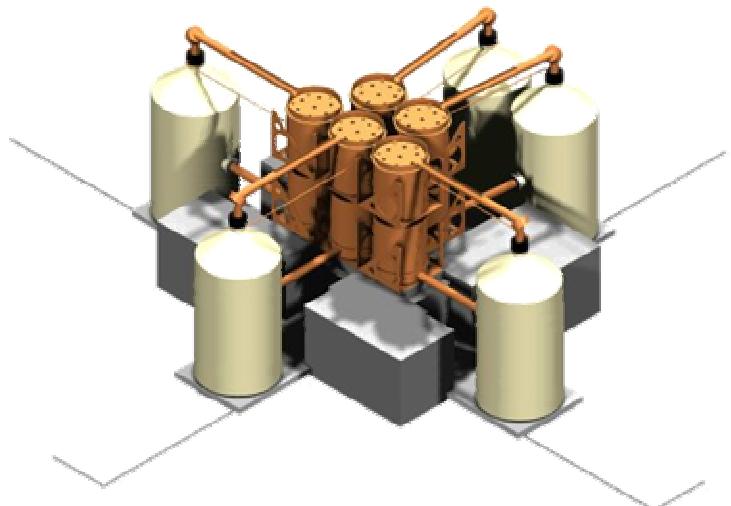
Materials screening

**Multi Element
Germanium Assay
(MEGA):
16+2 natural Ge**



High density
Materials qualification
Cryo design test
Geometry test
Powerful screening tool

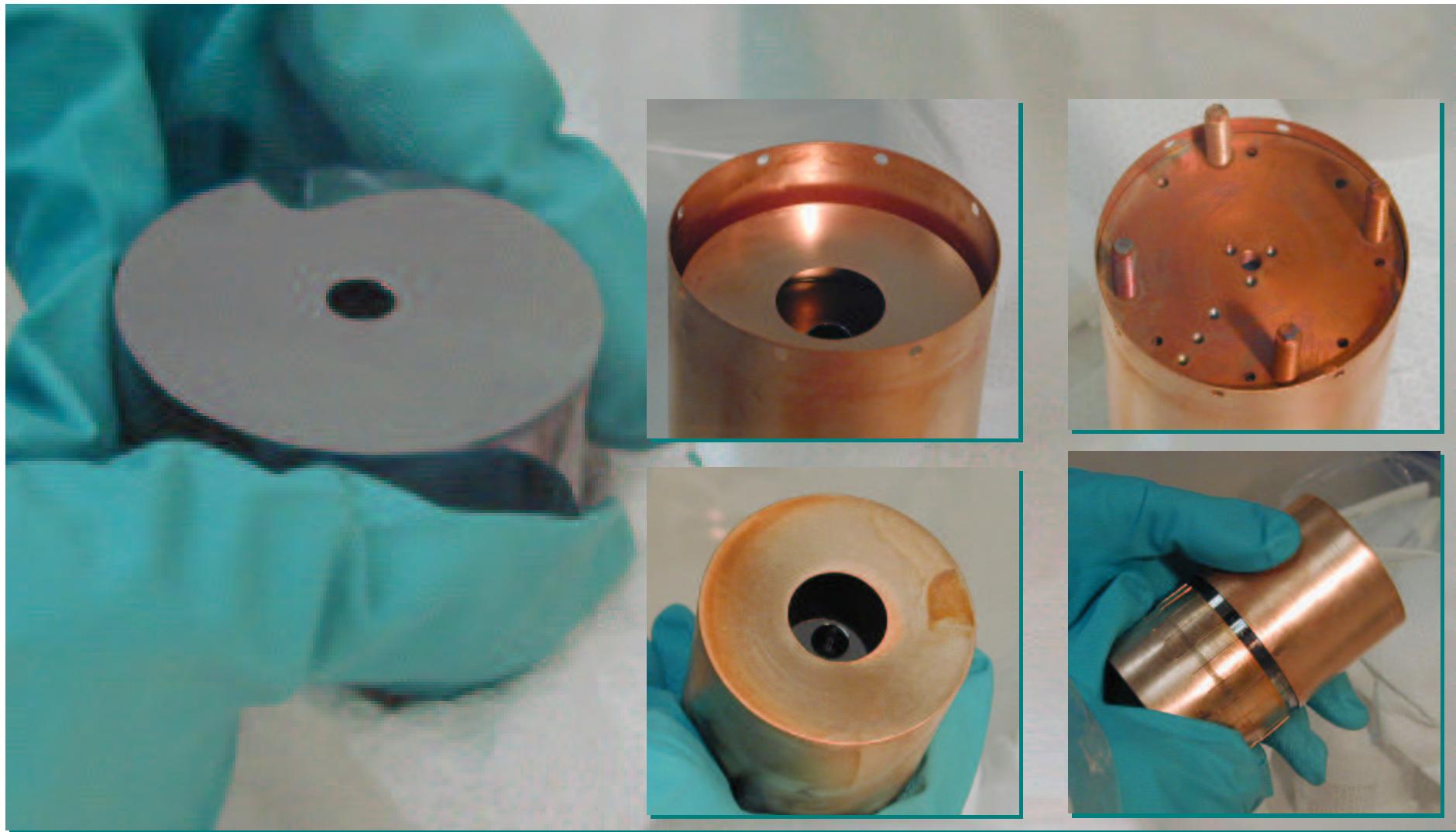
MAJORANA:
Few hundred Ge det.
All enriched/segmented
Ten multi-crystal modules



Full Experiment

Recent Crystal Packaging Test

(*MEGA*)



Summary of Proposals

	Proposed ton-year $= M * T * \epsilon$	Anticipated $\langle m_{ee} \rangle$, (QRPA)
CUORE	$0.21*5*1 = 1$	60 meV
EXO	$6.5*10*0.7 = 45$	13 meV
GENIUS	$1*2*1 = 2$	20 meV
MAJORANA	$0.5*10*1 = 5$	25 meV
MOON	$3.3*3*0.14 = 1.4$	30 meV

The $\langle m_{\beta\beta} \rangle$ limits depend on background assumptions and matrix elements which vary from proposal to proposal.

An exciting time for $\beta\beta$!

For at least
one neutrino:

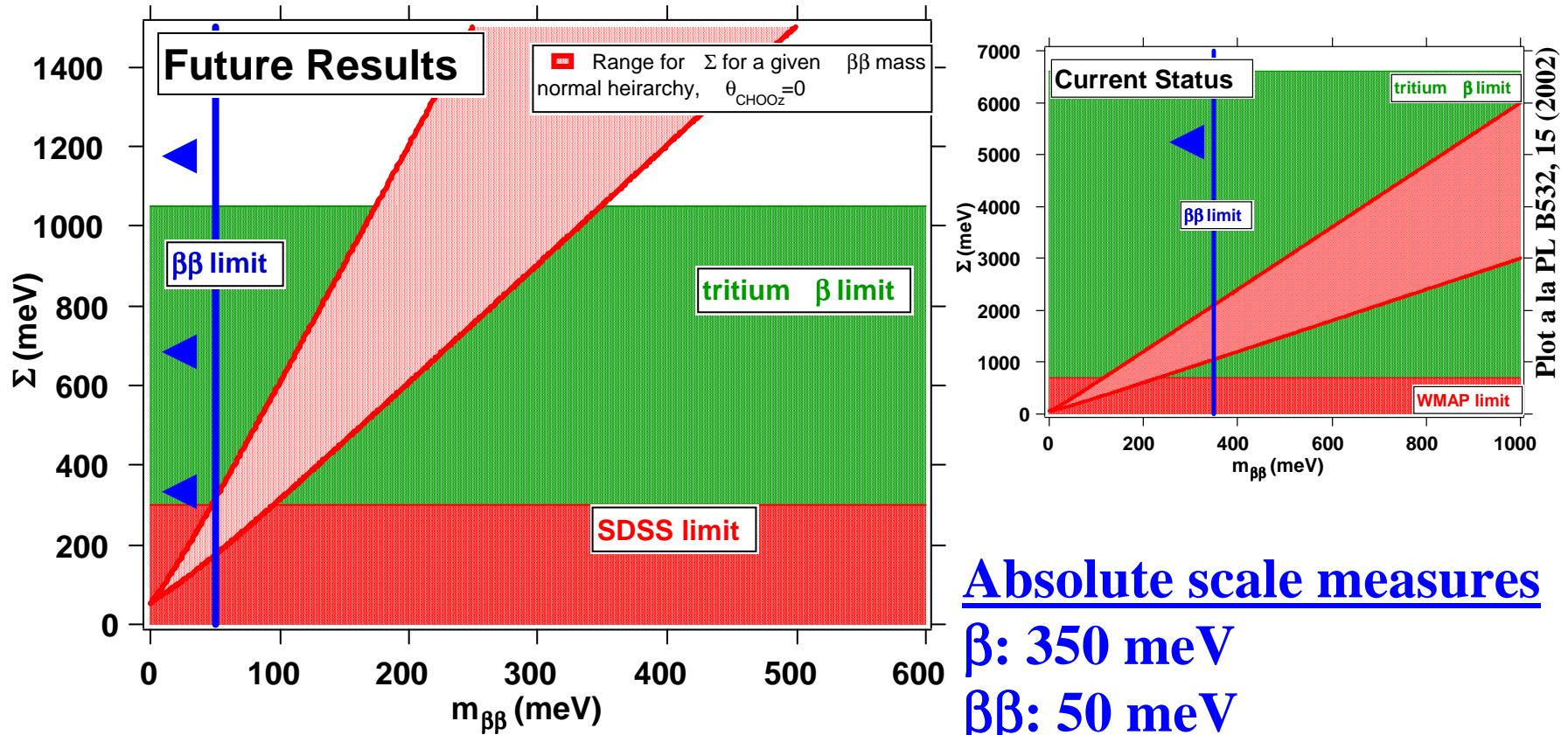
$$m_i > \sqrt{\delta m_{atmos}^2} \approx 50 \text{ meV}$$

For the next experiments:

$$\langle m_{\beta\beta} \rangle \leq 50 \text{ meV}$$

**$\langle m_{\beta\beta} \rangle$ in the range
near 50 meV is very interesting.**

Summary of Mass Measurements (with a guess at the future)



Absolute scale measures
 β : 350 meV
 $\beta\beta$: 50 meV
Cosmology: <100 meV